



Large Spin-2 Signals at the Cosmological Collider (CC)

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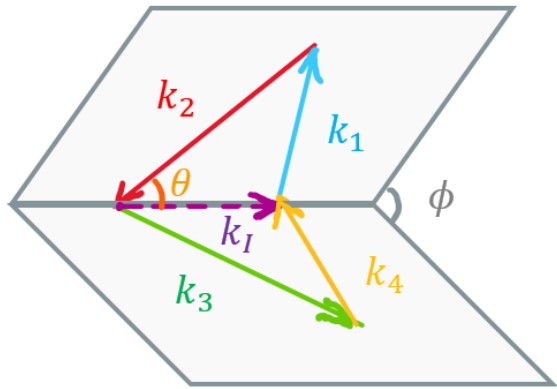
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Cosmological collider physics: ideal vs challenge



$$\langle \zeta^4 \rangle' \sim \lambda^2 \left(\frac{k_1 + k_2}{k_I} \right)^{i\sigma} P_S(\cos \theta) f(\phi) \times e^{-\pi m/H}$$

Labels for the equation:

- λ^2 : coupling
- $\frac{k_1 + k_2}{k_I}$: mass
- $i\sigma$: decay width
- $P_S(\cos \theta)$: spin
- $f(\phi)$: P&CP
- $e^{-\pi m/H}$: Boltzmann suppression (circled in red)

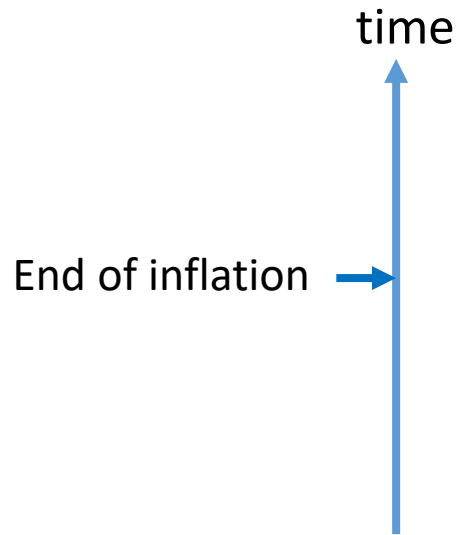
CC Signal strength:
Boltzmann suppression



Hard to observe heavy particles (e.g., $m > 5H$)



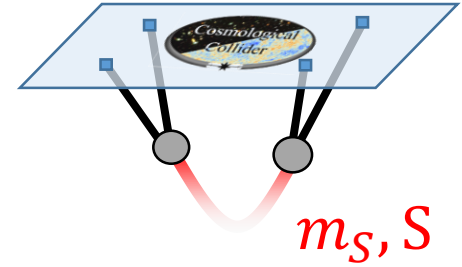
Amplification?



$$H \sim 10^{13} \text{ GeV}$$

[Chen & Wang, 2009]
 [Baumann & Green, 2011]
 [Noumi, Yamaguchi & Yokoyama]
 [Arkani-Hamed & Maldacena, 2015]

Large signal requires symmetry breaking



- Signal strength $\sim \sqrt{\langle n \rangle}$ ~~dS symmetries~~ $\rightarrow f(m_S) = e^{-\pi \frac{m_S}{H}} \ll 1$ for heavy particles



- Secret: Symmetries

$$dS: SO(4,1) = \begin{cases} 3 \text{ translations } P_i \\ 3 \text{ rotations } J_i \\ 1 \text{ dilation } D \\ \text{3 dS boosts } K_i \end{cases}$$



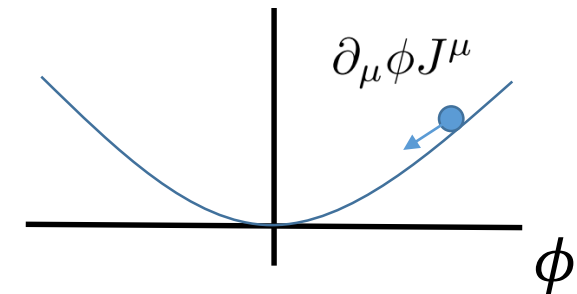
Dictates dS-invariant particle production $\sim e^{-\pi \mu_S} \sim e^{-\pi m_S/H}$

Broken by the rolling inflaton ϕ_0



$$f(m_S, \phi_0) \rightarrow e^{-\pi \frac{m_S - \phi_0/H}{H}}$$

Exponential amplification via a **chemical potential**



Chemical potential: A brief review

Q : What do we mean by chemical potential?

Statistical mechanics

$$Z = e^{-(H - \kappa N)/T}$$

+ Locality & Lorentz covariance

[Wang & Xianyu, 2020]
[Sou, Tong & Wang, 2021]

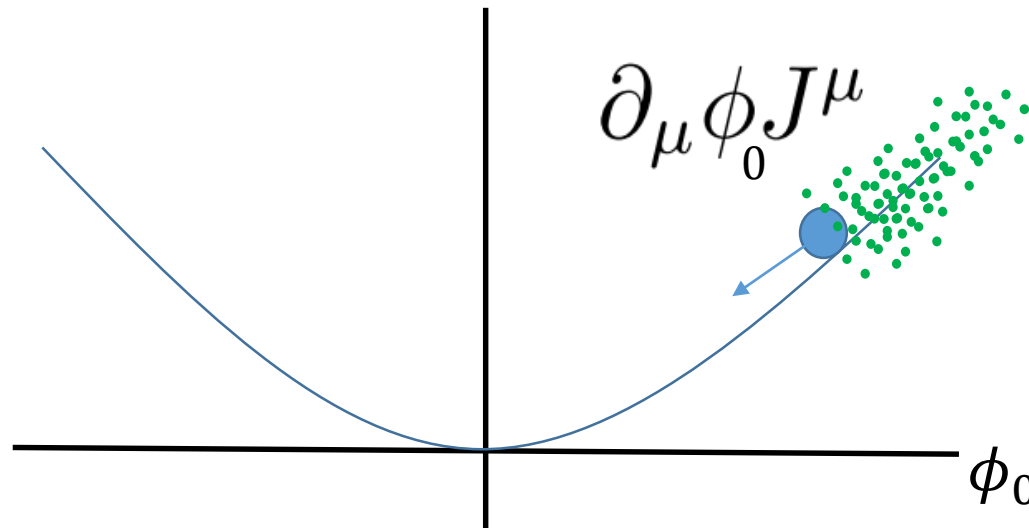
$$Z = \int \mathcal{D}\phi e^{i \int dt d^3x \sqrt{|g|} (\mathcal{L}(\phi, \partial\phi) + \Delta\mathcal{L}_{\text{chem}})}$$

with $\Delta\mathcal{L}_{\text{chem}} \equiv \kappa_\mu(x) J^\mu(x)$

A : A local operator in the Lagrangian that enhances particle production

Chemical potential by a background field

E.g., $\kappa_\mu \propto \partial_\mu \phi_0$



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- Chemical potential of (massive) lower spin particles are known:

{	$S = 0$	Schwinger effect	[Kobayashi & Afshordi, 2014] [Chua, Ding, Wang, & Zhou, 2018]
	$S = \frac{1}{2}$	Axial current	[Adshead & Sfakianakis, 2015] [Chen, Wang, & Xianyu, 2018]
	$S = 1$	Chern-Simons	[Barnaby & Peloso, 2011] [Liu et al., 2019] [Wang & Xianyu, 2020]



Huge **amplification** of CC signal

$$SNR \sim e^{-\pi m} \rightarrow e^{-\pi (m - \kappa)}$$

Reminder: Inflation may not be that far from the GUT/KK/string scale

How about higher spins? In particular, $S = 2$?

KK gravitons from GUT

[Kumar & Sundrum, 2019]

Open string spectrum

[Kim, Noumi, Takeuchi & Zhou, 2019]

[Noumi, Takeuchi & Zhou, 2020]

A spin-2 chemical potential?

- What do we need?
 1. An EFT operator with lowest possible dimension
 2. In the form $\Delta\mathcal{L}_{\text{chem}} \equiv \kappa_\mu J^\mu(h)$
 3. Quadratic: $O(h_{\mu\nu}^2)$
 4. Break dS boosts but not rotations, translations and dilation
 5. Give rise to a consistent ghost-free theory
- The **unique** spin-2 chemical potential:



$$S_c = \int d^4x \sqrt{-g} \left[\frac{\phi}{2\Lambda_c} \mathcal{E}^{\mu\nu\rho\sigma} \nabla_\mu h_{\nu\lambda} \nabla_\rho h_\sigma^\lambda \right] .$$

The spin-2 chemical potential

- The linear theory

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{4} h \partial^2 h + \frac{1}{4} m^2 h^2 + \frac{\phi}{2\Lambda_c} \epsilon \partial h \partial h \right]$$

Einstein-Hilbert Fierz-Pauli Chemical potential $\kappa = \frac{\dot{\phi}_0}{\Lambda_c}$

↖ ↑ ↗

↓

EoM: $(\square - (m^2 + 2H^2)) h_{\mu\nu} - \frac{\nabla_\alpha \phi}{\Lambda_c} \mathcal{E}^{\alpha\kappa\rho\sigma} (g_{\kappa\nu} \nabla_\rho h_{\mu\sigma} + g_{\kappa\mu} \nabla_\rho h_{\nu\sigma}) = 0,$

Constraints: $\nabla^\mu h_{\mu\nu} = 0, \quad h = 0.$

- Naïvely, DoF = 10 - 4 - 1 = 5?
- A non-local constraint kills the vector modes: $i\kappa a^2 h_{\pm 1} = 0.$ ➡ DoF = 10 - 4 - 1 - 2 = 3

- Ghost-free $\sqrt{\quad}$

$$\langle n_\lambda(\mathbf{k}) \rangle' \xrightarrow{\mu \gg 1} \begin{cases} e^{-\pi(\mu + \tilde{\kappa})}, \lambda = +2 & \text{(Suppressed)} \\ \text{☒}, \lambda = +1 \\ e^{-\pi\mu}, \lambda = 0 & \text{(Unchanged)} \\ \text{☒}, \lambda = -1 \\ e^{-\pi(\mu - \tilde{\kappa})}, \lambda = -2 & \text{(Enhanced)} \end{cases}$$

μ : mass
 $\tilde{\kappa}$: chemical potential

➡ CC signals \uparrow 😊 ~~Parity~~

Coupling to the visible sector

[Cheung et al., 2008]

[Lee, Baumann & Pimentel, 2016]

- Couple to inflatons and gravitons in the EFT of inflation

$$S_{h\gamma} = \int d\tau d^3x \rho \phi' \gamma'_{ij} h_{ij} ,$$

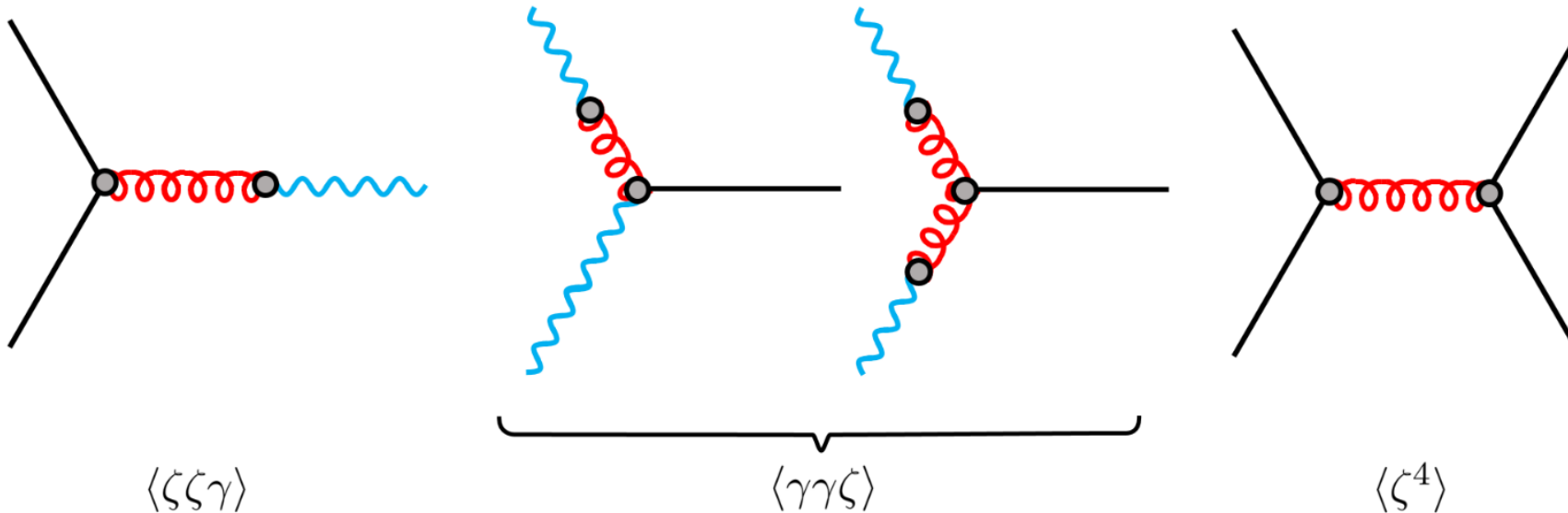
$$S_{h\phi^2} = - \int d\tau d^3x \frac{1}{M} h_{ij} \partial_i \phi \partial_j \phi .$$

➤ Dim-5

➤ Shift symmetric: $\phi \rightarrow \phi + C$

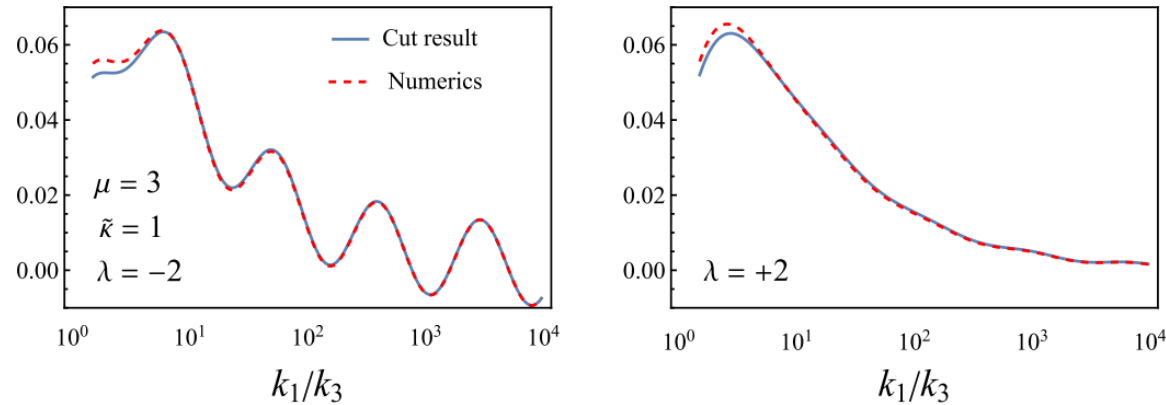
➤ ~~Spin-2 gauge & time diff~~

- Large non-Gaussianities from chemical potential

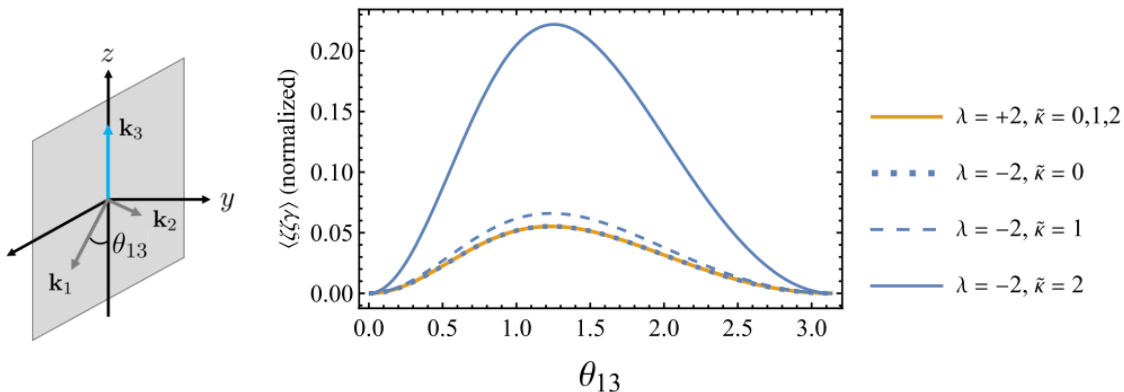


Summary of features in observables

- Helicity-dependent CC signals

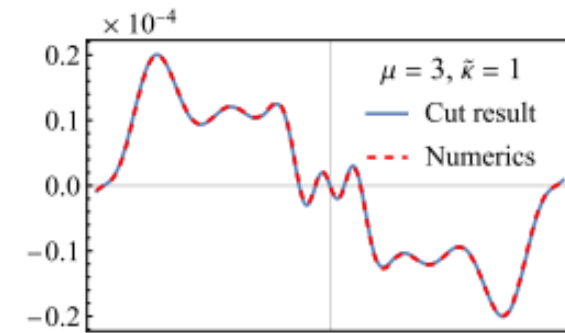


- Characteristic angular dependence

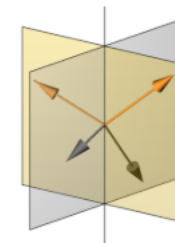


- Imaginary “decay plane” correlation

[Liu, Tong, Wang & Xianyu, 2019]



ψ_I



Non-perturbative ***imaginary*** part:

$$\text{Im} \langle \zeta^4 \rangle \propto e^{-\pi\mu} \sinh \pi\tilde{\kappa} \sin \psi_I$$

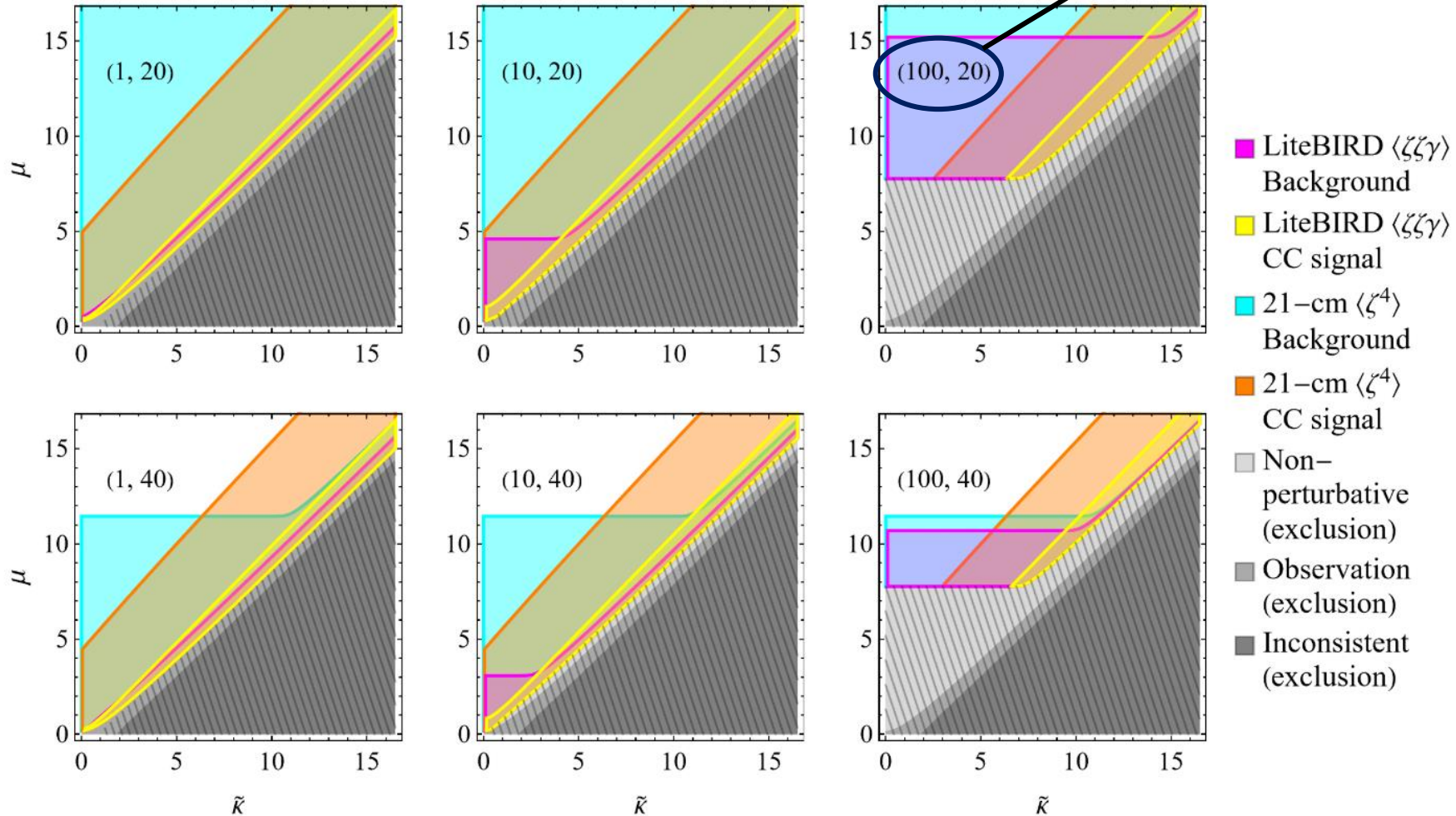
Can we see them?

Magic of the chemical potential:

Heavy spin-2 particles

$(\rho, M/H)$

- Parameter space for future CMB & 21-cm observations



[Shiraishi, 2019]
 [Flöss, Wild, Meerburg & Koopmans, 2021]

- ✓ CC physics
- ✓ Symmetry suppresses signal
- ✓ Amplification via chemical potential
- ✓ Generalization to spin-2
- ✓ Features observable in future experiments

Summary and outlooks

- ❑ Non-perturbative mixing between graviton and the spin-2 particle?
- ❑ Apparent discontinuity in the DoF count?
- ❑ UV completions?
- ❑ Going beyond spin-2?
- ❑ Probing the Regge trajectory?



Thanks for
listening!

Backup pages