## Dark Energy: Can we Vainshtein screen a fifth fundamental force?

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# The problem

Vainstein screening and UV completion

- Theories attempting to solve the cosmological constant problem come with extra scalar fields
- Scalar extensions come with extra mediated
   5<sup>th</sup> forces
- Do we see extra forces experimentally?



E.G.Adelberger J.H.Gundlach B.R.Heckel S.Hoedl, S.Schlamminger, Progr. in Part. and Nucl. Phys., 62, 1, 2009, 102-134; C. Will, Living Reviews in Relativity 2006.

# Canonical scalar fields with small mass disfavoured

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - \frac{1}{2} m^2 \phi^2 - \frac{\phi}{M} T^{\mu}{}_{\mu}$$

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- If scalar field has non-trivial self-interactions,
   5<sup>th</sup> force may be *screened* (= suppressed)
   where it is known to be small...
- ... while still being able to mediate a longrange force where it can be relevant at current cosmological scales

$$\mathcal{L} = -\frac{1}{2} Z^{\mu\nu}(\phi, \partial\phi, ...) \partial_{\mu} \partial_{\nu} \phi - V(\phi) + g(\phi) T^{\mu}{}_{\mu}$$

 $V(r) = -\alpha G \frac{M_s m}{r} e^{-\frac{r}{\lambda}}$  $V(r) = -\alpha(\phi_{bg})G\frac{f(M_s)m}{r}e^{-\frac{r}{\lambda(\phi_{bg})}}$ 

Standard gravity  $\Box \phi$ 

Cubic Galileons (with Vainshtein)

$$\Box \phi + (\Box \phi)^2 - (\partial_\mu \partial_\nu \phi)^2 = \frac{\rho}{M}$$

High derivative, nonlinear terms dominate

Standard terms dominate



M



$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$$



Typically, Vainshtein screening is taken to be an effective field theory...

...but if specific higher-order kinetic terms are large, then we are out of a perturbative regime...

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A. Padilla and I. D. Saltas, JCAP 06 (2018) 039; M. A. Luty, M. Porrati, R. Rattazzi, JHEP 0309 (2003) 029; A. Nicolis and R. Rattazzi, JHEP 0406 (2004) 059

 $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$ 



why are these higher-order kinetic term macroscopically important but others are not?

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 $\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2\Lambda^3}(\partial\phi)^2\Box\phi$ 



You can only really tell if you know the full UV sector

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### Vainshtein screening – more accurate picture



Hinterbichler, Rev. Mod. Phys. 84, 671-710 (2012)

### Vainshtein screening – more accurate picture

Λ<sub>3</sub> theories: classical nonlinearities kick in before quantum corrections; however, standard UV completion not possible



Hinterbichler, Rev. Mod. Phys. 84, 671-710 (2012)

### Vainshtein screening – more accurate picture

 $\Lambda_5$  theories: standard UV completion possible; however, quantum corrections kick in before classical nonlinearities



Hinterbichler, Rev. Mod. Phys. 84, 671-710 (2012)

## The goal

• Compare a low-energy theory with Vainshtein screening against its UV completion

• Does the screening survive the extension?

#### A UV-complete theory of massive Galileons

$$\mathcal{L} = -\frac{1}{2} (\partial \phi)^2 - \frac{1}{2} m^2 \phi^2 - \frac{1}{2} (\partial H)^2 - \frac{1}{2} M^2 H^2 - \frac{$$

Theory of two coupled massive scalar fields, one self-coupled Important higher-order kinetic terms

Does it screen?

C. de Rham, S. Melville, A. J. Tolley, and S.-Y. Zhou, JHEP 9 (2017) 72

A UV-complete theory of massive Galileons

$$\begin{cases} \Box \phi - m^2 \phi - \alpha \Box H = \frac{\rho}{M_{pl}} \\ \Box H - M^2 H - \alpha \Box \phi - \frac{\lambda}{3!} H^3 = 0 \end{cases}$$

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#### Low-energy theory

From this high-energy theory, it is possible to obtain the IR theory

$$\Box \pi - m^{2} \pi - \frac{\lambda \alpha^{4}}{3!} \frac{\Box (\Box \pi^{3})}{M^{8}} + \dots = \frac{\rho}{M_{pl}}$$

which relies on the series of operators

$$O_n \propto (-1)^{n+1} \lambda^n \alpha^{2n+1} \frac{\Box(\Box \pi^{2n+1})}{M^{6n+2}}$$

to converge.

# Does it look anything like the high-energy theory?



- Let's compute the scalar force for the highenergy and low-energy theories, to see which one, if any, displays Vainshtein screening
- Let's compute the O<sub>n</sub> operators to see if the series converge

C. Burrage, B. Coltman, A. Padilla, DS, T. Wilson, JCAP 02 (2021) 050 19

### The problem

• We need to solve complex equations characterised by large non-linearities, sharp transitions, complex boundary conditions

VS



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• We need to solve complex equations characterised by large non-linearities, sharp transitions, complex boundary conditions





# Method

 $\varphi$ enics: screening with the finite element method



- Finite-element code for the solution of the full equation of motion of models of screening
- Builds on FEniCS library

 Computes field profiles, fifth force, (horrible) high-order operators accurately

Step (1): discretise domain



Credits: FEniCS project

Functions can be approximated piecewise by polynomials:



Step (2): turn original (*strong*) equation into its weak form

• Example: Poisson equation

Strong problem: find u such that

$$-\nabla^2 u = f$$
$$u = u_D$$
$$-\frac{\partial u}{\partial n} = g$$

in  $\Omega$  (spatial domain) in  $\partial \Omega_D$  (Dirichlet boundary) in  $\partial \Omega_N$  (Neumann boundary), n outward normal direction to the boundary

 $J_{\partial\Omega}$ 

Weak formulation: find 
$$u$$
 such that  

$$-\int_{\Omega} \nabla^2 u \, v = \int_{\Omega} f v \, \forall v \in \hat{V} \qquad \forall v \in \hat{V},$$

$$\hat{V} \text{ space of 'test' functions}$$

Integrating by parts and applying Neumann boundary conditions...

$$\int_{\Omega} \nabla u \nabla v = -\int_{\partial \Omega_N} gv \, ds + \int_{\Omega} fv$$

$$a(u, v) := \int_{\Omega} \nabla u \nabla v \qquad \qquad \text{bilinear form}$$

$$L(v) := -\int gv \, ds + \int fv \qquad \qquad \text{linear form}$$

Step (3): turn complicated equation into algebraic system

Because a and L are bilinear and linear, respectively, they only need to be defined on basis vectors:

find u such that:

 $a(u,v) = L(v) \quad \forall v \in \hat{V}$ 

Expanding in the discretised basis:

$$a(u, e_i) = L(e_i) \quad \forall i = 1, \dots, N$$

$$\sum_{j=1}^{N} u_j a(e_j, \hat{e}_i) = L(\hat{e}_i) \quad \forall i = 1, \dots, N \quad \longrightarrow \quad AU = b$$

$$Coefficients of$$

Coefficients of u in discretised basis

Step (4): if equation is nonlinear, solve iteratively using e.g. Newton's method



 $F(u_k) + J(u_k)(u - u_k) = 0$ 

Starting from some initial guess  $u_k$ , Then substituting  $u \rightarrow u_k$  until convergence Features at the source-vacuum transition



### Terms in the equations across all regimes



### High order operators





Vainstein screening and UV completion

### Vainshtein screening – is it viable?



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# Summary and conclusions

### Summary and conclusions

- Vainshtein screening is an effective mechanism to hide the fifth force in high-density regions; it is characterised by complicated equations that are impossible to solve analytically and difficult to solve numerically;
- We developed  $\varphi$ enics, finite-element code for the full numerical solution of equations of screening
- We compared a theory of massive Galileons against its UV completion: the IR theory screens, the UV theory doesn't!
- It is unclear Vainshtein screening can be invoked within the limits of validity of the theories that make use of it

# Thank you for your attention!