# Detecting new forces in the gravitational wave background 

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## The era of gravitational waves



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## This talk in one slide

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1. SMBH GW background is a guaranteed discovery

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2. Long range forces can detectably modify spectrum

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1. SMBH GW background is a guaranteed discovery
2. Long range forces can detectably modify spectrum
3. SMBH GWs potentially probe many BSM scenarios

## The stochastic GW background



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## The stochastic GW background


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## Stochastic background spectrum



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Gravitational waves drive the evolution of the binary

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f_{\text {orbit }}(r)=\left(\frac{G\left(M_{1}+M_{2}\right)}{4 \pi^{2} r^{3}}\right)^{1 / 2} \\
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\frac{\mathrm{d} E_{\mathrm{GW}}}{\mathrm{~d} f_{\mathrm{GW}}} \propto f^{-1 / 3} \Longrightarrow \frac{\mathrm{~d} h_{c}}{\mathrm{~d} f_{\mathrm{GW}}} \propto f^{-2 / 3} \\
{[\text { Phinney, 2001] }}
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New physics can break this prediction

## Assumptions are made to be broken

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(1) $f \leftrightarrow r$ relation

Kepler's law


## Assumptions are made to be broken

(1) $f \leftrightarrow r$ relation Kepler's law
(2) All energy loss is gravitational


## Benchmark model

## Additional dynamics spoil the $-2 / 3$

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Toy model: charge BHs* under a new long-range force

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1 New force changes Kepler's law
2 New radiation takes energy
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Additional dynamics spoil the $-2 / 3$
Toy model: charge BHs* under a new long-range force *or their surroundings
(1) New force changes Kepler's law

2 New radiation takes energy

## $\underline{\text { Charge parameters }}$

$$
\underbrace{\alpha=\frac{Q_{1} Q_{2}}{G M_{1} M_{2}}}_{\text {Force }} \quad \underbrace{\gamma^{2}=\frac{1}{G}\left(\frac{Q_{1}}{M_{1}}-\frac{Q_{2}}{M_{2}}\right)^{2}}_{\text {Radiation }}
$$

## Whence charge?

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(1) Particle production in SMBH environment

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Any of these could signal new fundamental physics

Toy model assumption:
charge is pointlike relative to binary separation
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## Modified single-source spectrum

New force and radiation modify the spectrum

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\text { ator mass } \\
F=\frac{G M_{1} M_{2}}{r^{2}}(1-\underbrace{\alpha e^{-m r}(1+m r)}_{\text {new force }})
\end{gathered}
$$

## Modified single-source spectrum

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\begin{aligned}
& \frac{\mathrm{d} E_{\mathrm{GW}}}{\mathrm{~d} f_{\mathrm{GW}}}=-\pi^{2} \mu r^{2} f_{\mathrm{GW}}\left(\frac{2 f_{\mathrm{GW}}}{r} \frac{\mathrm{~d} r}{\mathrm{~d} f_{\mathrm{GW}}}+1\right) \underbrace{\frac{P_{\mathrm{GW}}}{P_{\mathrm{GW}}+P_{\mathrm{new}}}}_{\text {radiation }} \\
& \text { ator mass } \\
& \mathrm{t}
\end{aligned}
$$

$$
P_{\text {new }}=\frac{1}{3} G \gamma^{2} \mu^{2} r^{2} \omega^{4} \operatorname{Re}\left[\sqrt{1-\frac{m^{2}}{\omega^{2}}}\right] \begin{cases}\left(1-\frac{m^{2}}{2 \omega^{2}}\right) & \text { (scalar) } \\ 2\left(1+\frac{m^{2}}{2 \omega^{2}}\right) & \text { (vector) }\end{cases}
$$

## Modifying the force law $(\alpha \neq 0)$



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## New dipole radiation $(\gamma \neq 0)$



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## From single sources to $h_{c}$

$$
h_{c}^{2}(f)=\int \mathrm{d} z \mathrm{~d} M_{1} \mathrm{~d} M_{2} \frac{\mathrm{~d} n_{\mathrm{G}}}{\mathrm{~d} z M_{1} \mathrm{~d} M_{2}} \frac{f_{\mathrm{s}}}{1+z} \frac{\mathrm{~d} E_{\mathrm{GW}}}{\mathrm{~d} f_{\mathrm{s}}} \frac{3 H_{0}^{2}}{2 \pi^{2} \rho_{c}^{2}}
$$

## From single sources to $h_{c}$

Single-source spectrum

## From single sources to $h_{c}$



Source distribution
Single-source spectrum

## Observables

Force law ( $|\alpha|>0$ )


Dipole radiation ( $|\gamma|>0$ )


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$f[\mathrm{eV}]$
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Force law ( $|\alpha|>0$ )
$f[\mathrm{eV}]$

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3 Sensitivity curves: this is happening NOW

## Current data

Interpret the NANOGrav 12.5-yr result in this framework

[Arzoumanian et al., 2020]

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## Beyond the benchmark

Probe any new physics that affects binary dynamics

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## Beyond the benchmark

Probe any new physics that affects binary dynamics
(1) Charged clouds

2 Dark matter spikes
(3) Superradiance

## Conclusions

Supermassive black holes are our new laboratories

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SGWB discovery
is imminent

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SGWB discovery is imminent

Long-range forces
are detectable

## Conclusions

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## Conclusions

## Supermassive black holes are our new laboratories



## References I

Z. Arzoumanian et al. The NANOGrav 12.5 yr Data Set: Search for an Isotropic Stochastic Gravitational-wave Background. Astrophys. J. Lett., 905(2):L34, 2020. doi: 10.3847/2041-8213/abd401.
C. J. Moore, R. H. Cole, and C. P. L. Berry. Gravitational-wave sensitivity curves. Class. Quant. Grav., 32(1):015014, 2015. doi: 10.1088/0264-9381/32/1/015014.
E. S. Phinney. A Practical theorem on gravitational wave backgrounds. 7 2001.
A. Sesana. Systematic investigation of the expected gravitational wave signal from supermassive black hole binaries in the pulsar timing band. Mon. Not. Roy. Astron. Soc., 433:1, 2013. doi: $10.1093 / \mathrm{mnrasl} / \mathrm{slt} 034$.

