# Cosmology with Marked Power Spectra

Elena Massara



**Cosmology from Home 2022** 

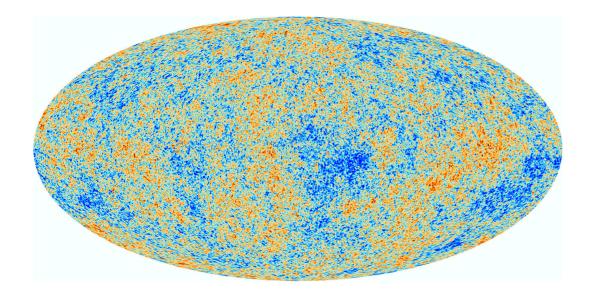
### THE LARGE SCALE STRUCTURE

### The distribution of matter in the Universe is sensitive to:

- properties of dark matter
- nature of dark energy
- neutrino mass scale
- initial condition of the Universe

 $\Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu$ 

### NON-GAUSSIAN density field



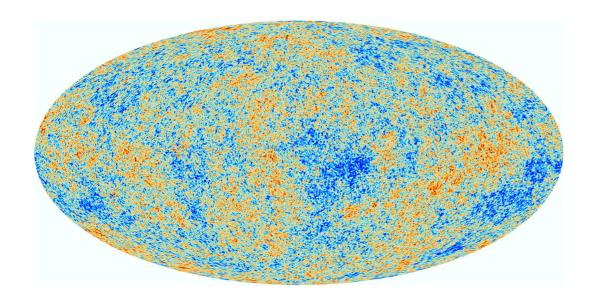
### $\delta(\mathbf{k}) \sim N(0, P(\mathbf{k}))$

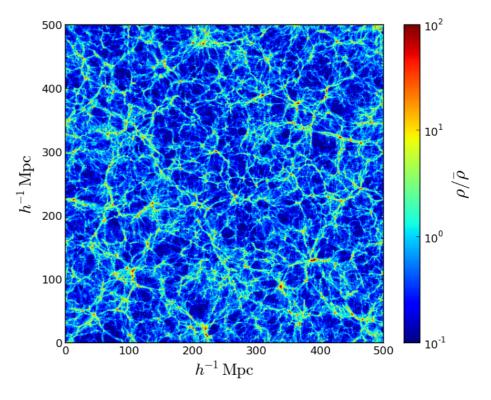
All information contained in 2-pt statistics:

- correlation function
- power spectrum

Higher order statistics are not needed to describe the field

### NON-GAUSSIAN density field





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NOT all information contained in 2-pt statistics

 $\delta(\mathbf{k}) \nsim N(0, P(\mathbf{k}))$ 

Higher order statistics contain information to describe the field

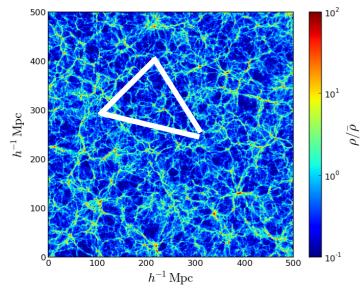
### **NON-GAUSSIAN** statistics

A variety of statistics have been proposed to retrieve the cosmological information beyond the two point functions

#### **Higher-order statistics:**

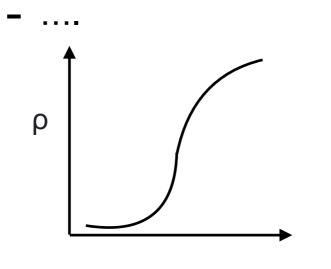
- bispectrum
- trispectrum





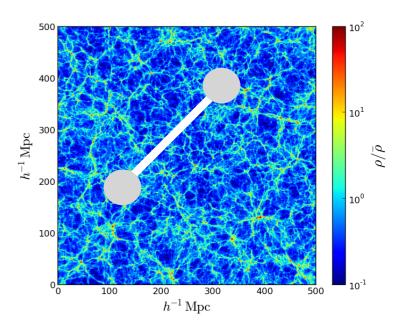
### Different summary statistics:

- peaks
- voids
- scattering transforms
- minimum spanning tree



#### Non-linear transformations of the field:

- log-transformations
- <u>marked power</u>
   <u>spectra</u>



#### And many others ...

### COSMOLOGY with VOIDS

## Low-density regions are good laboratories to study cosmology because

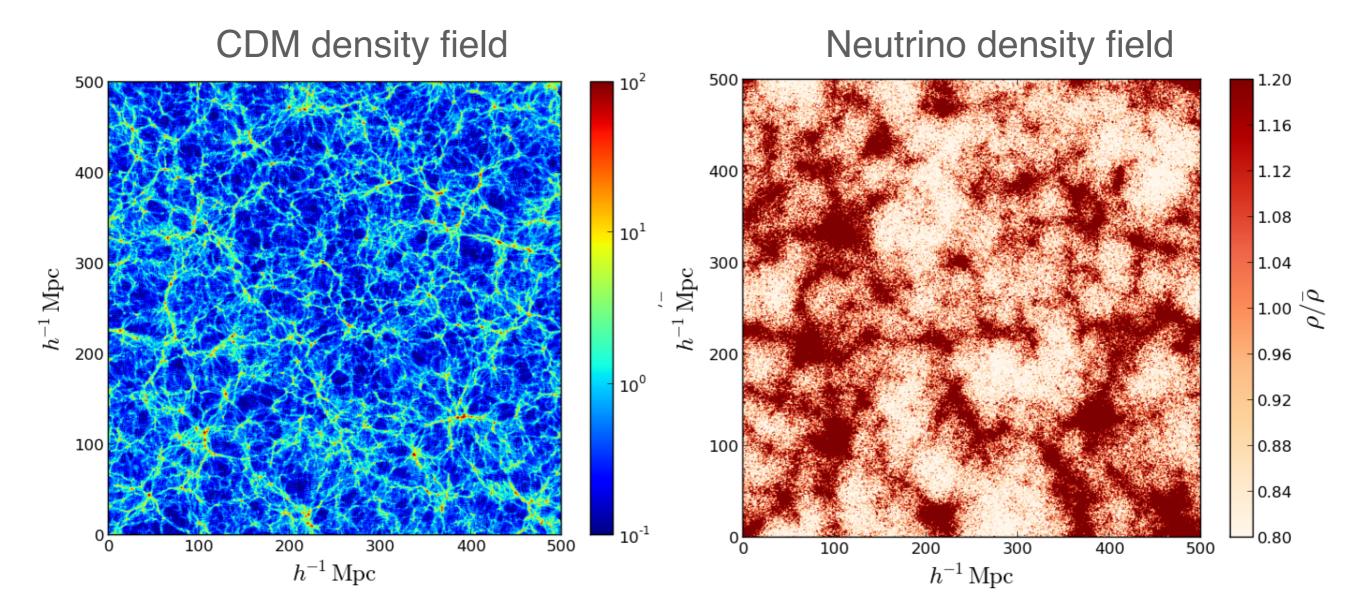
• They are unvirialized, thus they are expected to retain most of their initial cosmological information

### COSMOLOGY with LOW-DENSITY regions

## Low-density regions are good laboratories to study cosmology because

- They are unvirialized, thus they are expected to retain most of their initial cosmological information
- They are sensitive to diffuse components such as
  - neutrinos
  - dark energy

### COSMOLOGY with LOW-DENSITY regions



### COSMOLOGY with LOW-DENSITY regions

## Low-density regions are good laboratories to study cosmology because

- They are unvirialized, thus they are expected to retain most of their initial cosmological information
- They are sensitive to diffuse components such as
  - neutrinos
  - dark energy
- Screening mechanism are inefficient in them

### LOW DENSITY REGIONS

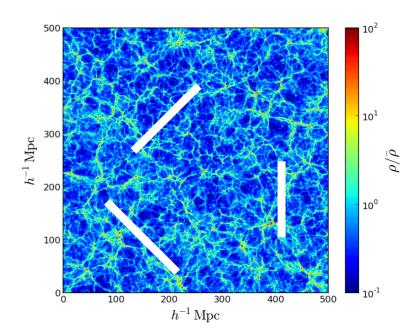
Low-density regions are good probe to study cosmology

- 1. Do 2-pt functions depend on low-density regions?
- 2. Can we modify standard 2-pt functions to incorporate more information from low-density regions?

### **CORRELATION FUNCTION**

#### **Correlation function**

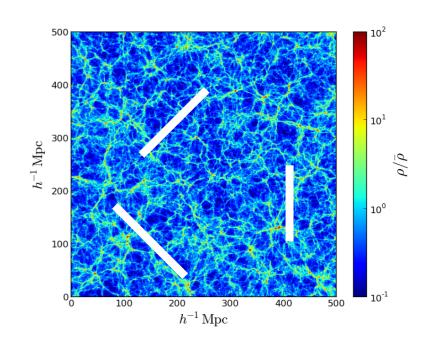
$$1 + \xi(r) = \frac{V}{N^2} \sum_{i,j=1}^{N} \delta_{D}(|\vec{x_i} - \vec{x_j}| - r)$$



### **CORRELATION FUNCTION**

#### **Correlation function**

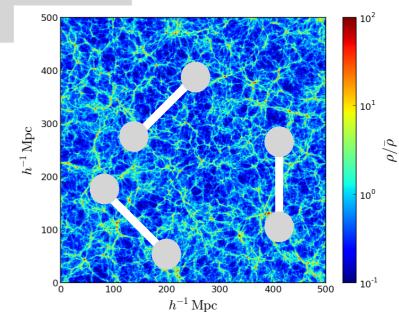
$$1 + \xi(r) = \frac{V}{N^2} \sum_{i,j=1}^{N} \delta_{D}(|\vec{x_i} - \vec{x_j}| - r)$$



#### Marked correlation function

$$1 + M(r,\phi) = \frac{V}{N^2} \sum_{i,j=1}^{N} \frac{\delta_{D}(|\vec{x_i} - \vec{x_j}| - r) m(\vec{x_i},\phi) m(\vec{x_j},\phi)}{\bar{m}^2} = \bar{m}^{500}$$

- 1. m depends on the local density around each point
- 2. m up-weights low-density regions and down-weights high-density regions



### MARKED CORRELATION FUNCTION

- -

$$1 + M(r,\phi) = \frac{V}{N^2} \sum_{i,j=1}^{N} \frac{\delta_{D}(|\vec{x_i} - \vec{x_j}| - r) m(\vec{x_i},\phi) m(\vec{x_j},\phi)}{\bar{m}^2}$$

$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1+\delta_s}{1+\delta_s+\delta_R(\vec{x})}\right]^p$$

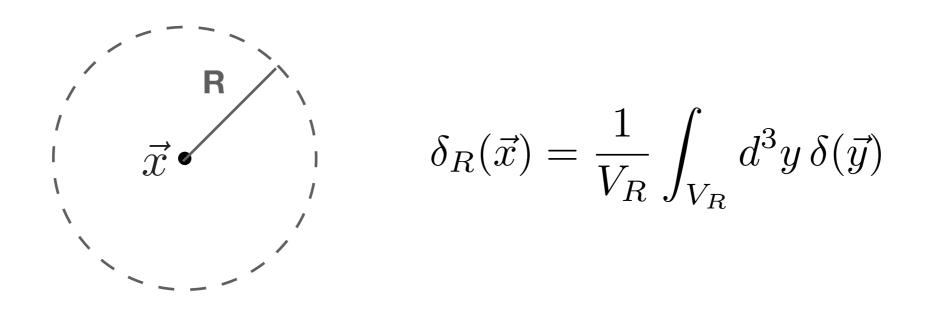
(M. White 2016)

### MARKED CORRELATION FUNCTION

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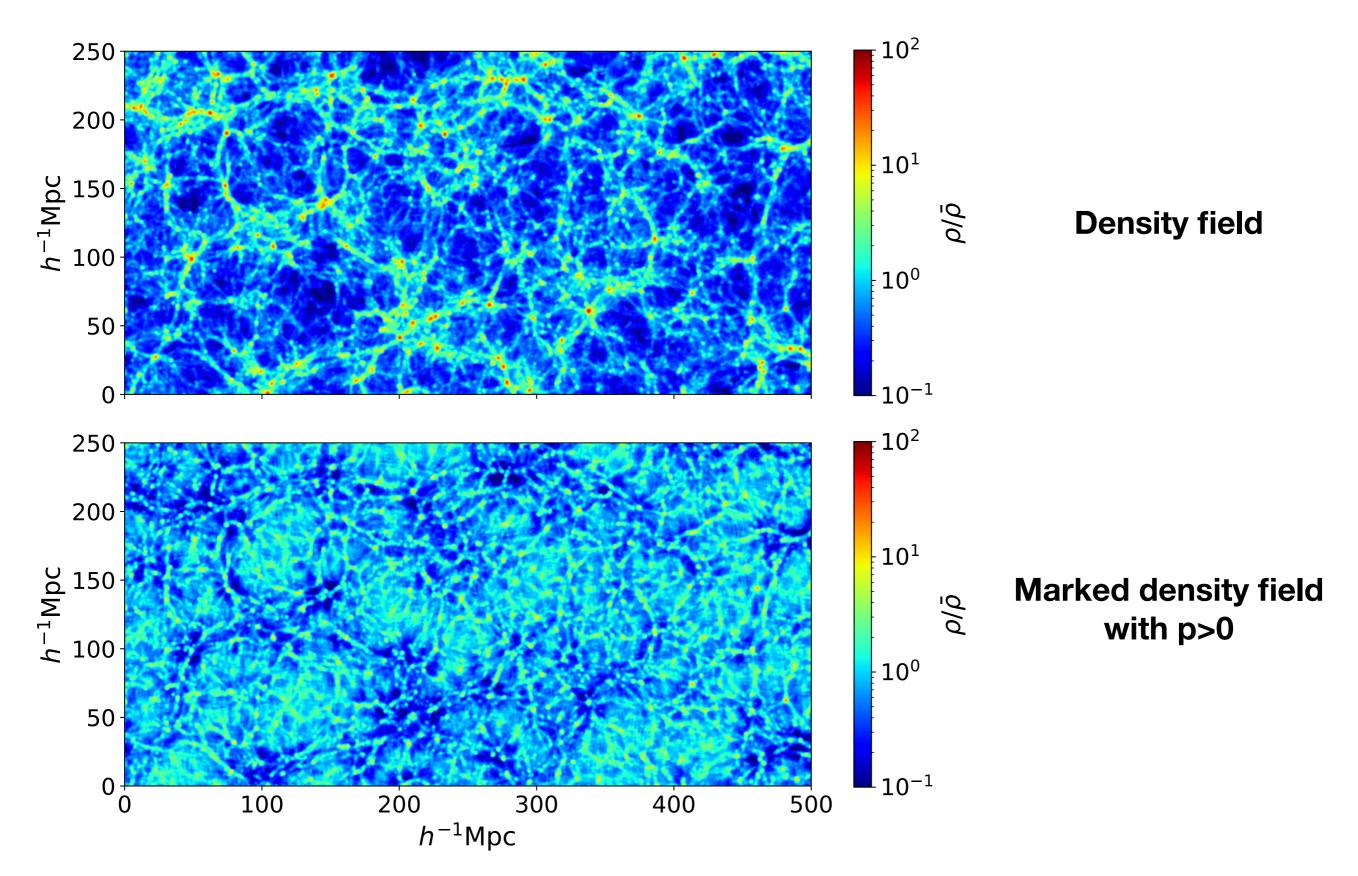
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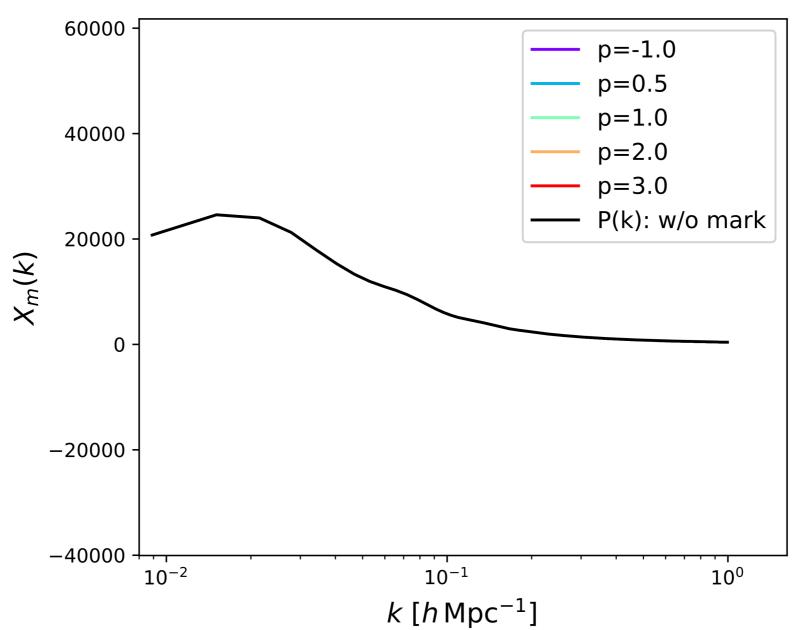
p > 0 up-weight galaxies in **low** density regions p < 0 up-weight galaxies in **high** density regions

### MARKED DENSITY FIELD



**EM** et al. 2020

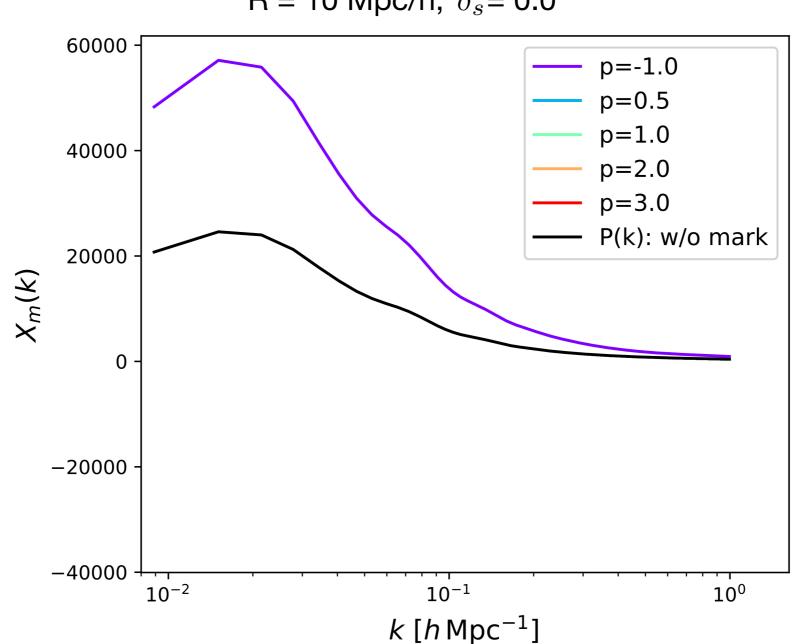
#### Marked-standard density cross-power spectrum



R = 10 Mpc/h,  $\delta_s$  = 0.0

**EM** et al. 2020

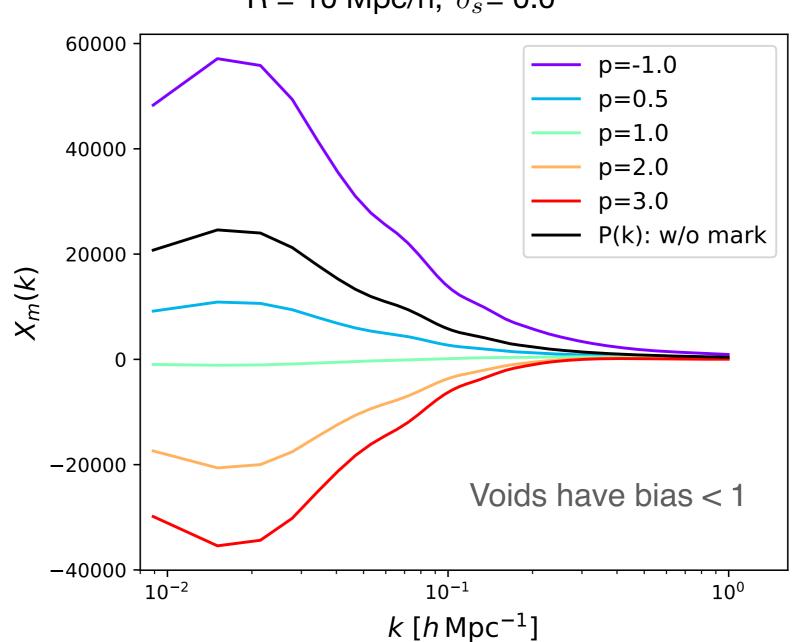




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**EM** et al. 2020





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## INFORMATION CONTENT in MARKED POWER SPECTRA of the MATTER FIELD

### FISHER ANALYSIS

Cosmological parameters:

$$\vec{\theta} = \{\Omega_m, \Omega_b, h, n_s, \sigma_8, M_\nu\}$$

Data vector (observables):

$$\vec{d} = \{P(k_1), P(k_2), ..., P(k_n)\}$$

Error on each parameter:

$$\sigma(\theta_{\alpha}) \leq \sqrt{(F^{-1})_{\alpha\alpha}}$$

Fisher matrix:

$$F_{\alpha,\beta} = \frac{\partial \vec{d}}{\partial \theta_{\alpha}} C^{-1} \frac{\partial \vec{d}}{\partial \theta_{\beta}}$$

### QUIJOTE SIMULATIONS

Villaescusa-Navarro, Hanh, EM et al 2019

- https://github.com/franciscovillaescusa/Quijote-simulations
- Set of 43,100 full N-body simulations
- 1 Gpc/h box size, 512<sup>3</sup> CDM particles (512<sup>3</sup> neutrinos)
- More than 7000 models with different  $\Omega_m, \Omega_b, h, \sigma_8, n_s, M_\nu, \omega$
- 1 Pb of publicly available data



### QUIJOTE SIMULATIONS

#### Villaescusa-Navarro, Hanh, EM et al 2019

#### Boxes to compute the covariances

Name	$\Omega_m$	$\Omega_b$	h	$n_s$	$\sigma_8$	$M_{\nu}$ [eV]	realizations	ICs
Fiducial	0.3175	0.049	0.6711	0.9624	0.834	0	15 ,000	2LPT

Table 1. Description of the N-body simulations used in the Fisher analysis.

### QUIJOTE SIMULATIONS

#### Villaescusa-Navarro, Hanh, EM et al 2019

#### Boxes to compute the numerical derivatives

Name	$\Omega_m$	$\Omega_b$	h	$n_s$	$\sigma_8$	$M_{\nu}$	realizations	ICs
						[eV]		
Fiducial ZA	0.3175	0.049	0.6711	0.9624	0.834	0	500	Zel'dovich
$\Omega_m^+$	0.3275	0.049	0.6711	0.9624	0.834	0	500	2LPT
$\Omega_m^-$	0.3075	0.049	0.6711	0.9624	0.834	0	500	2LPT
$\Omega_p^{++}$	0.3175	0.051	0.6711	0.9624	0.834	0	500	2LPT
$\Omega_p^{}$	0.3175	0.047	0.6711	0.9624	0.834	0	500	2LPT
$h^+$	0.3175	0.049	0.6911	0.9624	0.834	0	500	2LPT
$h^-$	0.3175	0.049	0.6511	0.9624	0.834	0	500	2LPT
$n_s^+$	0.3175	0.049	0.6711	0.9824	0.834	0	500	2LPT
$n_s^-$	0.3175	0.049	0.6711	0.9424	0.834	0	500	2LPT
$\sigma_8^+$	0.3175	0.049	0.6711	0.9624	0.849	0	500	2LPT
$\sigma_8^-$	0.3175	0.049	0.6711	0.9624	0.819	0	500	2LPT
$M_{\nu}^+$	0.3175	0.049	0.6711	0.9624	0.834	0.1	500	Zel'dovich
$M_{\nu}^{++}$	0.3175	0.049	0.6711	0.9624	0.834	0.2	500	Zel'dovich
$M_{\nu}^{+++}$	0.3175	0.049	0.6711	0.9624	0.834	0.4	500	Zel'dovich

Table 1. Description of the N-body simulations used in the Fisher analysis.

**EM** et al. 2020

#### The Mark

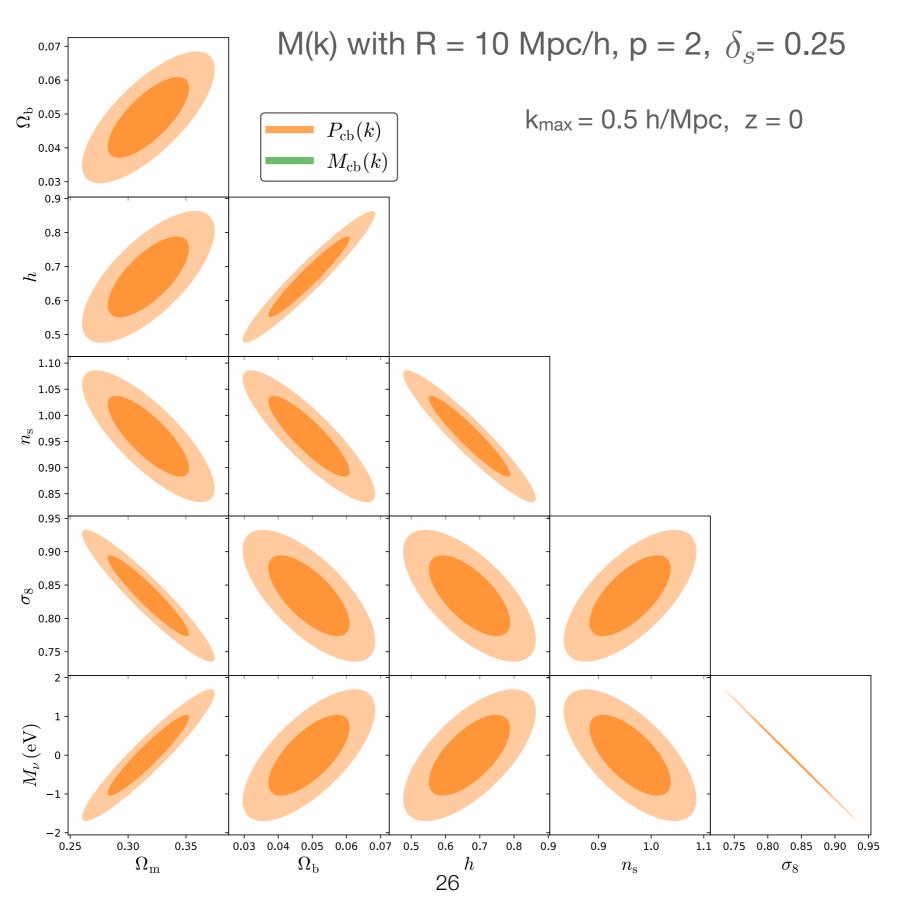
$$m(\vec{x}, \phi = R, p, \delta_s) = \left[\frac{1+\delta_s}{1+\delta_s+\delta_R(\vec{x})}\right]^p$$

#### Considered values for the mark parameters

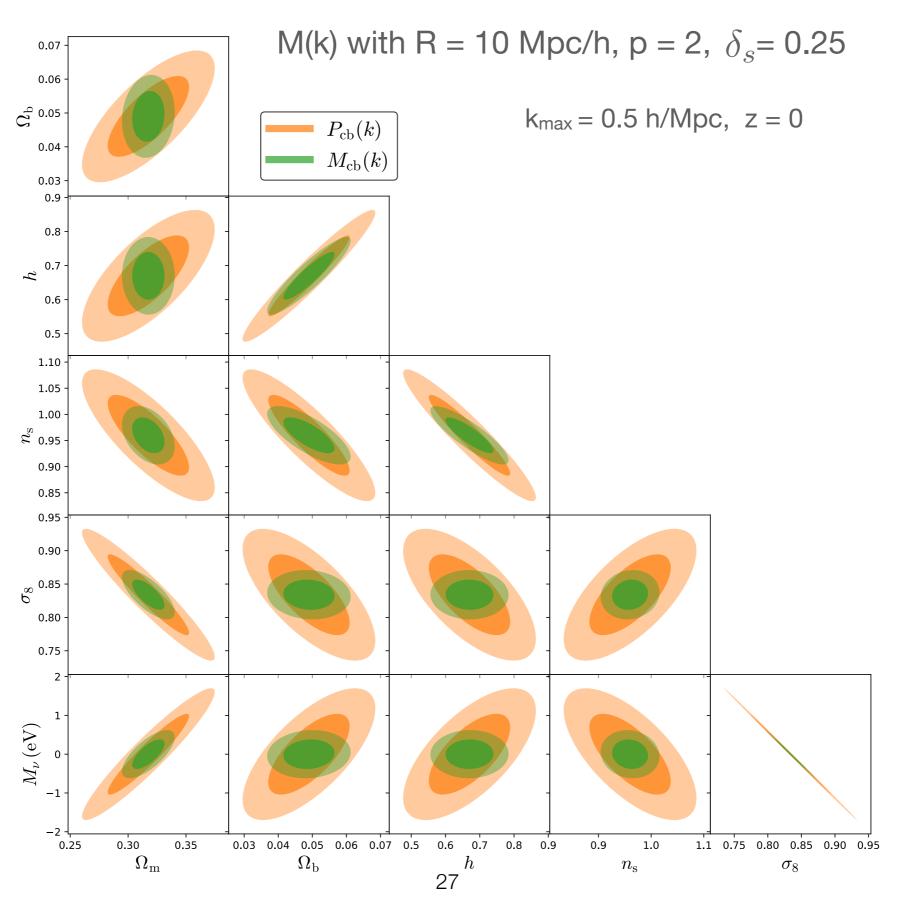
R = [5, 10, 15, 20, 30] Mpc/h  
p = [-1, 0.5, 1, 2, 3]  
$$\delta_s$$
 = [0, 0.25, 0.5, 0.75, 1]

#### <u>**125**</u> marked power spectra compute on the matter fields</u> cb (cdm) and m (cdm+neutrinos)

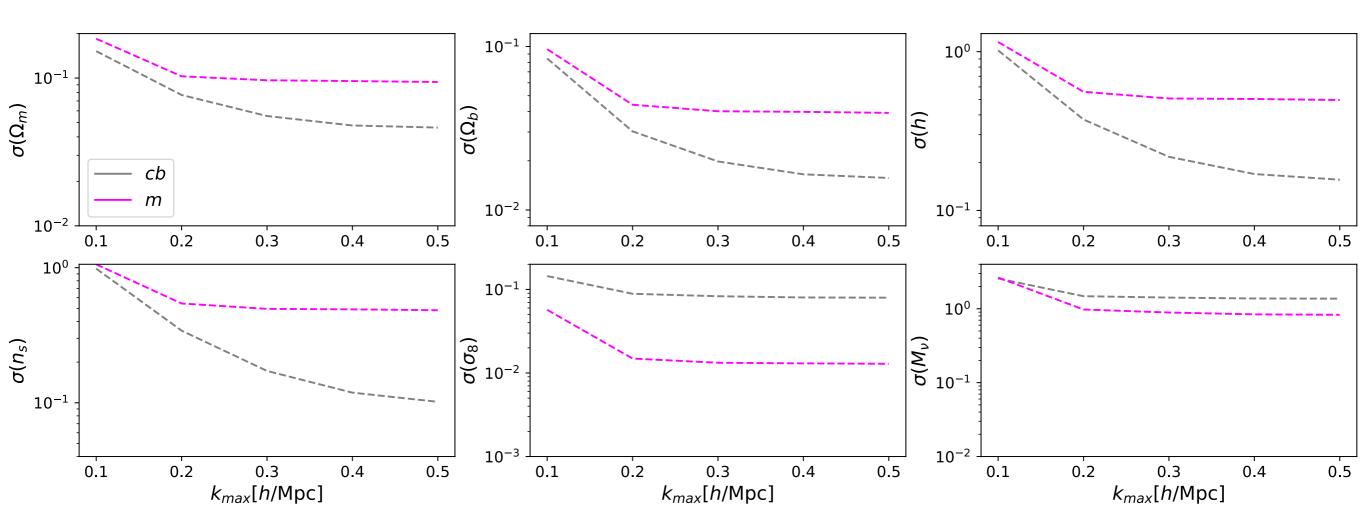
### Forecast for statistics of the cold dark matter



### Forecast for statistics of the cold dark matter



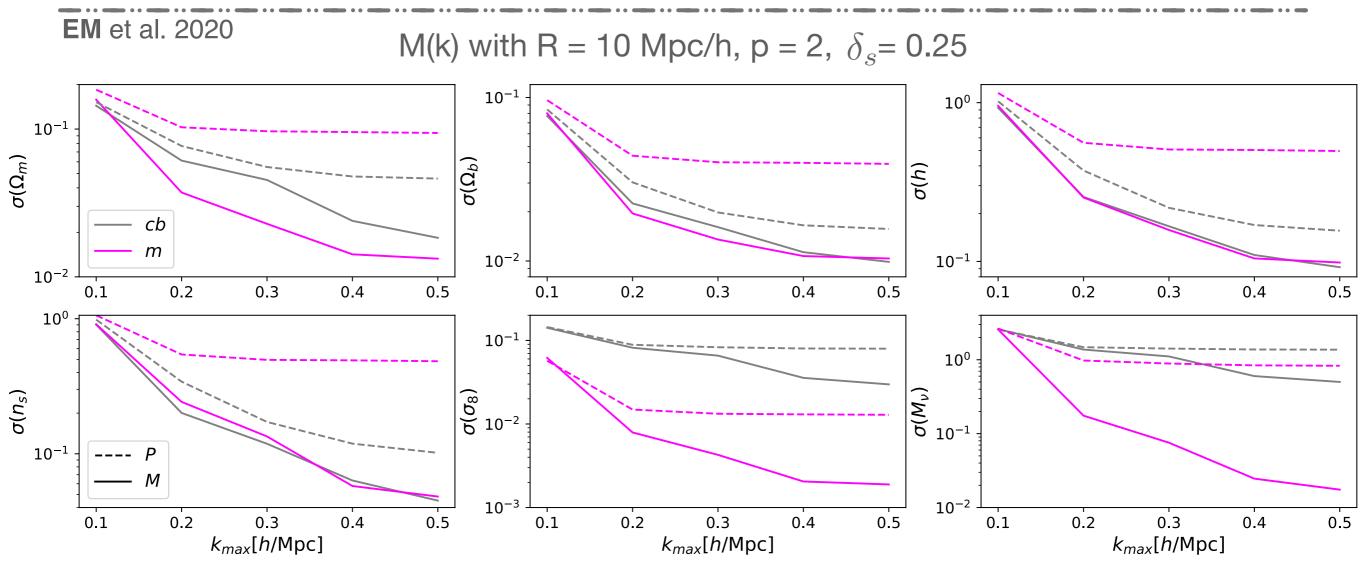
### Marginalized errors



Marginalized errors for  $k_{max} = 0.5 h/Mpc$ 

Parameter	$P_{cb}$	$M_{cb}$	$P_{cb}$ / $M_{cb}$	$P_m$	$M_m$	$P_m$ / $M_m$
$\Omega_m$	0.046			0.094		
$\Omega_b$	0.016			0.039		
h	0.16			0.50		
$n_s$	0.10			0.48		
$\sigma_8$	0.080			0.013		
$M_{ u}$	1.4			0.83		

### Marginalized errors



Marginalized errors for  $k_{max} = 0.5 h/Mpc$ 

Parameter	$P_{cb}$	$M_{cb}$	$P_{cb}$ / $M_{cb}$	$P_m$	$M_m$	$P_m$ / $M_m$
$\Omega_m$	0.046	0.018	2.5	0.094	0.013	7.2
$\Omega_b$	0.016	0.0099	1.6	0.039	0.010	3.9
h	0.16	0.092	1.7	0.50	0.098	5.1
$n_s$	0.10	0.045	2.2	0.48	0.048	10
$\sigma_8$	0.080	0.030	2.7	0.013	0.0019	6.8
$M_{ u}$	1.4	0.50	2.8	0.83	0.017	48

## INFORMATION CONTENT in MARKED POWER SPECTRA of the <u>GALAXY FIELD</u>

### Molino galaxy catalogs

Hahn et al. 2021

Built upon the Quijote simulations using Halo Occupation Distribution (HOD) framework from Zheng et al. (2007):

Mean central galaxy occupation

Mean satellite galaxy occupation

$$\langle N_c \rangle = \frac{1}{2} \left[ 1 - \operatorname{erf} \left( \frac{\log M_h - \log M_{\min}}{\sigma_{\log M}} \right) \right] \qquad \langle N_s \rangle = \langle N_c \rangle \left( \frac{M_h - M_0}{M_1} \right)$$

5 additional parameters to describe the BIAS scheme of GALAXIES

This prescription allow us to compute the **redshift-space** multiples (monopole and quadrupole) of the marked power spectrum of the galaxy field

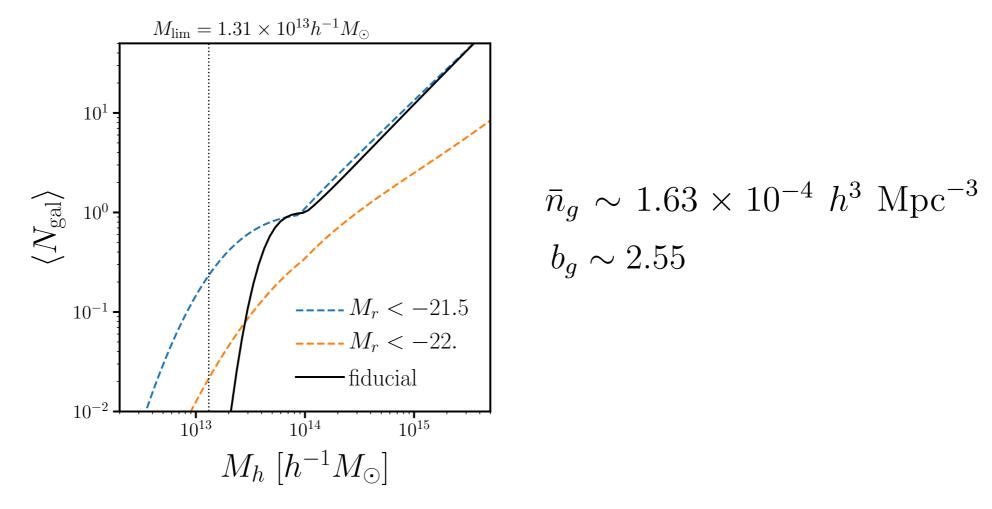
### Halo Occupation Distribution (HOD)

Hahn et al. 2021

#### Molino mock catalog

Fiducial:



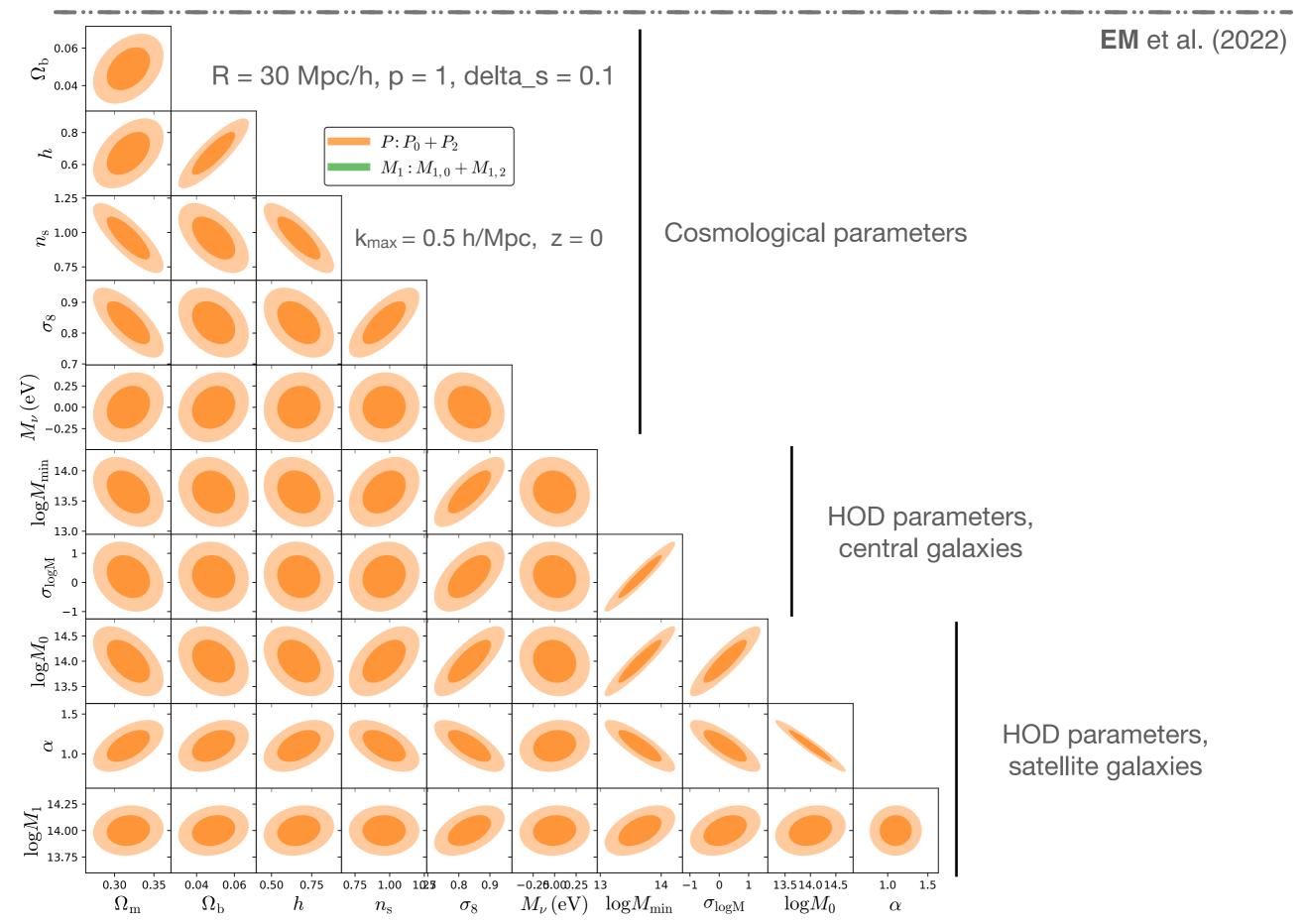


Variation to

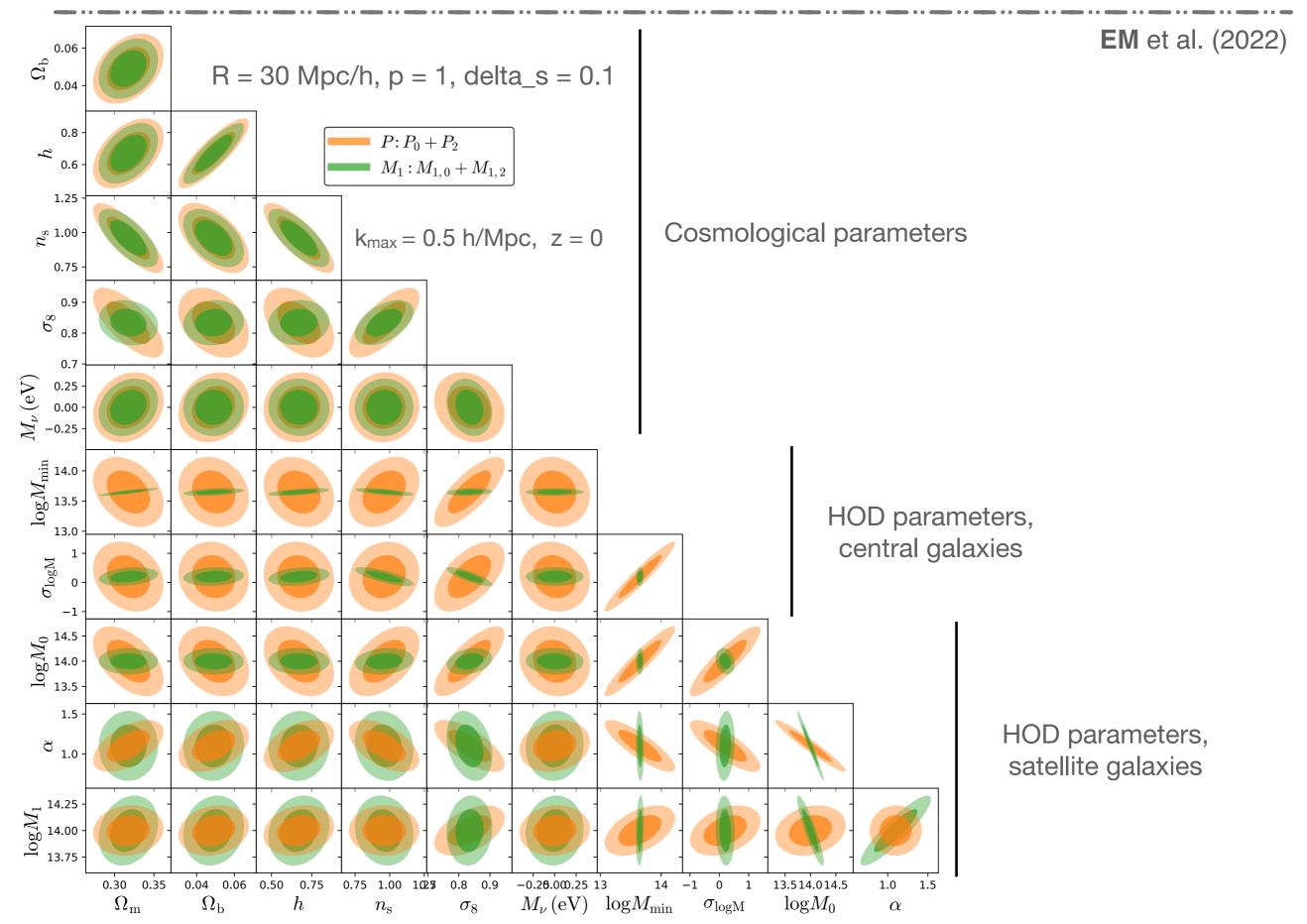
compute

derivatives:  $\{\Delta \log M_{\min}, \Delta \sigma_{\log M}, \Delta \log M_0, \Delta \alpha, \Delta \log M_1\} = \{0.05, 0.2, 0.2, 0.2, 0.2\}$ 

### Forecast for the redshift-space Galaxy field



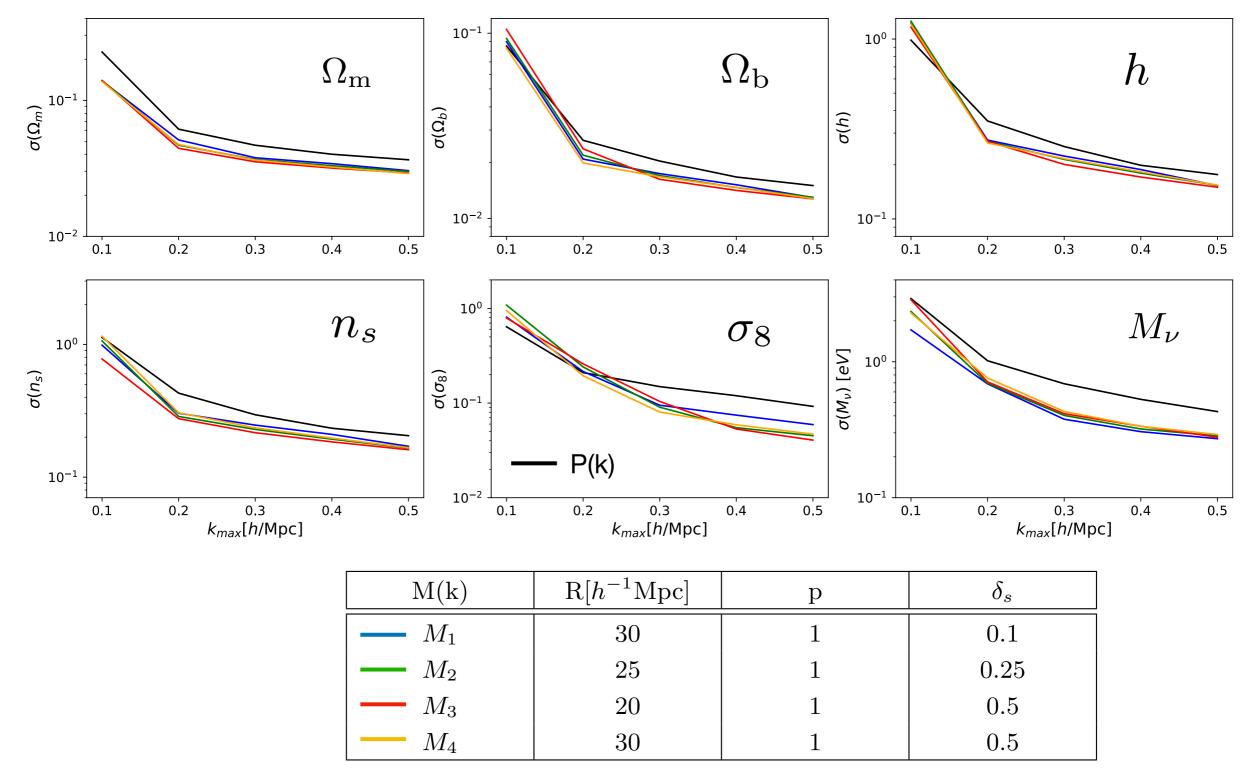
### Forecast for the redshift-space Galaxy field



### Marginalized errors - galaxy field

#### EM et al. (2022)

Constraints using monopole and quadrupole of different statistics up to k = 0.5 h/Mpc



**Table 2.** Values for the mark parameters  $(R, p, \delta_s)$  of selected marked power spectra  $M_1, M_2, M_3$ , and  $M_4$ .

### Marginalized errors - galaxy field

**EM** et al. (2022)

The combination of **4 marked power spectra and the standard power spectrum** can largely improve the cosmological constraints coming from the power spectrum alone.

Θ	$P/(P + \Sigma_i M_i)$
	0+2
$\Omega_m$	2.4
$\Omega_b$	2.5
h	2.6
$n_s$	3.3
$\sigma_8$	6.1
$M_{\nu}$	3.0

Future surveys (DESI, EUCLID, Roman) will observe a **LARGER number density** of galaxies that will allow them to **trace the inner part of voids** better than the Molino catalogs. Thus, they will be able to exploit the potential the M(k) even further!

### WHY DO MARKED POWE SPECTRA CONTAIN INFORMATION BEYOND THE STANDARD POWER SPECTRUM?

Philcox, EM et al. 2020

The marked density field:

$$\delta_M(\mathbf{x}) \equiv \frac{1}{\bar{m}} m(\mathbf{x}) \left[ 1 + \delta(\mathbf{x}) \right] - 1$$

with mark:

$$m(\mathbf{x}) = \left(\frac{1+\delta_s}{1+\delta_s+\delta_R(\mathbf{x})}\right)^p \equiv \left(1+\frac{\delta_R(\mathbf{x})}{1+\delta_s}\right)^{-p}$$

Philcox, EM et al. 2020

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Taylor expanding the mark:

$$\delta_M(\mathbf{x}) = \frac{1}{\bar{m}} \left[ 1 + \delta(\mathbf{x}) \right] \left[ 1 - C_1 \delta_R(\mathbf{x}) + C_2 \delta_R^2(\mathbf{x}) - C_3 \delta_R^3(\mathbf{x}) \right] - 1 + \mathcal{O}\left(\delta^4\right).$$

Philcox, EM et al. 2020

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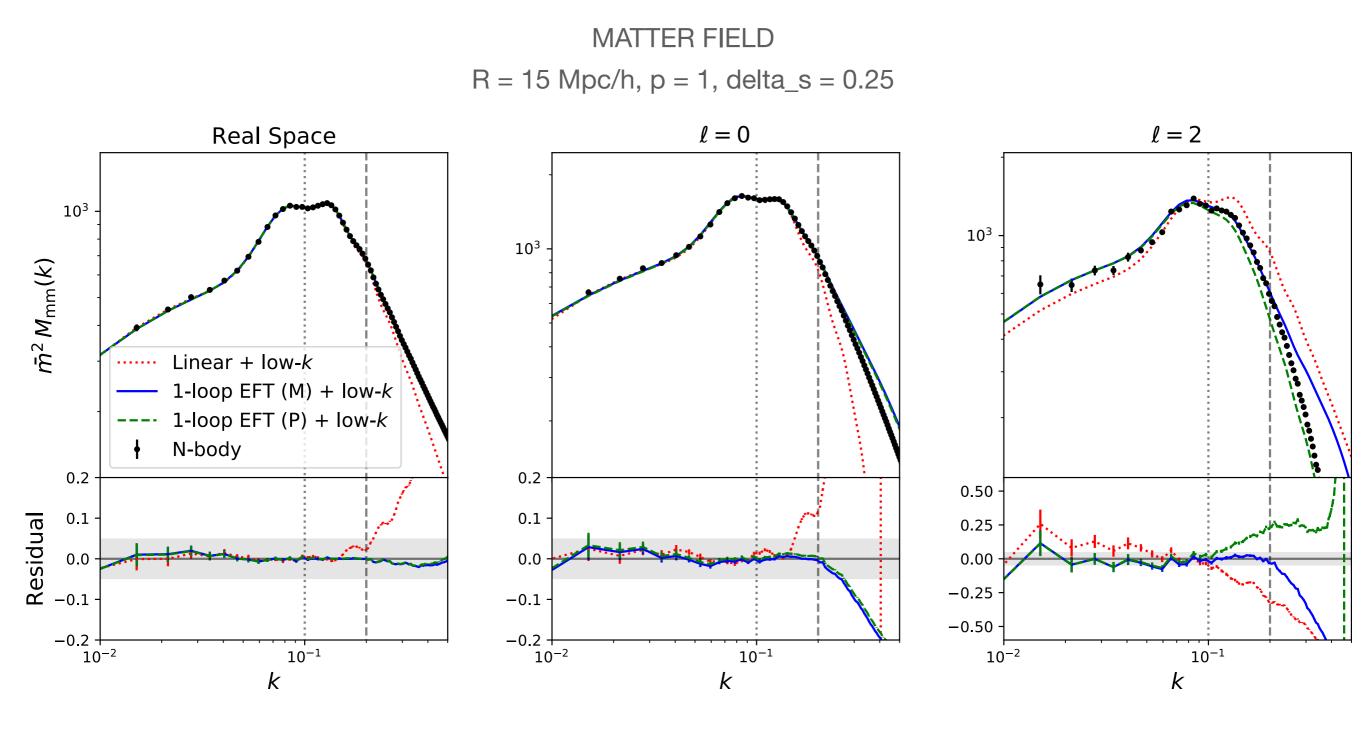
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#### 2-point function of the marked field contains higher order statistics of the original field

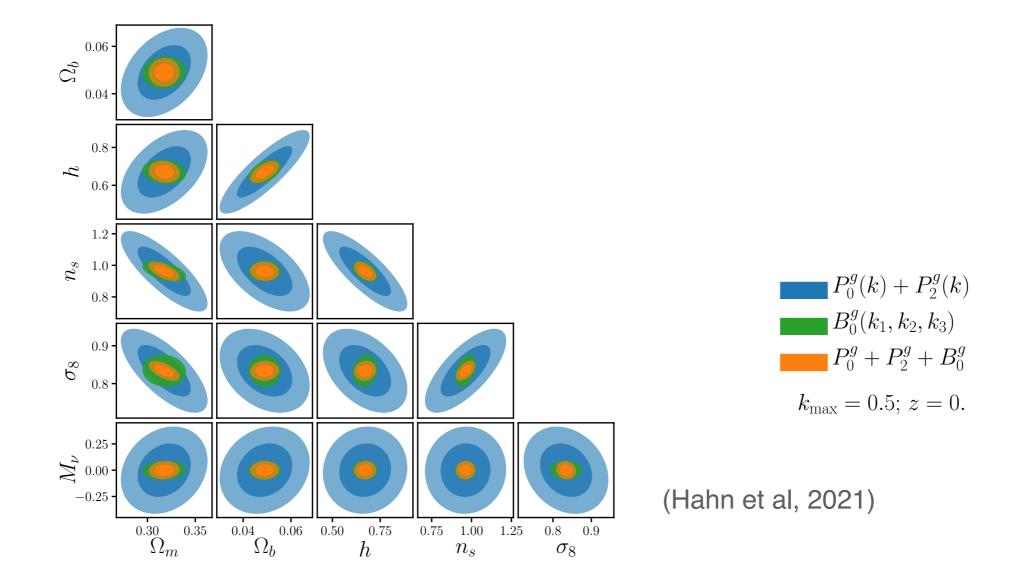
### EFT for MARKED POWER SPECTRA

Philcox, EM et al. 2020, Philcox, Aviles, EM et al. 2021



### THE INFORMATION CONTENT

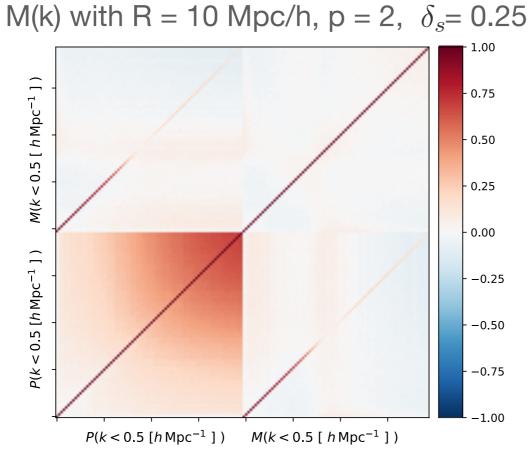
• Marked power spectra contain higher order statistics of the density field



### THE INFORMATION CONTENT

• Marked power spectra contain higher order statistics of the density field

 The covariance matrix of some marked power spectra M(k) that up-weight low-density regions is almost diagonal



Other nonlinear transformations, such as the log-transformation, have shown to make the field more Gaussian (Neyrinck et al, 2009, 2010, 2011)

### THE INFORMATION CONTENT

• Marked power spectra contain higher order statistics of the density field

• The **covariance** matrix of some marked power spectra M(k) that up-weight low-density regions is almost **diagonal** 

 Marked power spectra that up-weight low-density regions incorporate information from voids

### CONCLUSIONS

- Results from Fisher analyses: marked power spectra that up-weight low density regions improve parameter constraints beyond the standard power spectrum.
- 6x tighter constraints for sigma8 and 2-3x for the other cosmological parameters when considering combinations of marked and standard power spectra of the galaxy field.
- Upcoming surveys (DESI, EUCLID, Roman) will probe larger volumes and higher galaxy number density, that will allow them to better explore low-density regions and improve the performance of marked power spectra.
- Next step: cosmological analysis with marked power spectra in available surveys. We are building a simulation-based inference framework that will allow us to forward modeling survey systematics and geometry.