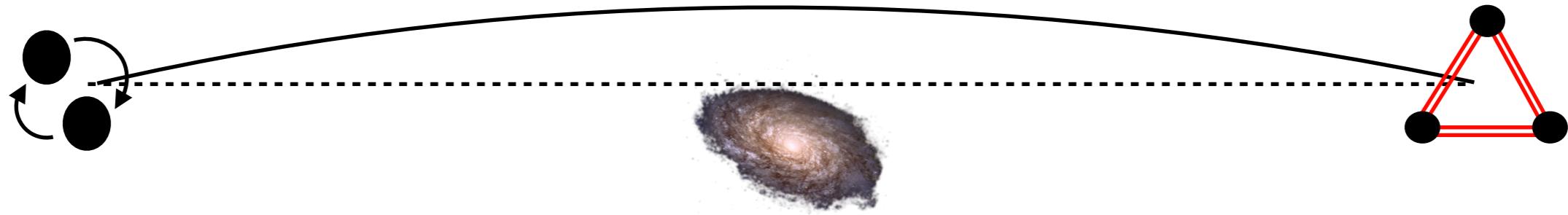


Polarization distortions of lensed gravitational waves

Charles Dalang

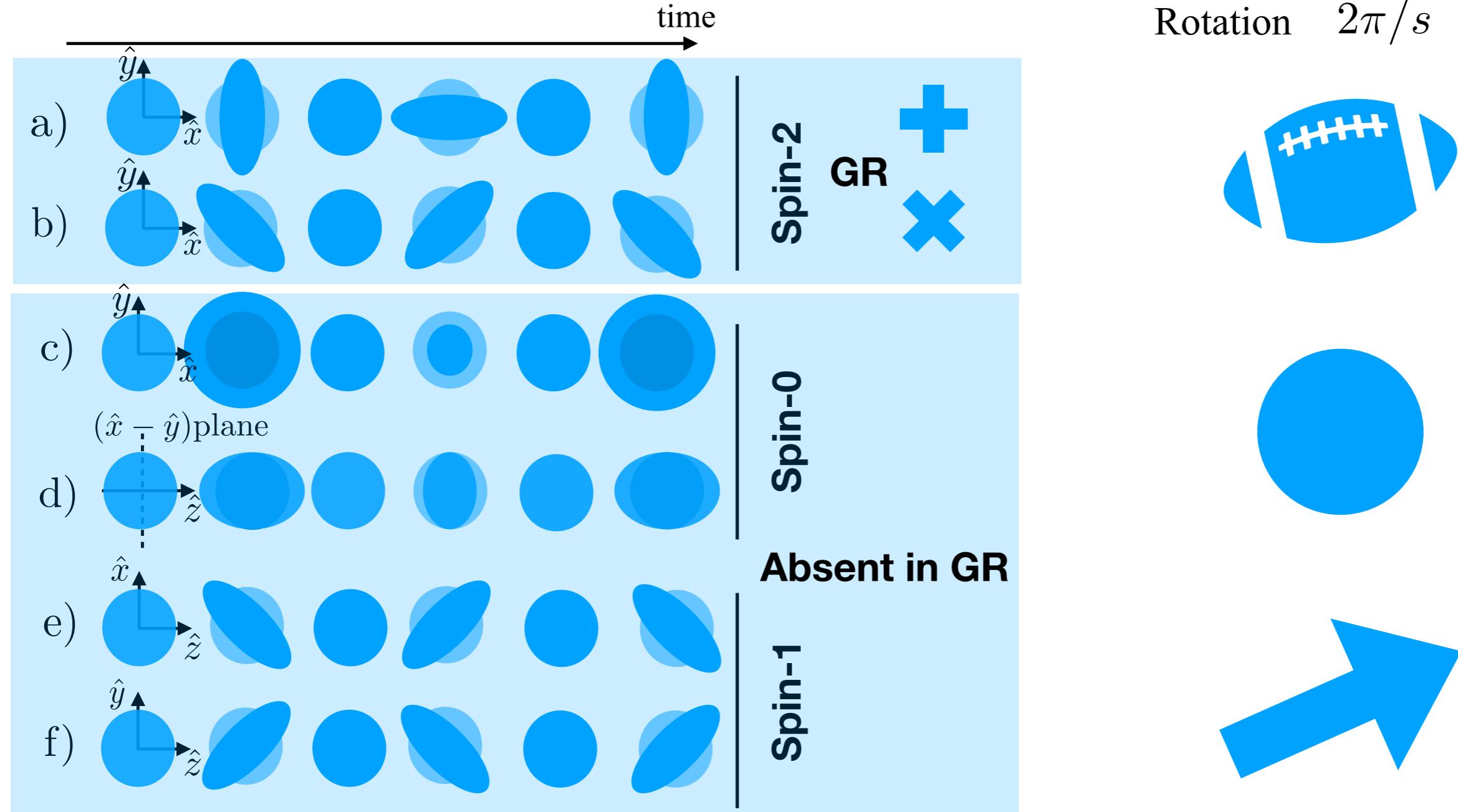
Département de Physique Théorique
Université de Genève
Switzerland



References

- C.D., Giulia Cusin and Macarena Lagos, [2104.10119]
Giulia Cusin and Macarena Lagos, [1910.13326] Phys. Rev. D.

GW extra polarization modes are a *smoking gun* signature of beyond GR physics.



GW propagation in curved spacetime

Misner, Thorne, Wheeler, 1973

Giulia Cusin, Macarena Lagos, 2020

C.D., Giulia Cusin, Macarena Lagos, 2021

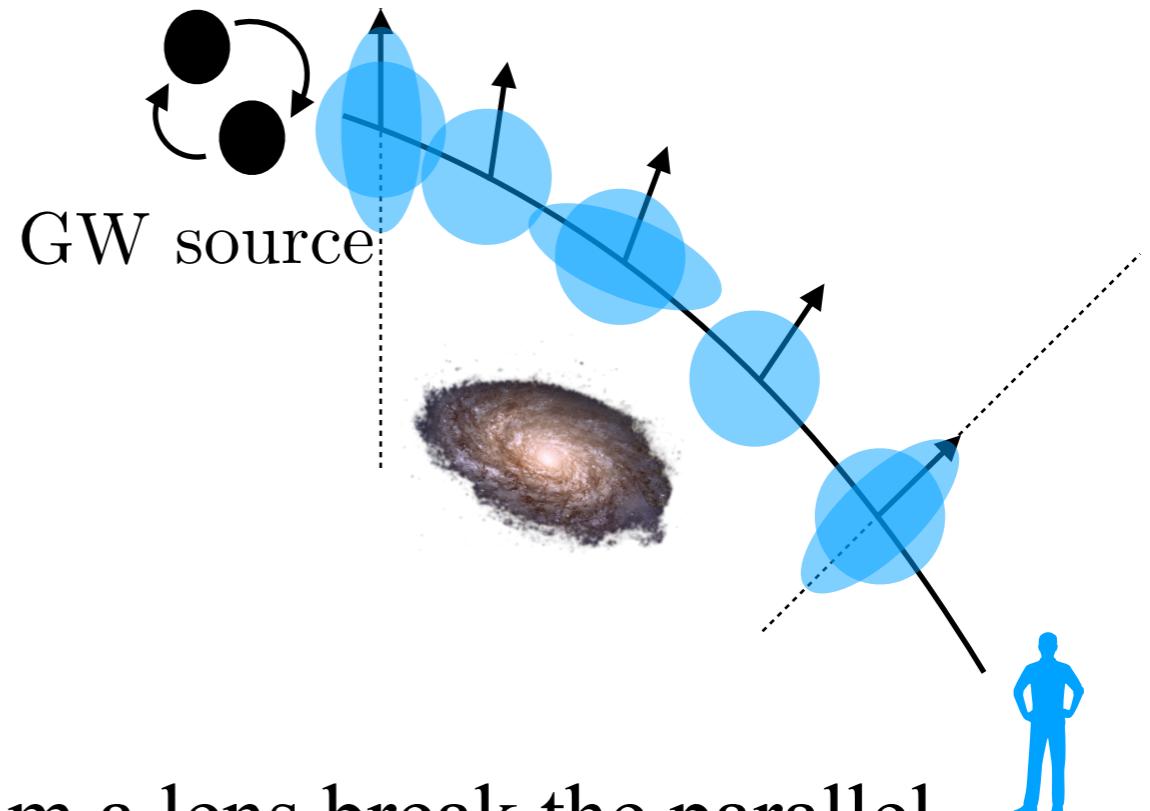
General Relativity curved space GW Eq. of motion

$$\bar{\nabla}^\alpha \bar{\nabla}_\alpha h_{\mu\nu} - \cancel{2\bar{R}_{\alpha\mu\nu\beta} h^{\alpha\beta}} = 0$$

Effective mass-like term
Short wavelength approximation



Parallel transport of + and × modes
in curved space



Does the presence of curvature from a lens break the parallel transport of GW polarizations?

Geometric optics and beyond

Giulia Cusin, Macarena Lagos, 2020

Wave ansatz

$$h_{\mu\nu} = \Re \left(\left(\underbrace{\varepsilon_{\mu\nu}^{(0)}}_{\text{amplitude}} + \underbrace{\omega^{-1} \varepsilon_{\mu\nu}^{(1)}}_{\text{amplitude correction}} + \dots \right) e^{i\omega\Phi} \right)$$

$k_\mu \equiv \Phi_{,\mu}$
wavevector

large ω



Equation of motion

$$0 = -\omega^2 k_\beta k^\beta \varepsilon_{\mu\nu}^{(0)} + i\omega^1 [2k^\beta \varepsilon_{\mu\nu;\beta}^{(0)} + k^\beta_{;\beta} \varepsilon_{\mu\nu}^{(0)}]$$
$$+ \omega^0 \left[i[2k^\beta \varepsilon_{\mu\nu;\beta}^{(1)} + k^\beta_{;\beta} \varepsilon_{\mu\nu}^{(1)}] + \square \varepsilon_{\mu\nu}^{(0)} - 2\varepsilon_{(0)}^{\alpha\beta} \bar{R}_{\alpha\mu\nu\beta} \right]$$

Solve order by order

Giulia Cusin, Macarena Lagos, 2020

$$\mathcal{O}(\omega^2)$$

$$k_\beta k^\beta = 0 \quad \text{GWs follow nul geodesics!}$$

$$\mathcal{O}(\omega^1) \quad [2k^\beta \varepsilon_{\mu\nu;\beta}^{(0)} + k^\beta{}_{;\beta} \varepsilon_{\mu\nu}^{(0)}] = 0$$

The polarization is parallel transported along k^μ

$$\mathcal{O}(\omega^0)$$

$$2k^\beta \varepsilon_{\mu\nu;\beta}^{(1)} + k^\beta{}_{;\beta} \varepsilon_{\mu\nu}^{(1)} = -i \left[2\varepsilon_{(0)}^{\alpha\beta} \bar{R}_{\alpha\mu\nu\beta} - \bar{\square} \varepsilon_{\mu\nu}^{(0)} \right]$$

The background gravitational field sources amplitude corrections of the GW!

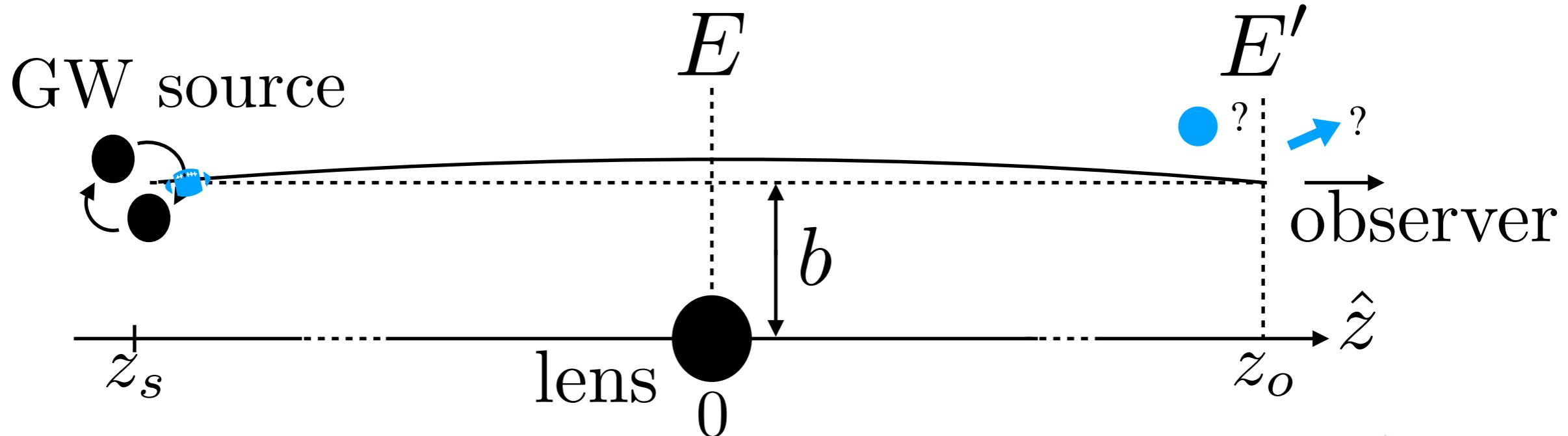
Point-like lens

C.D., Giulia Cusin, Macarena Lagos, 2021

Schematically,

$$\varepsilon_{\mu\nu}^{(1)}(\lambda) = \underbrace{\frac{\varepsilon_{\mu\nu}^{(1)}(\lambda_s)D(\lambda_s)}{D(\lambda)}}_{=0} - \frac{i}{D(\lambda)} \int_{\lambda_s}^{\lambda} d\lambda' D \left(\varepsilon_{(0)}^{\alpha\beta} \bar{R}_{\alpha\mu\nu\beta} - \frac{1}{2} \square \varepsilon_{\mu\nu}^{(0)} \right)$$

$$ds^2 = -(1 + 2\Psi)dt^2 + (1 - 2\Psi)d\mathbf{x}^2 \quad \Psi = \frac{R_s}{2\|\mathbf{x}\|}$$

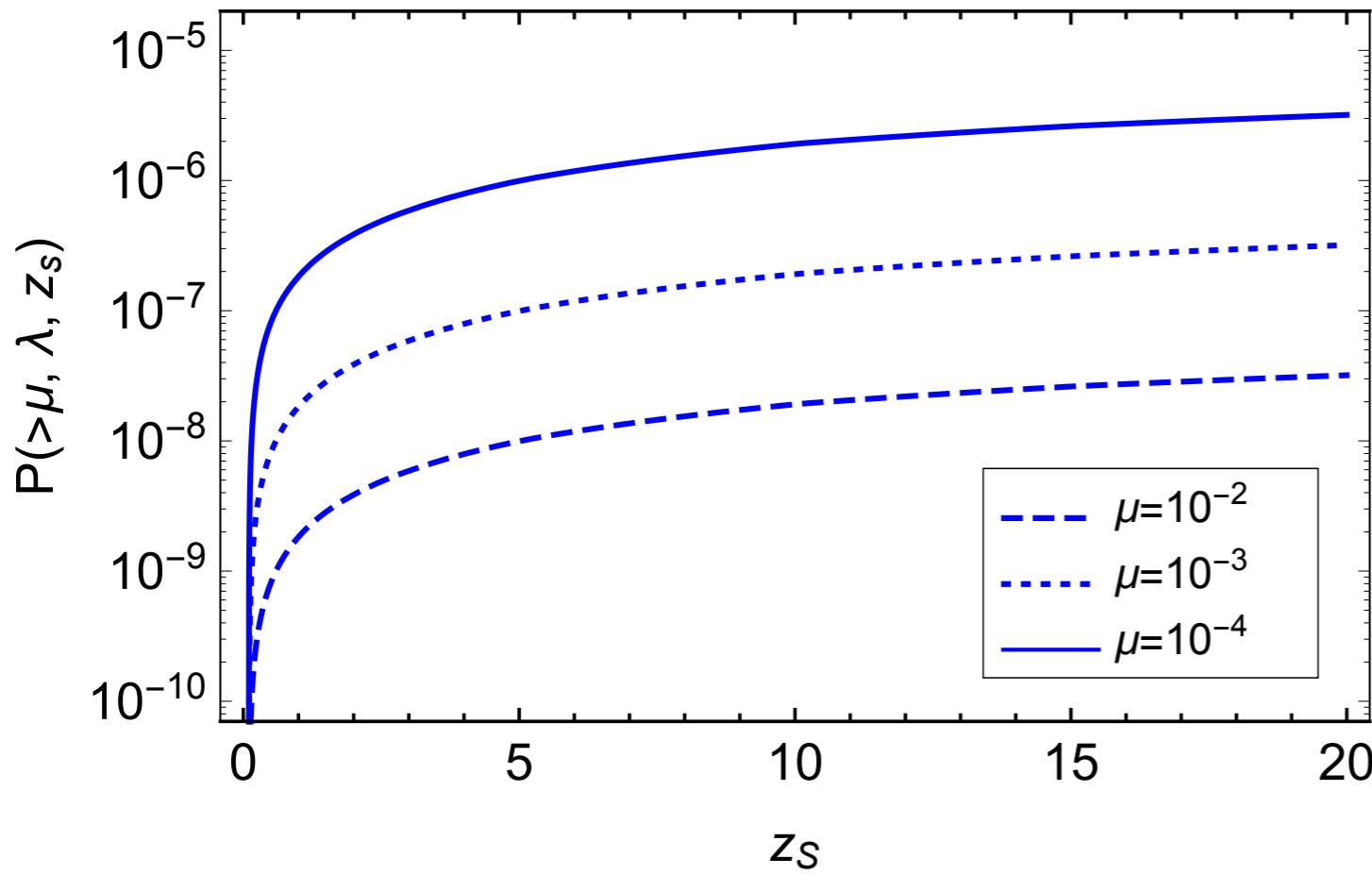


Some + and \times leak into scalar polarizations (c), (d) due to the presence of the lens!

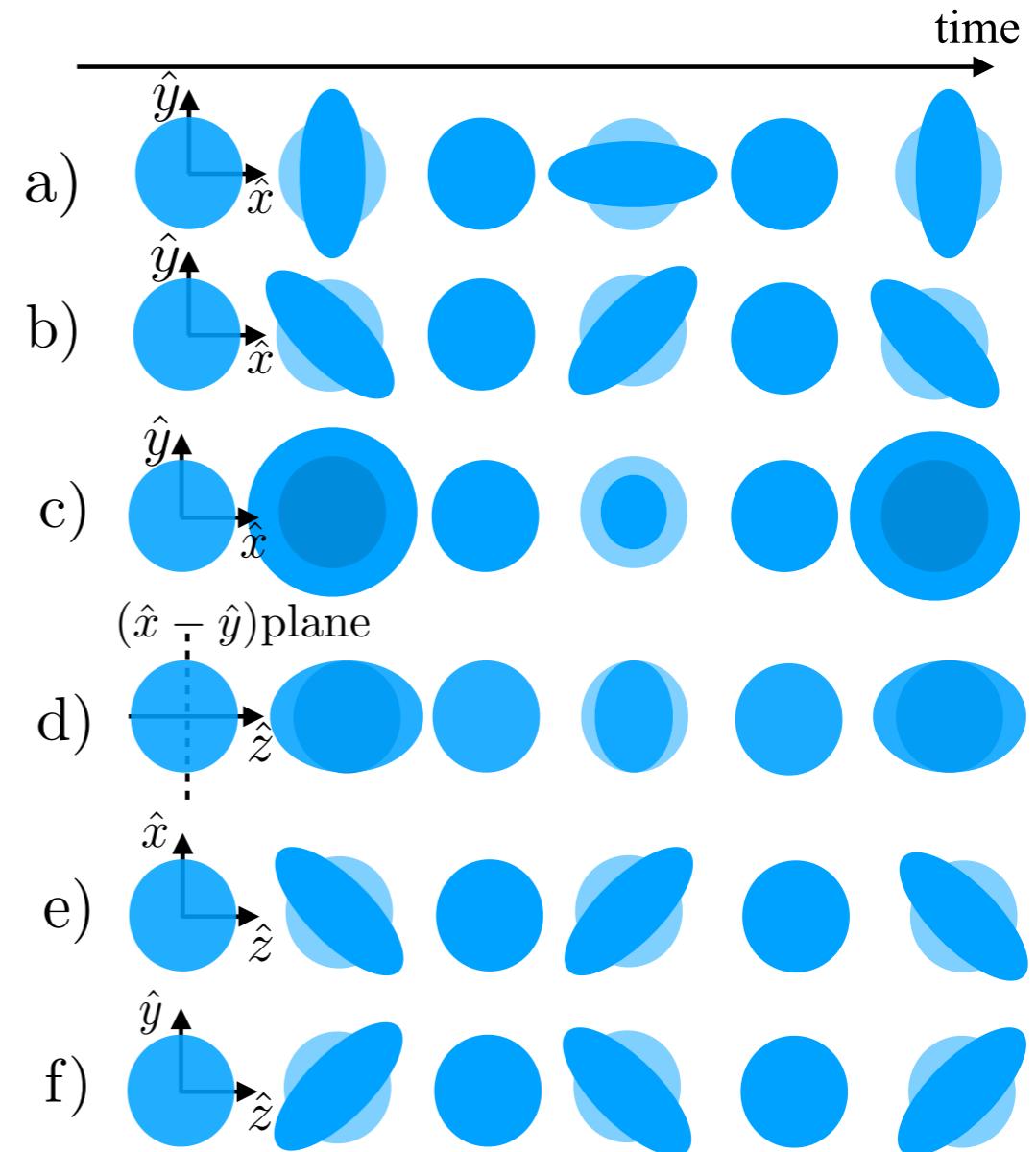
C.D., Giulia Cusin, Macarena Lagos, 2021

Probability to create scalar modes
 μ^{-1} smaller than the tensor modes

$$\frac{A_s}{A_t} = \frac{R_s \lambda}{b^2} = \mu \ll 1$$

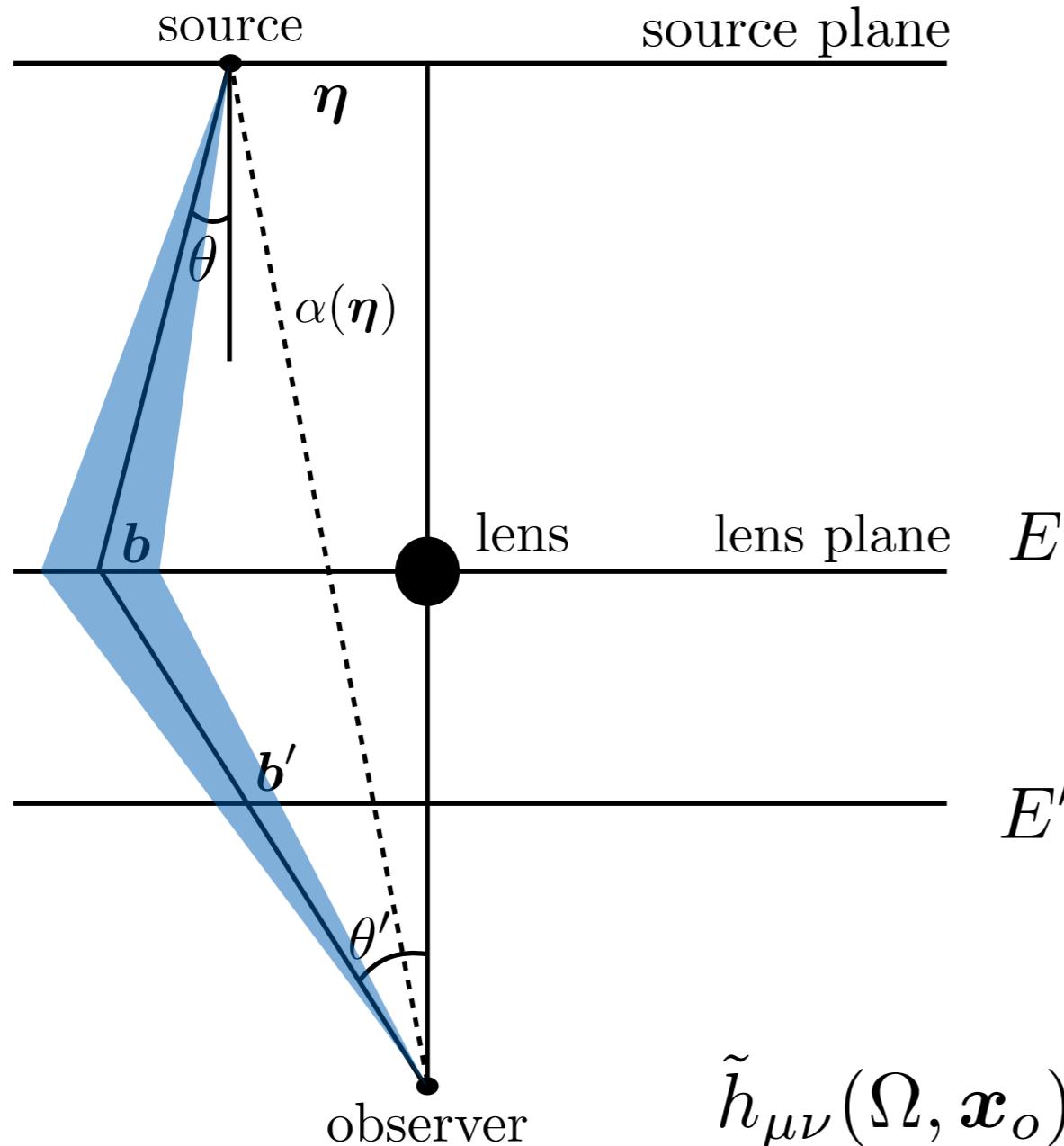


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Several paths interacting destructively?

C.D., Giulia Cusin, Macarena Lagos, 2021



Two effects compete:

- Geometric time delay
- Shapiro time delay

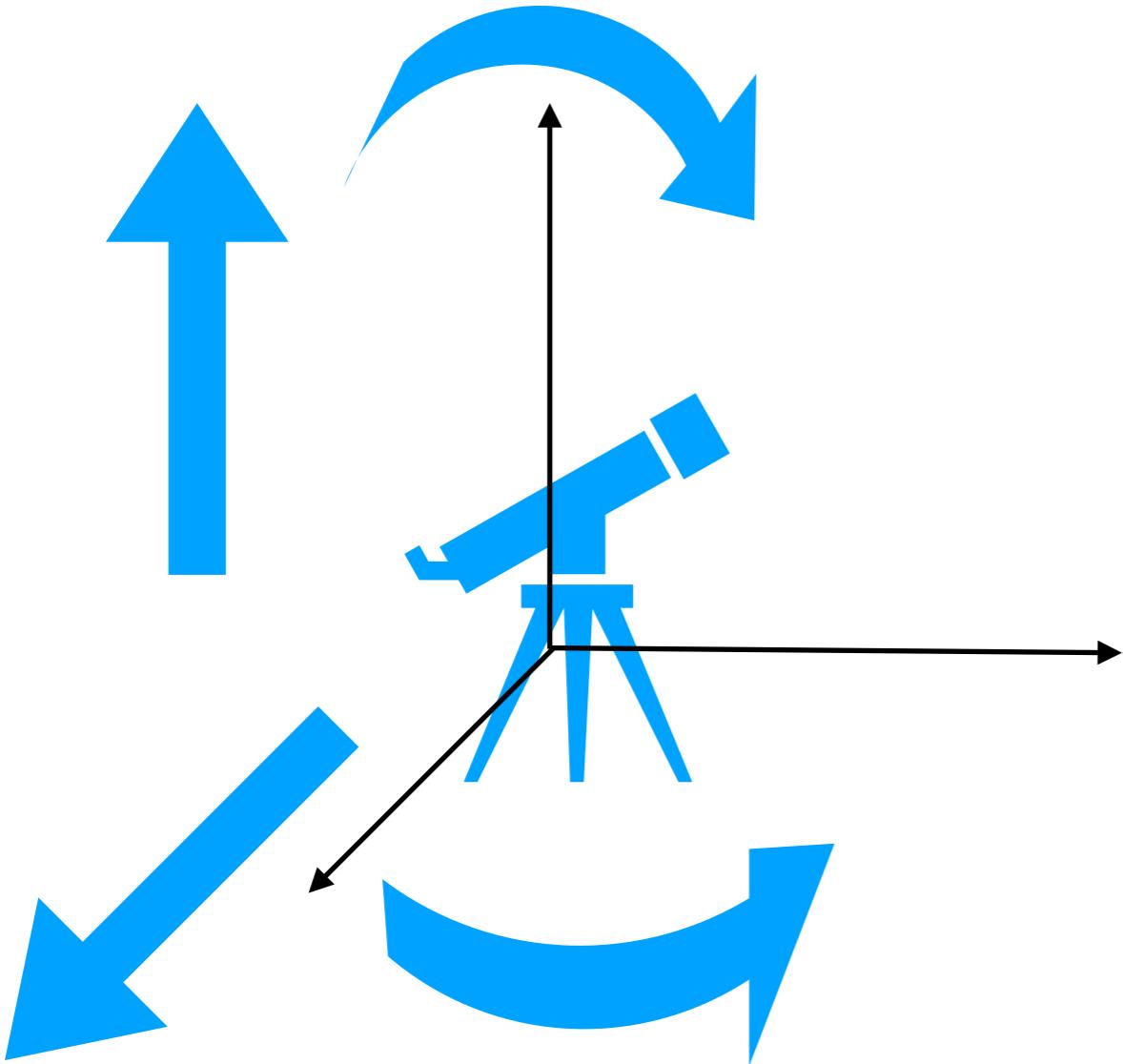
$$\tilde{h}_{\mu\nu}(\Omega, \mathbf{x}_o) = \underbrace{F_s \cdot F_{\mu\nu\alpha\beta}}_{\text{Standard amplification factor}} \tilde{h}_{\text{nolens}}^{\alpha\beta}(\Omega, \mathbf{x}_o)$$

polarization distortions

Observer dependence?

C.D., Giulia Cusin, Macarena Lagos, 2021

What about observers related by Lorentz transformations?



- 3 boosts
- 3 rotations
- Polarization content is observer dependent but there does not exist any observer who feels only the plus and cross polarizations.

Conclusion

- Parallel transport of the GW polarization is mostly a good approximation but it is not exact.
- Additional apparent polarization may appear in lensing scenarios, even in GR.

Thank you!

I am happy to answer any questions you may have :)