Including relativistic and primordial Non-Gaussianity contributions in cosmological simulations by modifying the initial condition

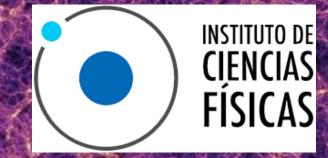


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**Bispectrum Power spectrum** Tree level 1.050 10<sup>9</sup> B Rc ng B Rc 1.025 ΒG B(Δk, k, k)[h <sup>-6</sup>Mpc<sup>6</sup>] 108 1.000 Sim Rc ng Sim Rc - Sim G 0.975 P/P<sub>nw</sub> Pert Rc ng 107 Pert Rc 0.950 Linear ٠ Pert PNG+c<sub>s</sub><sup>2</sup> 0.925 GRamses 4 106 0.900 0.875 105 16  $|\frac{P-P_G}{P_G}| \times 100$  $\left|\frac{B-B_G}{2}\right| \times 100$ 14 Pert Rc\_ng Pert Rc ng 12 Sim Rc\_ng --- Sim Rc BG 10 BRc ng B Rc 8 Tree level 0.025 0.050 0.075 0.100 0.125 0.150 0.175 0.200 6 .  $10^{-2}$ 10-1 k [Mpc / h]<sup>-1</sup>

*k* [Mpc / h]<sup>-1</sup> M. Enríquez, J. C. Hidalgo and O. Valenzuela, Cosmological simulations with relativistic and primordial non-Gaussianity contributions as initial conditions, 2109.13364.

### Newtonian and relativistic correspondance

#### Newtonian

#### Relativistic

$$\frac{d\delta_N}{d\tau} = -(1+\delta_N)\nabla^2\nu_N$$

$$\frac{d(\nabla^2 \nu_N)}{d\tau} + \mathcal{H} \nabla^2 \nu_N + \partial^i \partial_k \nu_N \partial^k + \partial_i \nu_N + 4\pi G a^2 \bar{\rho} \delta_N = 0$$

$$\delta' = -(1+\delta)\theta$$

$$\theta' + \mathcal{H}\theta + \theta^i_j \theta^j_i + 4\pi G a^2 \bar{\rho} \delta = 0$$

$$\theta = \nabla^2 \nu_N = \nabla^2 \nu_l^{(1)}$$
$$\theta_k i = \partial^i \partial_k \nu_N = \partial^i \partial_k \nu_l^{(1)}$$

#### **GR.** Constrain equation

$$\theta^2 - \theta^i_j \theta^j_i + 4\mathcal{H}\theta + R = 16\pi G a^2 \bar{\rho}\delta$$

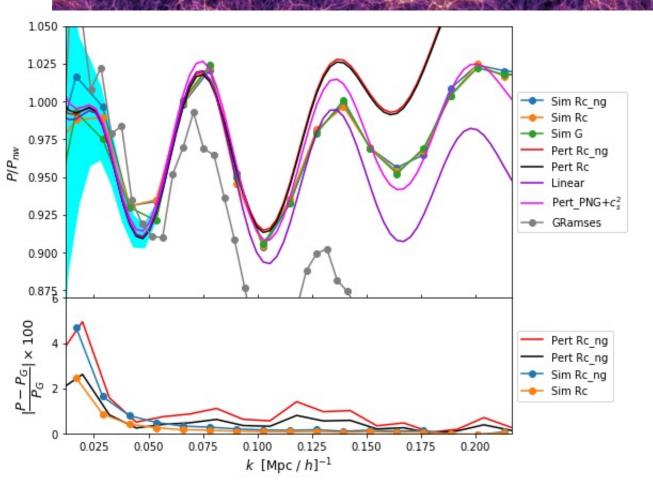
# Density field with GR and N-G contributions

$$\delta = \delta^{(1)} + \frac{1}{2}\delta^{(2)} + \frac{1}{6}\delta^{(3)} + \dots$$

$$\frac{1}{2}\delta^{(2)} = \frac{D_{+}(\eta)}{10\mathcal{H}^{2}D_{+IN}}\frac{24}{5}\left[-(\nabla\zeta^{(1)})^{2}\left(\frac{5}{12}+f_{\rm NL}\right)+\zeta^{(1)}\nabla^{2}\zeta^{(1)}\left(\frac{5}{3}-f_{\rm NL}\right)\right]$$

$$\frac{1}{6}\delta^{(3)} = \frac{D_{+}(\eta)}{10\mathcal{H}_{IN}^{2}D_{+IN}}\frac{108}{25} \left[2\zeta^{(1)}(\nabla\zeta^{(1)})^{2}\left(-g_{\rm NL} + \frac{5}{9}f_{\rm NL} + \frac{25}{54}\right) + \zeta^{(1)2}\nabla^{2}\zeta^{(1)}\left(-g_{\rm NL} + \frac{10}{3}f_{\rm NL} - \frac{50}{27}\right)\right]$$

## Results



- We computed the difference with no-wiggle power spectrum.
- The blue shadow corresponds to a Euclid like observed area.
- The difference in the supression with GRAMSES is due to the counter term\* as:

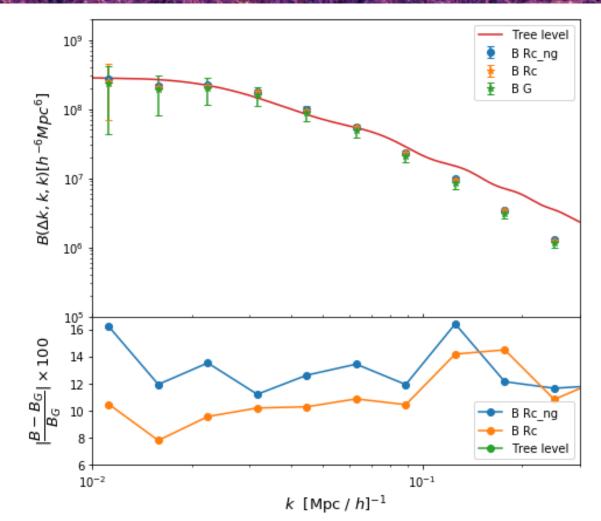
$$P_{ctr,1loop} \equiv -2k^2 c_s^2 P_{11}$$

 PNG and relativistic contibutions difference up to 4.5% (fnl=-4.2 gnl=-7000)

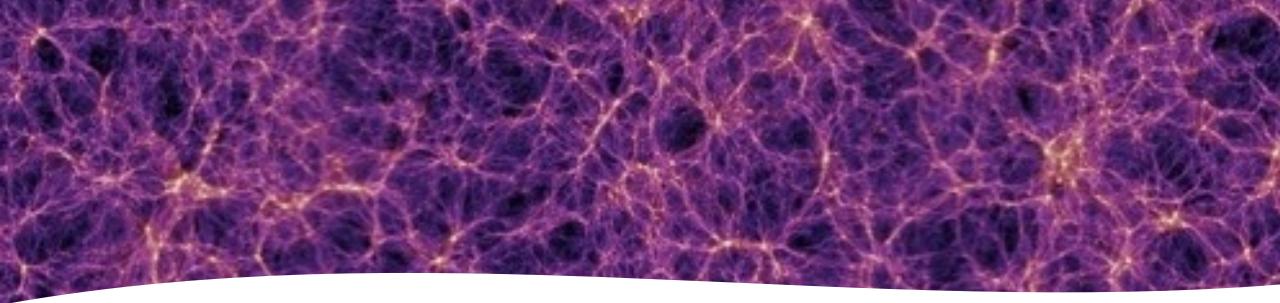
\*T. Baldauf, L. Mercolli, and M. Zaldarriaga, "Effective field theory of large scale structure at two loops: The apparent scale dependence of the speed of sound," Phys. Rev. D, vol. 92, p. 123007, Dec 2015.

### Results

- We computed the bispectrum and compare it with the Tree level bispectrum,, where we found a maximum difference por relativistic contributions with PNG at large scales of 16%
- This calculations were made in Pylians\*
- The difference at high k's (i.e. small scales) are because we used just one realization, further tests are needed.



\*F. Villaescusa-Navarro, Pylians:Python libraries for the analysis of numerical simulations,2018. ascl: 1811.008



Summary and future work

- Can be added relativistic corrections at initial conditions with Newtonian evolution.
- The non-Gaussianity and the relativistic corrections are observable at high order statistics as the bispectrum.
- There are different triangulations in the bispectrum that are need to be take into account.
- Look at the effects on simulations with larger sensitivity to identify physical effects.
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