

Instabilities appearing in dark energy models

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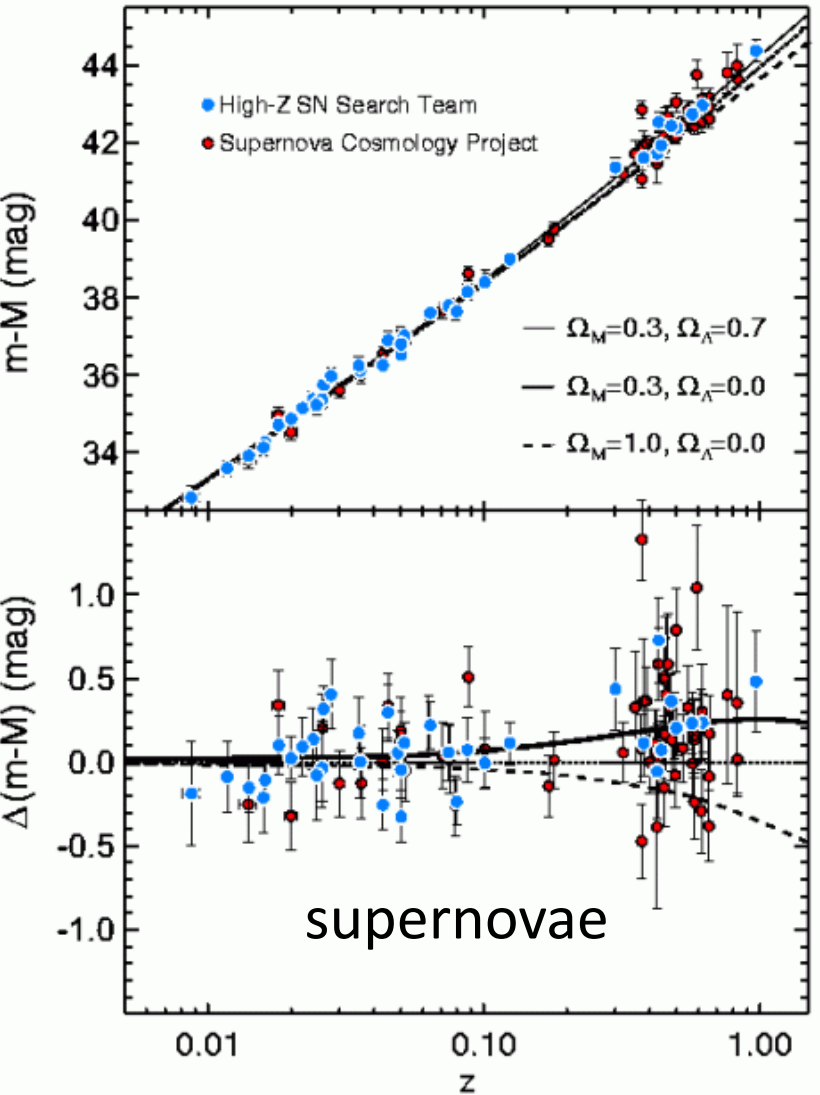
arXiv:2205.01055, arXiv:2204.13098, arXiv: 2107.14215

In collaboration with Julian Adamek, Jean-Pierre Eckmann, Martin Kunz, Pan Shi, Peter Wittwer, Hatem Zaag.

UiO : Universitetet i Oslo

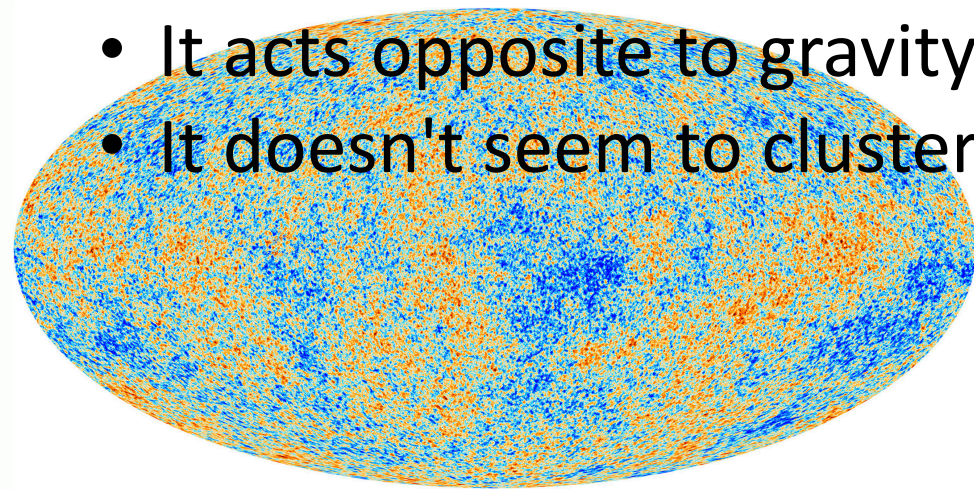


Accelerating expansion of the Universe:

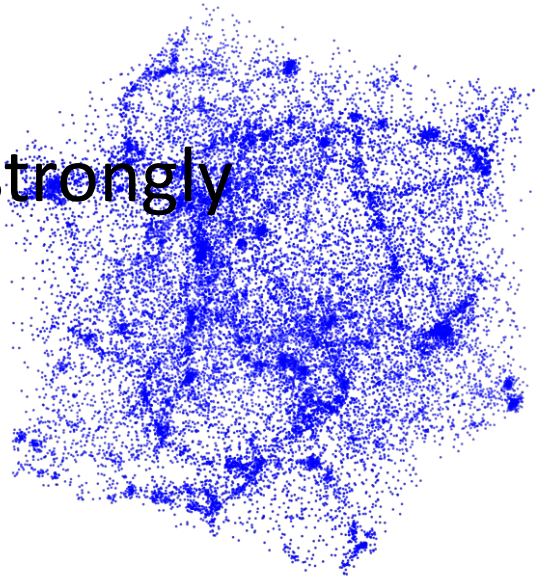


$$\Omega_{DE} \sim 0.7$$

- It acts opposite to gravity
- It doesn't seem to cluster strongly



The CMB


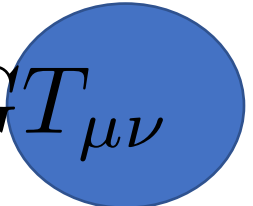


Large scale structure

The reason behind the cosmic acceleration:

➤ Possibilities:

- Cosmological constant /Dark energy /beyond general relativity

Modified gravity  $G_{\mu\nu} = 8\pi G$  $T_{\mu\nu}$ Dark energy

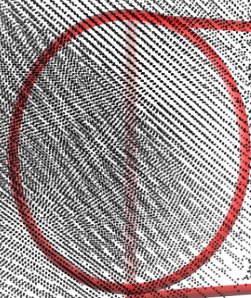
➤ Effective field theory of dark energy:

- Convenient way to describe DE and MG: a general action for class of theories
- Connecting to observations minimally

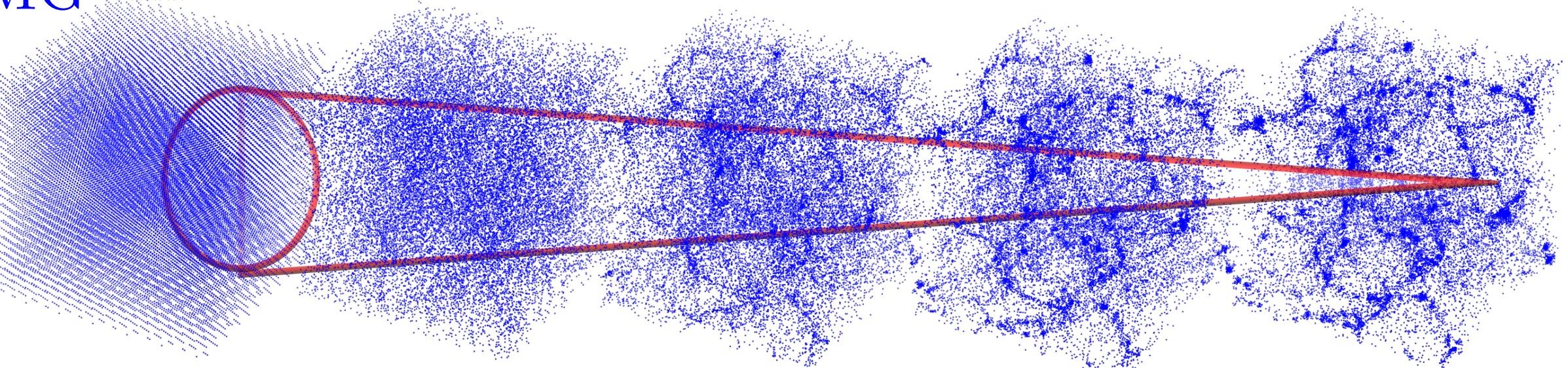
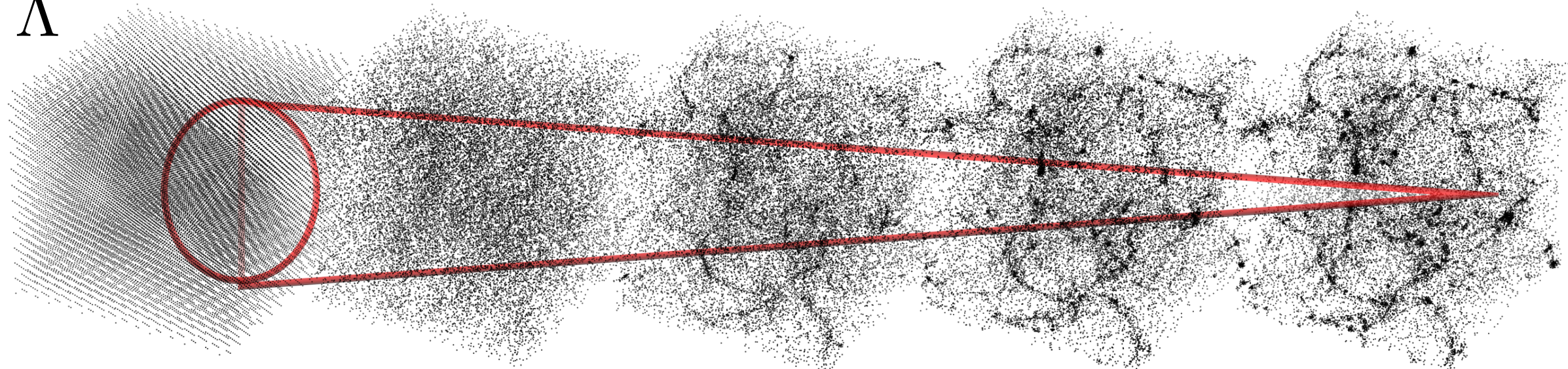
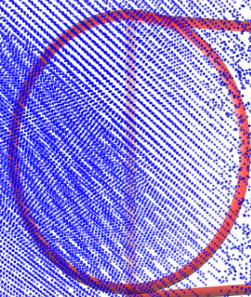
➤ DE/MG theories are studied well at the level of the background and linear

Are these theories really viable?

Λ z_{ini}




MG z_{ini}



Importance of non-linear modelling:

- Non-linear evolution of the k-essence leads to an instability

$$S_{\text{DE}} = \int d^4x \sqrt{-g} P(X, \phi)$$


$$w, c_s^2$$

- The instability:

- In the cosmological context (using k-evolution N-body code)

(Hassani, Adamek, Kunz, Shi, Wittwer; arXiv: 2107.14215 & arXiv: 2204.13098)

- In 1+1 dimensions

(Eckmann, Hassani, Zaag; arXiv: 2205.01055)

Non-linear scalar field equation: $\pi = \frac{\delta\phi}{\phi}$

$$\begin{aligned} & \partial_\tau^2 \pi + \mathcal{H}(1 - 3w)\partial_\tau \pi + \left(\partial_\tau \mathcal{H} - 3w\mathcal{H}^2 + 3c_s^2(\mathcal{H}^2 - \partial_\tau \mathcal{H}) \right) \pi \\ & - \partial_\tau \Psi + 3\mathcal{H}(w - c_s^2)\Psi - 3c_s^2 \partial_\tau \Phi - c_s^2 \nabla^2 \pi \\ & = \mathcal{N}(\pi, \partial_\tau \pi, \vec{\nabla} \pi, \vec{\nabla} \partial_\tau \pi, \nabla^2 \pi) \end{aligned}$$

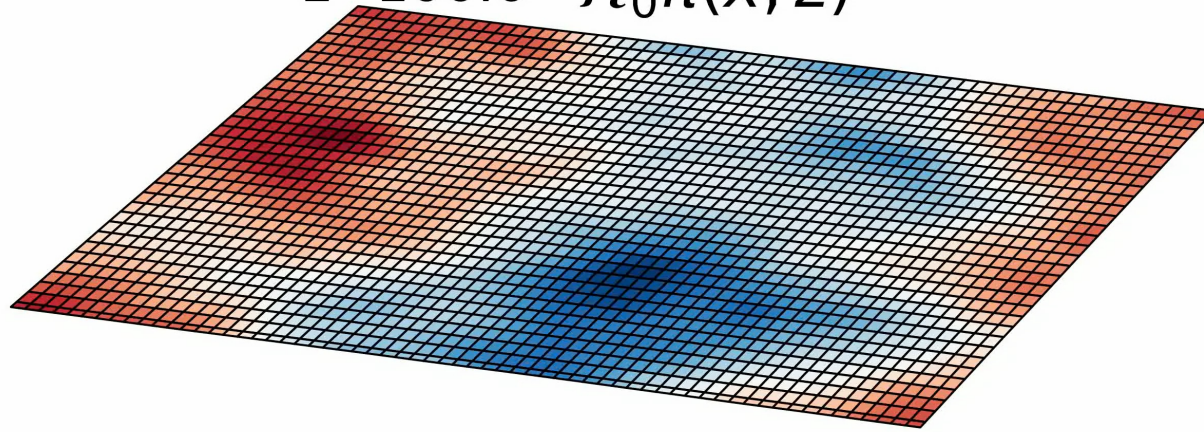
$$\begin{aligned} \mathcal{N}(\pi, \partial_\tau \pi, \vec{\nabla} \pi, \vec{\nabla} \partial_\tau \pi, \nabla^2 \pi) &= -\frac{\mathcal{H}}{2} (5c_s^2 + 3w - 2) \boxed{(\vec{\nabla} \pi)^2} \\ &+ 2(1 - c_s^2) \vec{\nabla} \pi \cdot \vec{\nabla} \partial_\tau \pi - \left[(c_s^2 - 1)(\partial_\tau \pi + \mathcal{H}\pi - \Psi) \right. \\ &+ c_s^2(\Phi - \Psi) + 3\mathcal{H}c_s^2(1 + w)\pi \left. \right] \nabla^2 \pi + (2c_s^2 - 1) \vec{\nabla} \Psi \cdot \vec{\nabla} \pi \\ &- c_s^2 \vec{\nabla} \Phi \cdot \vec{\nabla} \pi + \frac{3(c_s^2 - 1)}{2} \partial_i \left(\partial_i \pi (\vec{\nabla} \pi)^2 \right). \end{aligned}$$



$$c_s^2 < 10^{-4.7}$$

$$\mathcal{H}_0 \pi$$

$z=100.0 \quad \mathcal{H}_0 \pi(\vec{X}, z)$



No instability for large c_s^2

Non-linear scalar field equation:

$$\begin{aligned} & \partial_\tau^2 \pi + \mathcal{H}(1 - 3w)\partial_\tau \pi + \left(\partial_\tau \mathcal{H} - 3w\mathcal{H}^2 + 3c_s^2(\mathcal{H}^2 - \partial_\tau \mathcal{H}) \right) \pi \\ & - \partial_\tau \Psi + 3\mathcal{H}(w - c_s^2)\Psi - 3c_s^2 \partial_\tau \Phi - c_s^2 \nabla^2 \pi \\ & = \mathcal{N}(\pi, \partial_\tau \pi, \vec{\nabla} \pi, \vec{\nabla} \partial_\tau \pi, \nabla^2 \pi) \end{aligned}$$

3+1 D



Spherical symmetry

$$\frac{d^2 \pi}{d\tau^2} = +c_s^2 \frac{\partial^2 \pi}{\partial r^2} - \alpha \left(\frac{d\pi}{dr} \right)^2$$

1+1 D

$$c_s^2 \rightarrow 0 \quad \frac{d^2 \pi}{d\tau^2} = -\alpha \left(\frac{d\pi}{dr} \right)^2$$

This PDE is unstable. The curvature of the scalar field around minima blows up in finite time.

$$\frac{d^2 \pi}{d\tau^2} = -\alpha \left(\frac{d\pi}{d\tau} \right)^2$$

One special solution

$$\pi_s(\tau, r) = \frac{\kappa(\tau)}{2} r^2$$

$$\partial_\tau^2 \kappa(\tau) = \alpha [\kappa(\tau)]^2$$

We can think of this ODE as
Newton's second law

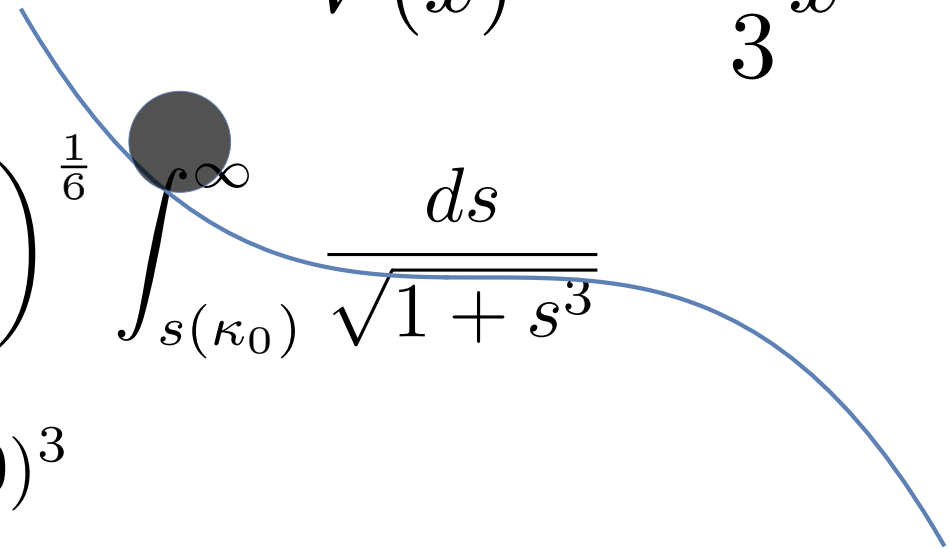
$$F(x) = x^2$$

$$V(x) = -\frac{4}{3} x^3$$

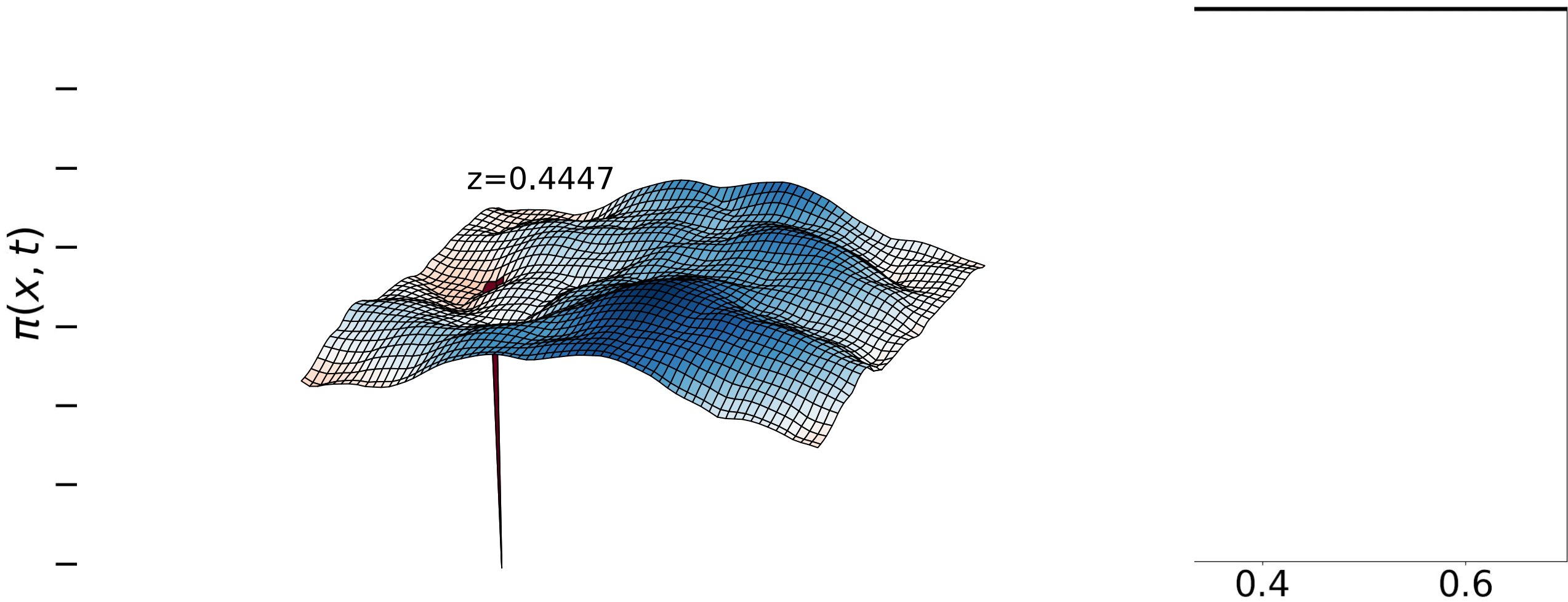
$$\kappa(\tau) = \frac{3}{2(\tau - \tau_b)^2}$$

$$\tau_b = \left(\frac{3}{8\alpha} \right)^{\frac{1}{3}} \left(\frac{1}{C} \right)^{\frac{1}{6}} \int_{s(\kappa_0)}^{\infty} \frac{ds}{\sqrt{1+s^3}}$$

$$C = \kappa'(0)^2 - \frac{8}{3} \kappa(0)^3$$



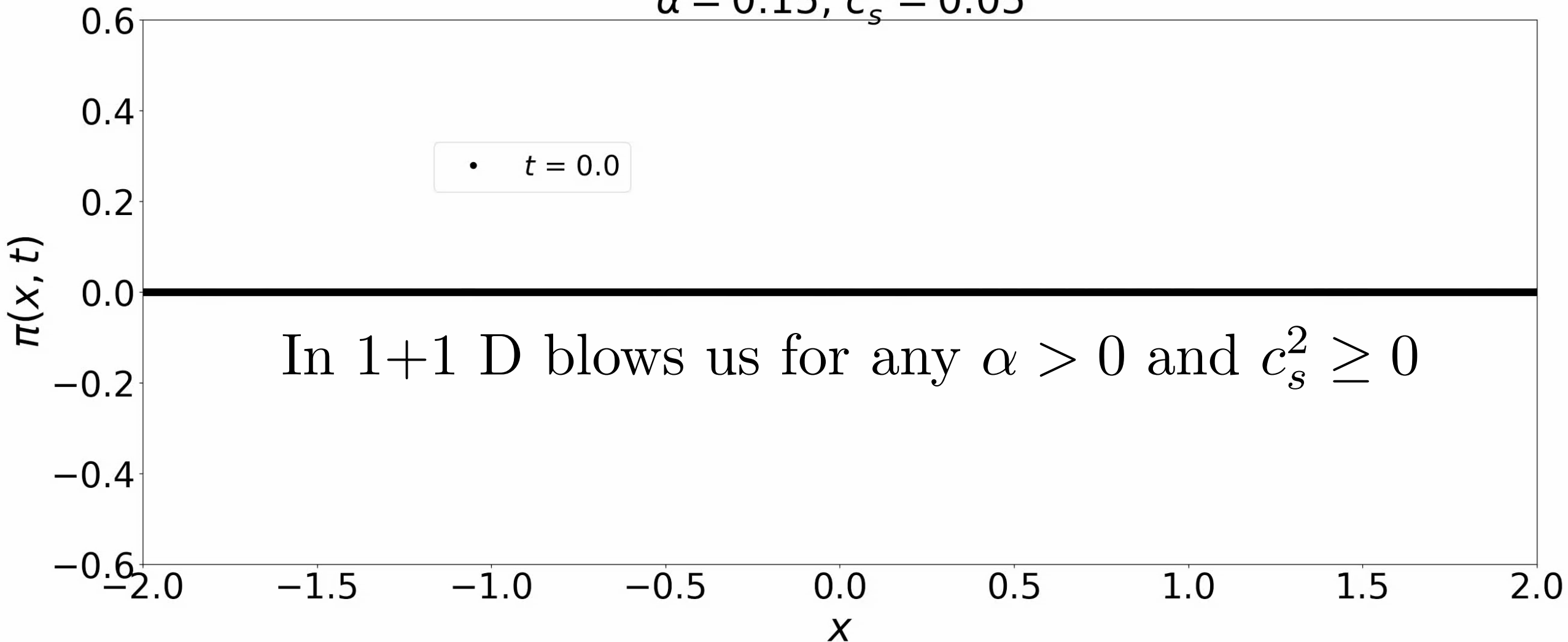
V-type blowup $\pi(0, x) = 0, \partial_\tau \pi(0, x) = \Phi(x)$



(Hassani, Adamek, Kunz, Shi, Wittwer; arXiv: 2107.14215 & arXiv: 2204.13098)

M-type blowup

$$\alpha = 0.15, c_s^2 = 0.05$$



$$\frac{d^2 u}{d\tau^2} = +c_s^2 \frac{\partial^2 u}{\partial x^2} - \alpha \left(\frac{du}{dx} \right)^2$$

(Eckmann, Hassani, Zaag; arXiv: 2205.01055)

Take home messages:

- Non-linearities can be **VERY** important
- Cosmological N-body simulations bound : $c_s^2 \geq 10^{-4.7}$
- From mathematical point of view any c_s^2 leads to an instability!
- Can k-essence be a cosmic acceleration candidate?

Thank you 😊