

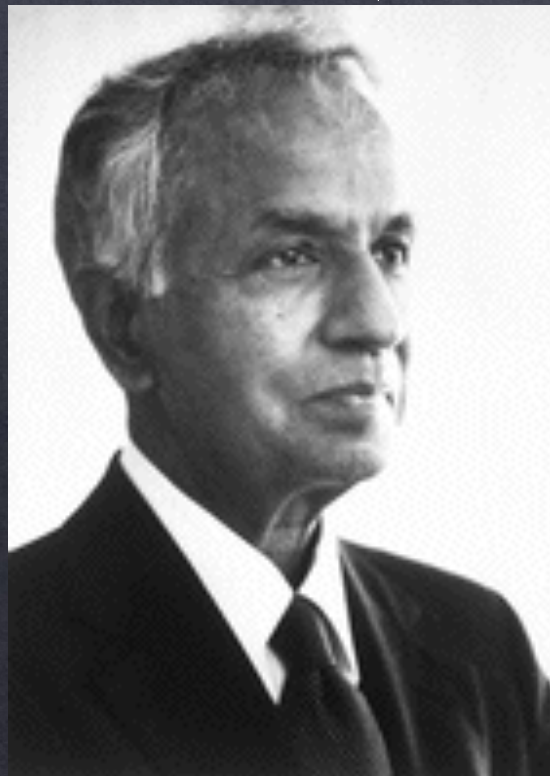
Symmetries of Black Hole Perturbation Theory

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w/ Lam Hui, Austin Joyce, Riccardo Penco, & Luca Santoni
arXiv:2105.01069
arXiv:210x.xxxxxx

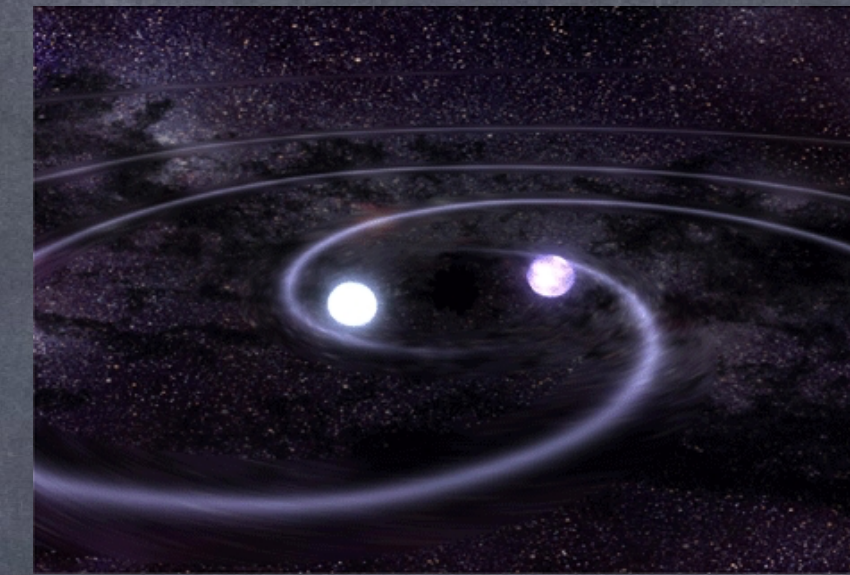
"The black holes of nature are the most perfect macroscopic objects there are in the universe: the only elements in their construction are our concepts of space and time."

-Chandrasekhar



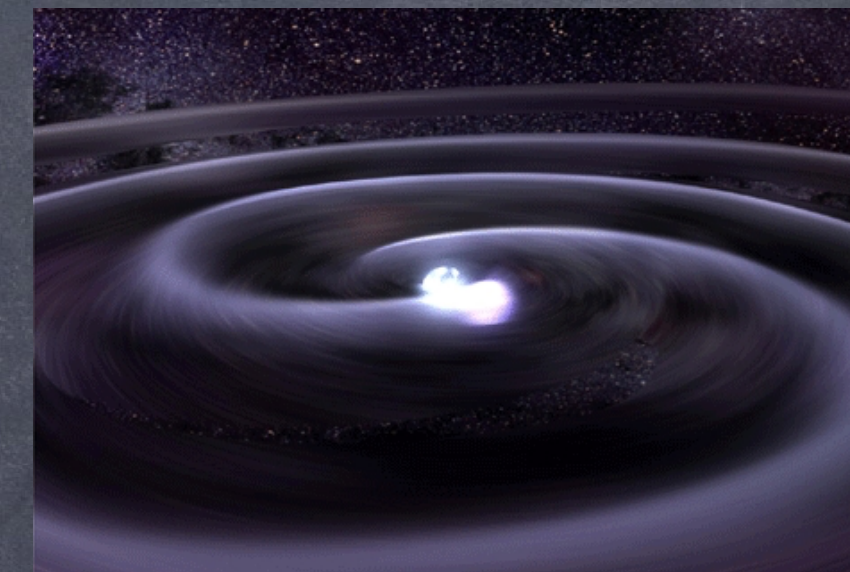
Motivation

- Black holes are **simple** but **mysterious**
- "A new era of gravitational wave astronomy"TM
- Binary BH mergers probe gravity in a **wide variety of regimes**
- This talk: **linear** perturbations; **Schwarzschild** for simplicity
 - Many/(all?) results apply to **Kerr**



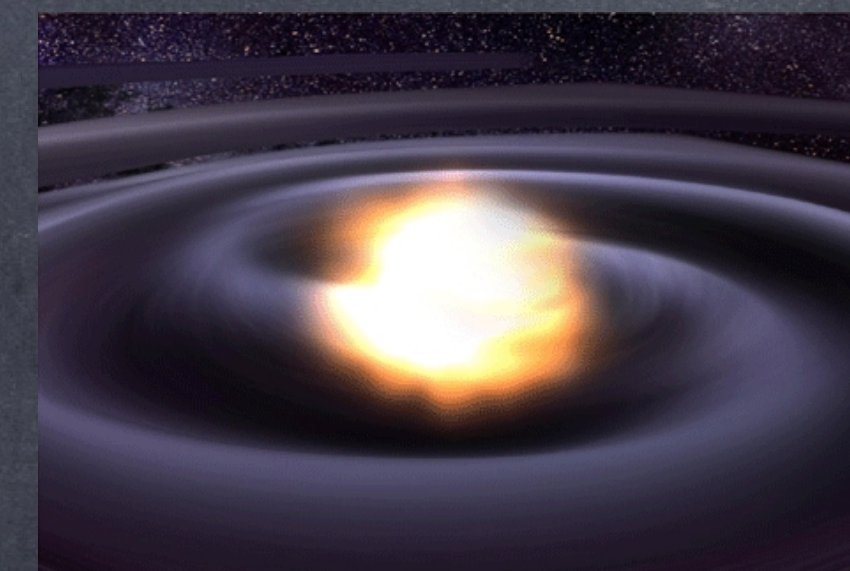
Inspiral

Post-Newtonian; EFT;
effective one-body;
perturbation theory



Merger

Numerical relativity



Ringdown

Perturbation theory

Image: NASA

Black Hole Spectroscopy

- Ringdown dominated by **quasinormal modes**
- QNMs: decaying waves with BCs:
 - Ingoing at horizon
 - Outgoing at infinity
- Discrete spectra: allows **spectroscopy à la atoms**
- Applications: no-hair (Isi+1905.00869), tests of GR, BH mimickers, ...

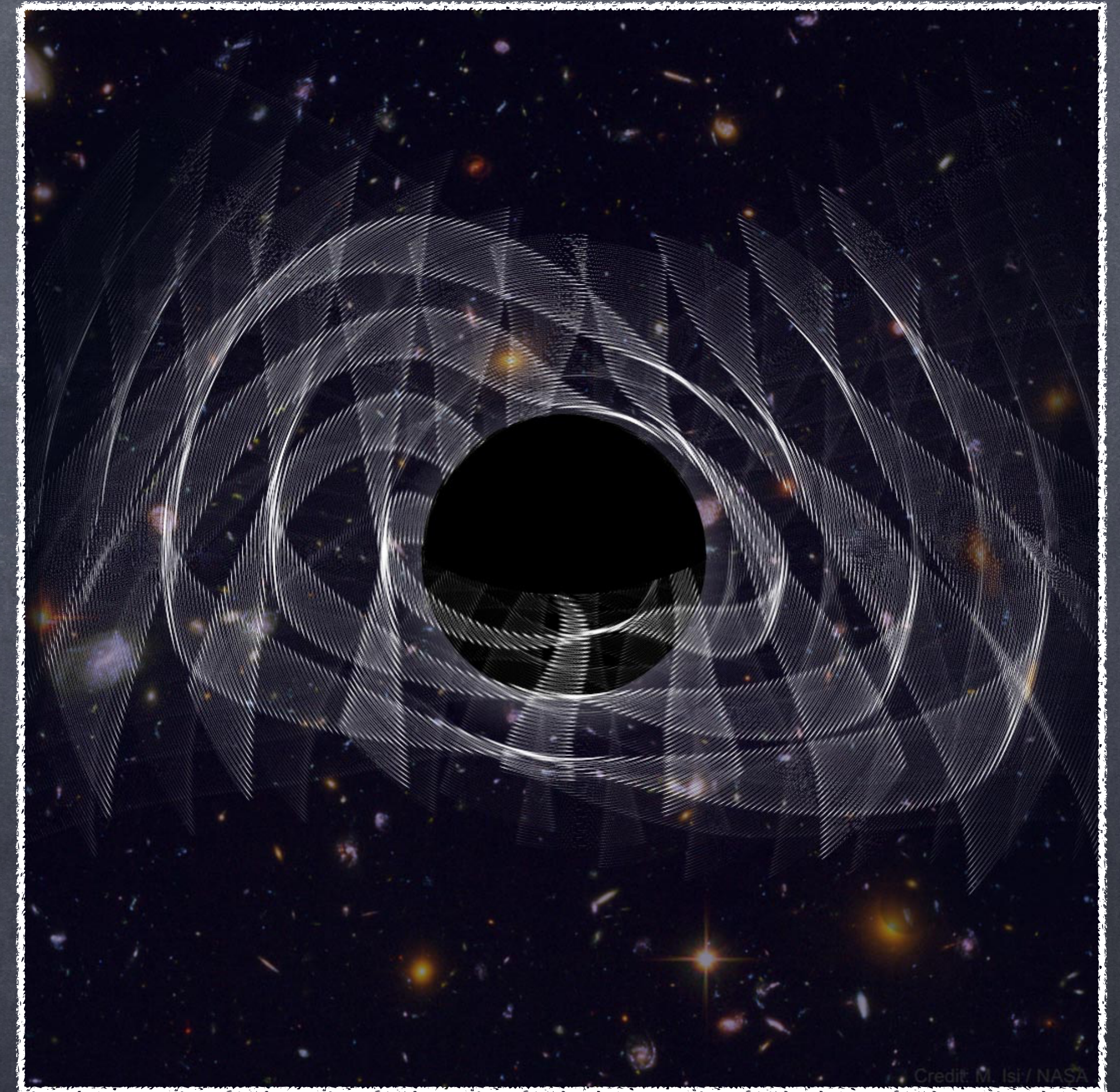


Image: M. Isi

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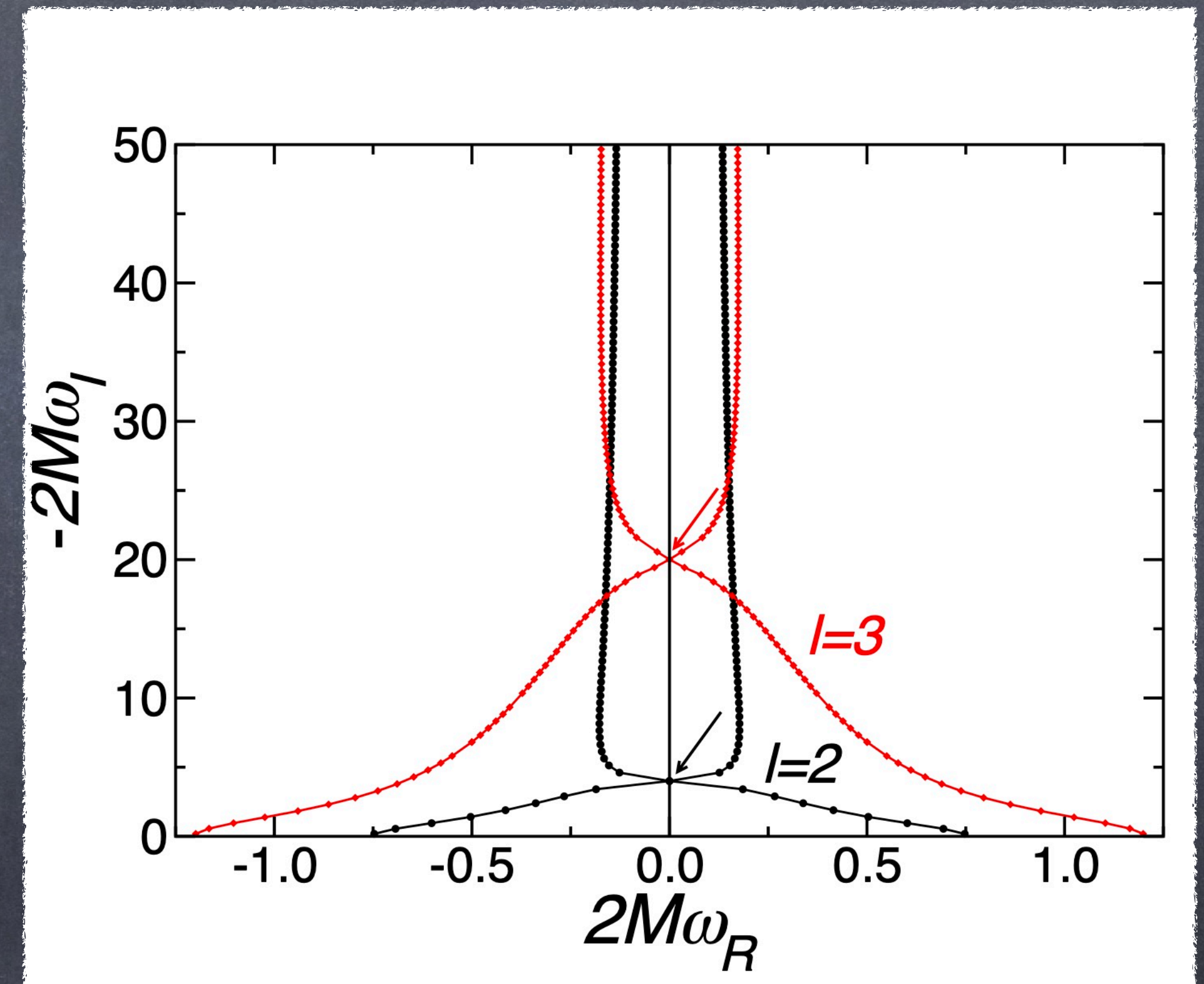


Figure: Berti, Cardoso, & Starinets (0905.2975, CQG review)

Tidal Responses & Love Numbers

- During inspiral, companion object sources a tidal field
- Gravitational response to tidal deformation encoded in Love numbers
 - Measures internal structure
 - Observable at **SPN**
- Wilson coefficients of point-particle EFT encoding finite size/structure:

$$S_{\text{pp}} = \int d\tau \left(-m + \frac{1}{2} \lambda_E E_{ij}^2 + \frac{1}{2} \lambda_B B_{ij}^2 + \dots \right)$$

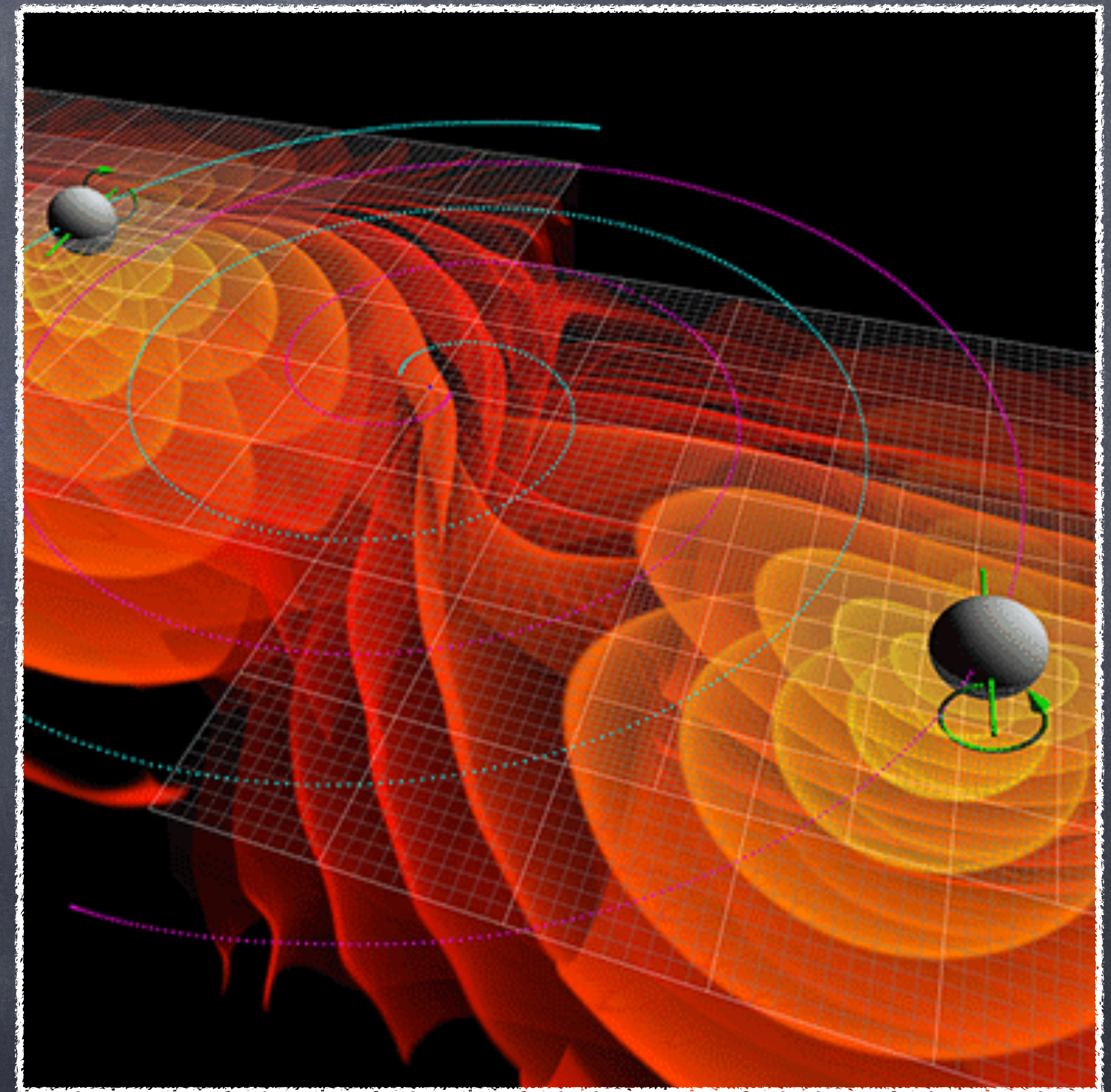


Image: C. Henze/NASA Ames Research Center

Hints of Symmetries

(in $D=4$ GR)

Symmetry of Love? Black hole Love numbers = 0

"Fine tuning": vanishing EFT coefficients (Porto, 1606.08895)

Symmetry of spectra? Both GW polarizations have the
same QNM spectrum

Caused by a duality of the Einstein equations

A Black Hole Perturbation Theory Primer

(for cosmologists)

Expand metric around background

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + \frac{1}{M_{\text{Pl}}} h_{\mu\nu}$$

Split by behavior under parity
(viz scalar/vector/tensor decomposition)

$$h_{\mu\nu} = h_{\mu\nu}^{\text{even}} + h_{\mu\nu}^{\text{odd}}$$

Decompose into ($m=0$) spherical harmonics
(viz Fourier expansion)

$$h_{\mu\nu}(t, r, \theta, \phi) = \sum_{\ell} h_{\mu\nu}^{\ell}(t, r) \Theta(\theta)$$

$\nearrow Y_{\ell 0}, \sin \theta \partial_{\theta} Y_{\ell 0}$

Graviton has 2 dof: 1 in even sector and 1 in odd

At infinity, these correspond to $+/\times$ polarizations

Regge-Wheeler and Zerilli

- Encode dynamical d.o.f. in **master variables**:

- Even: Zerilli Ψ_+

- Odd: Regge-Wheeler Ψ_-

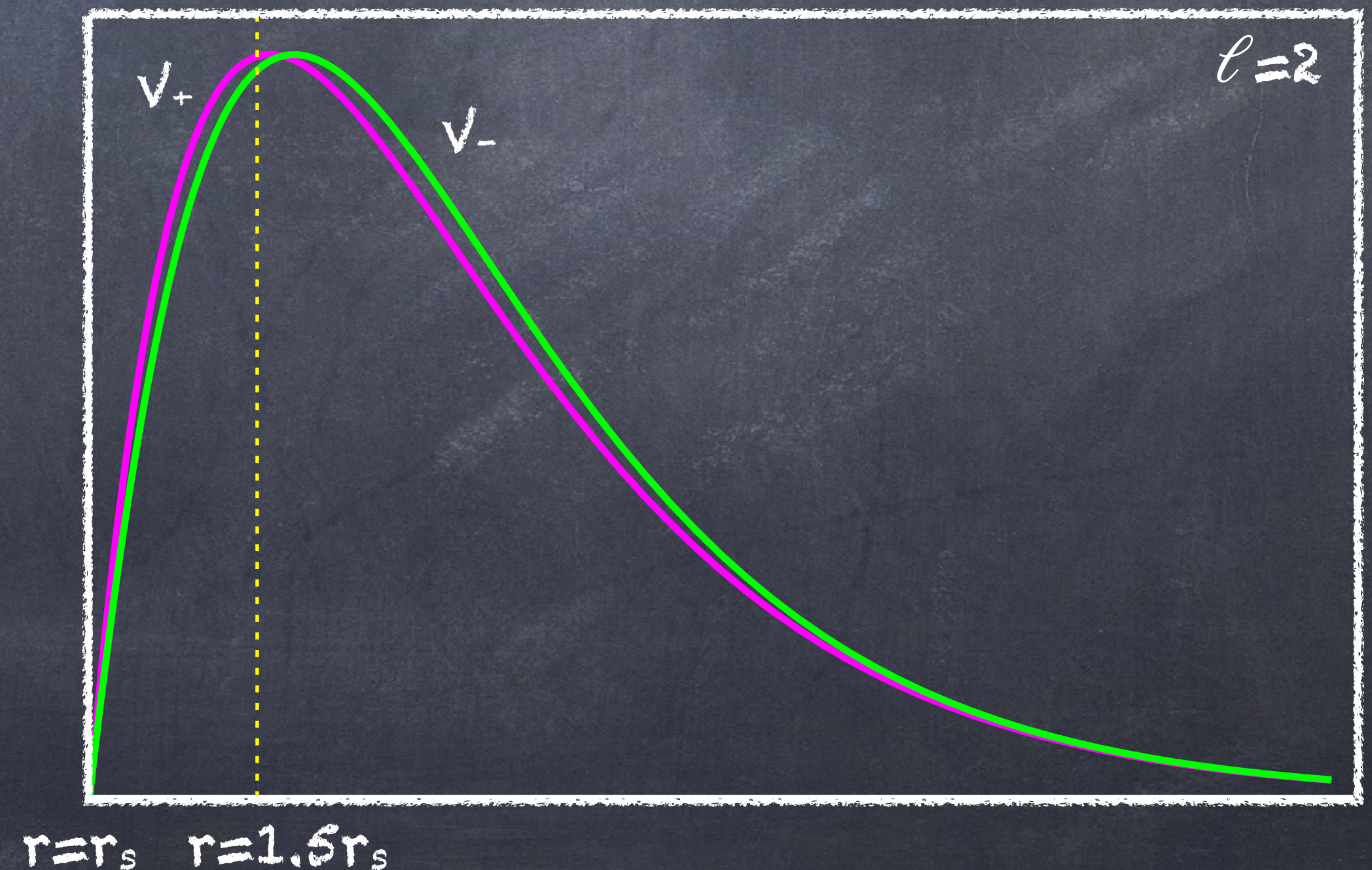
- These obey simple **Schrödinger-like** equations

- Tortoise coordinate:

$$dr \equiv \left(1 - \frac{r_s}{r}\right) dr_\star$$

Master equation:

$$\frac{\partial^2 \Psi_\pm(t, r)}{\partial t^2} - \frac{\partial^2 \Psi_\pm(t, r)}{\partial r_\star^2} + V_\pm(r) \Psi_\pm(t, r) = 0$$



Chandrasekhar's Duality

- Secret Link between Regge-Wheeler and Zerilli potentials:

$$V_{\pm}(r) = W^2(r) \mp \frac{dW(r)}{dr_{\star}} + \beta$$

with $W(r)$ the **superpotential** and β a constant

- This directly implies **isospectrality**
Chandrasekhar (1980s)

$$V_+ = \frac{1 - \frac{r_s}{r}}{r^3} \frac{9r_s^3 + 12\lambda^2 r_s r^2 + 8\lambda^2(1 + \lambda)r^3 + 18\lambda r_s^2 r}{(2\lambda r + 3r_s)^2} \quad 2\lambda \equiv (\ell - 1)(\ell + 2)$$

$$V_- = \left(1 - \frac{r_s}{r}\right) \left(\frac{\ell(\ell + 1)}{r^2} - \frac{3r_s}{r^3}\right)$$

Our question: where does this property in GR come from?

Linearized Einstein-Hilbert

To study symmetries, we want to work at the level of the **action**:

$$S = \frac{M_{\text{Pl}}^2}{2} \int d^4x \sqrt{-g} R \Big|_{g_{\mu\nu} = \bar{g}_{\mu\nu} + \frac{1}{M_{\text{Pl}}} h_{\mu\nu}}$$

Why? Allows us to calculate
e.g. Noether currents

Procedure

1. Fix gauge
2. Integrate out **non-dynamical** (auxiliary) dofs
3. Canonically normalize:
rescale **field** and **coordinates**

parity and spherical harmonics decoupling:

$$S = \sum_{\pm} \sum_{\ell} \int dt dr \mathcal{L}_{\pm}^{\ell}(t, r)$$

Linearized Einstein-Hilbert

Canonical normalization

- Coordinate choice: $r \rightarrow r_\star$
- Field redefinition: $h_{\mu\nu}^\pm \rightarrow \Psi_\pm$

End result

$$S = \frac{1}{2} \sum_{\ell=2}^{\infty} \int dt dr_\star \sum_{\pm} \left[\left(\frac{\partial \Psi_\pm}{\partial t} \right)^2 - \left(\frac{\partial \Psi_\pm}{\partial r_\star} \right)^2 - V_\pm(r) \Psi_\pm^2 \right]$$

Chandrasekhar Duality in Action

Replacing the potentials with the superpotential:

$$S = \frac{1}{2} \sum_{\ell=2}^{\infty} \int dt dr_{\star} \sum_{\pm} \left[\left(\frac{\partial \Psi_{\pm}}{\partial t} \right)^2 - \left(\frac{\partial \Psi_{\pm}}{\partial r_{\star}} \right)^2 - \left(W^2 \mp \frac{dW}{dr_{\star}} + \beta \right) \Psi_{\pm}^2 \right]$$

the action is invariant under

$$\delta \Psi_{\pm} = \left(\frac{\partial}{\partial r_{\star}} \mp W(r) \right) \Psi_{\mp}$$

We see that Chandrasekhar duality is a **symmetry** of Einstein-Hilbert (i.e., **off-shell**)

(NB: can also write symmetry at the level of metric perturbations)

Application: Tidal Response

- Noether current for static solutions:

$$J^{r_\star} = \partial_{r_\star} \Psi_+ \partial_{r_\star} \Psi_- + W(\Psi_+ \partial_{r_\star} \Psi_- - \Psi_- \partial_{r_\star} \Psi_+) - (W^2 + \beta) \Psi_+ \Psi_- = \text{const.}$$

- Regularity at the horizon: $J^{r_\star} = 0$
- At infinity: $\Psi_\pm \propto r^{\ell+1} + \lambda_\pm r^{-\ell}$; $J^{r_\star} \propto (\lambda_+ - \lambda_-)$
- Duality implies equal Love numbers
- Not quite vanishing, but helpful: odd sector much simpler than even

Minkowski Limit

Electric-magnetic duality

Duality in the flat limit: $\delta\Psi_+ = -\Psi_-$
 $SO(2)$ $\delta\Psi_- = \Psi_+$

implies, on-shell,

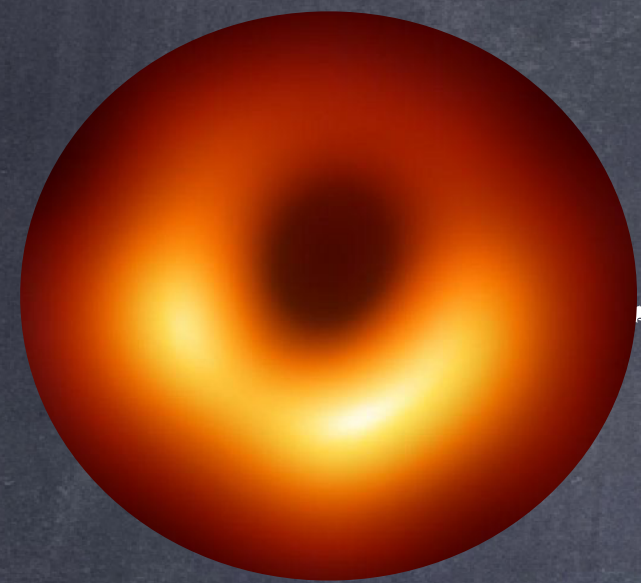
Gravitational electric-magnetic duality:

$$\begin{aligned}\delta R_{\mu\nu\alpha\beta} &= \tilde{R}_{\mu\nu\alpha\beta} \\ \delta \tilde{R}_{\mu\nu\alpha\beta} &= -R_{\mu\nu\alpha\beta}\end{aligned}$$
$$\tilde{R}_{\mu\nu\alpha\beta} = \frac{1}{2}\epsilon_{\mu\nu\rho\sigma}R^{\rho\sigma}{}_{\alpha\beta}$$

EM duality previously known symmetry of Einstein-Hilbert
around Minkowski, $(A)dS$

Vanishing Love, Mathematically

static Regge-Wheeler: $\partial_{r_\star}^2 \Psi - V(r)\Psi = 0$



Horizon

Infinity

regular

$\psi \sim \text{const.}$

~~blowing up~~

$\psi \sim \ln(1 - r_s/r)$

growing $\psi \sim r^\ell$ (tidal field)

decaying $\psi \sim \frac{1}{r^{\ell+1}}$ (static response)

Love number: coefficient of decaying term

Math: why do these diff eqs have this property?

Physics: why does GR give us diff eqs like this?

See also Charalambous,
Dubovsky, and Ivanov
(2103.01234)

Spin-0 Warm-Up

Klein-Gordon on Schwarzschild: $\nabla^2 \phi = 0 \longrightarrow \sum_{\ell} H_{\ell} \phi_{\ell} = 0$

$$H_{\ell} = -\Delta \left[\partial_r (\Delta \partial_r) - \ell(\ell + 1) \right] \quad \Delta(r) \equiv r(r - r_s)$$

Admits **raising** and **lowering** operators: $D_{\ell}^{+} \equiv -\Delta \partial_r + \frac{\ell + 1}{2}(r_s - 2r)$
 $D_{\ell}^{-} \equiv \Delta \partial_r + \frac{\ell}{2}(r_s - 2r)$

in the sense that $H_{\ell+1}(D_{\ell}^{+} \phi_{\ell}) = 0$

$$H_{\ell-1}(D_{\ell}^{-} \phi_{\ell}) = 0$$

Note: This example
contains the salient
features of spin-s on Kerr

Turning the Ladder Sideways

- The raising/lowering symmetry D_ℓ^\pm is unusual: it relates solutions at different levels ℓ

- **Want:** a symmetry for each level individually

- **Strategy:** lower to $\ell = s$ and use translation symmetry

- Horizontal symmetry: $\delta\phi_\ell = Q_\ell\phi_\ell$

$$Q_0 \equiv \Delta\partial_r$$

$$Q_1 \equiv D_0^+ Q_0 D_1^-$$

- **Conserved charge:** $P_\ell \equiv \Delta\partial_r(D_1^- D_2^- \cdots D_\ell^- \phi_\ell)$

$$P_0 = \Delta\partial_r\phi_0$$

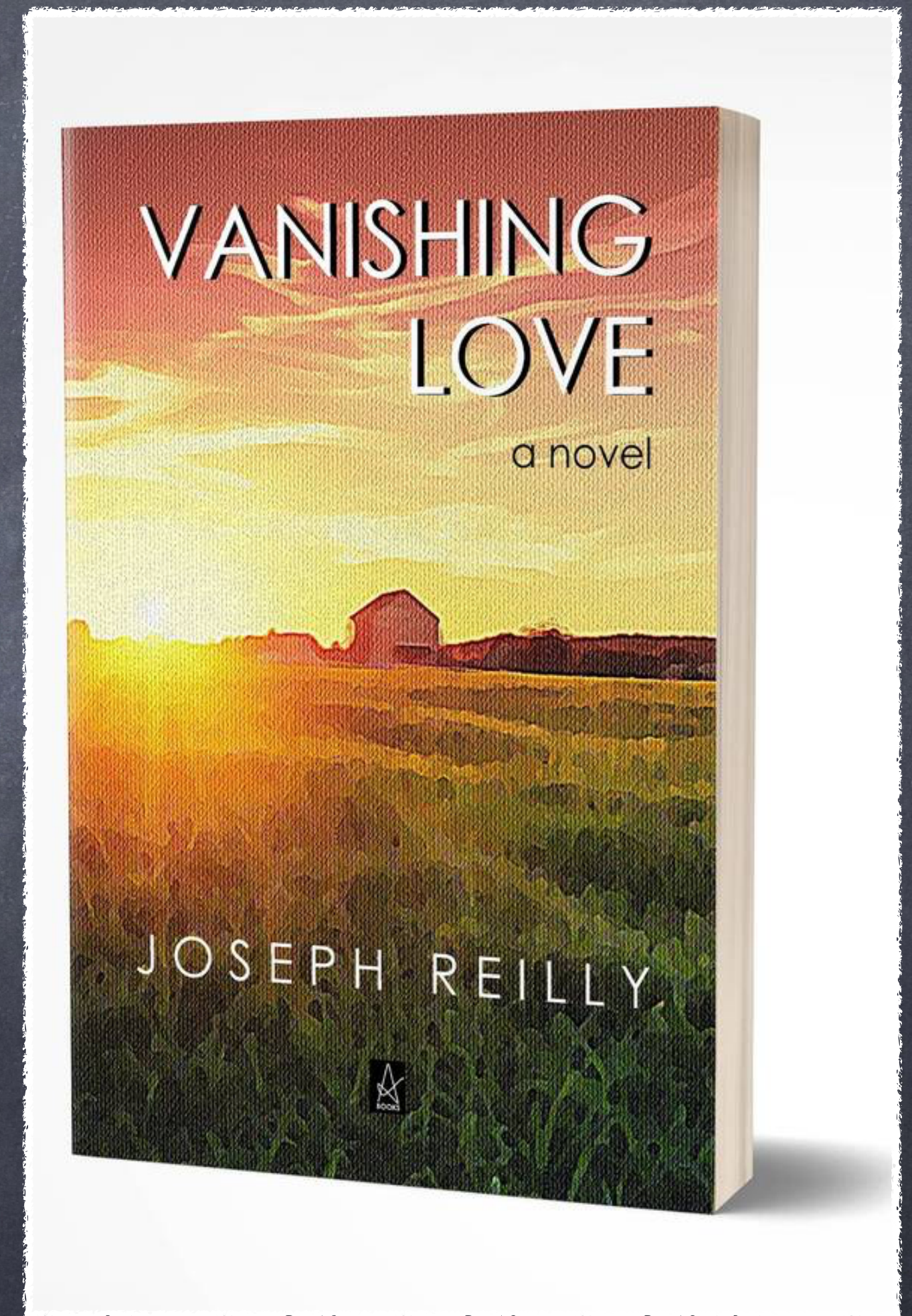
$$\vdots$$

$$Q_\ell \equiv D_{\ell-1}^+ Q_{\ell-1} D_\ell^-$$

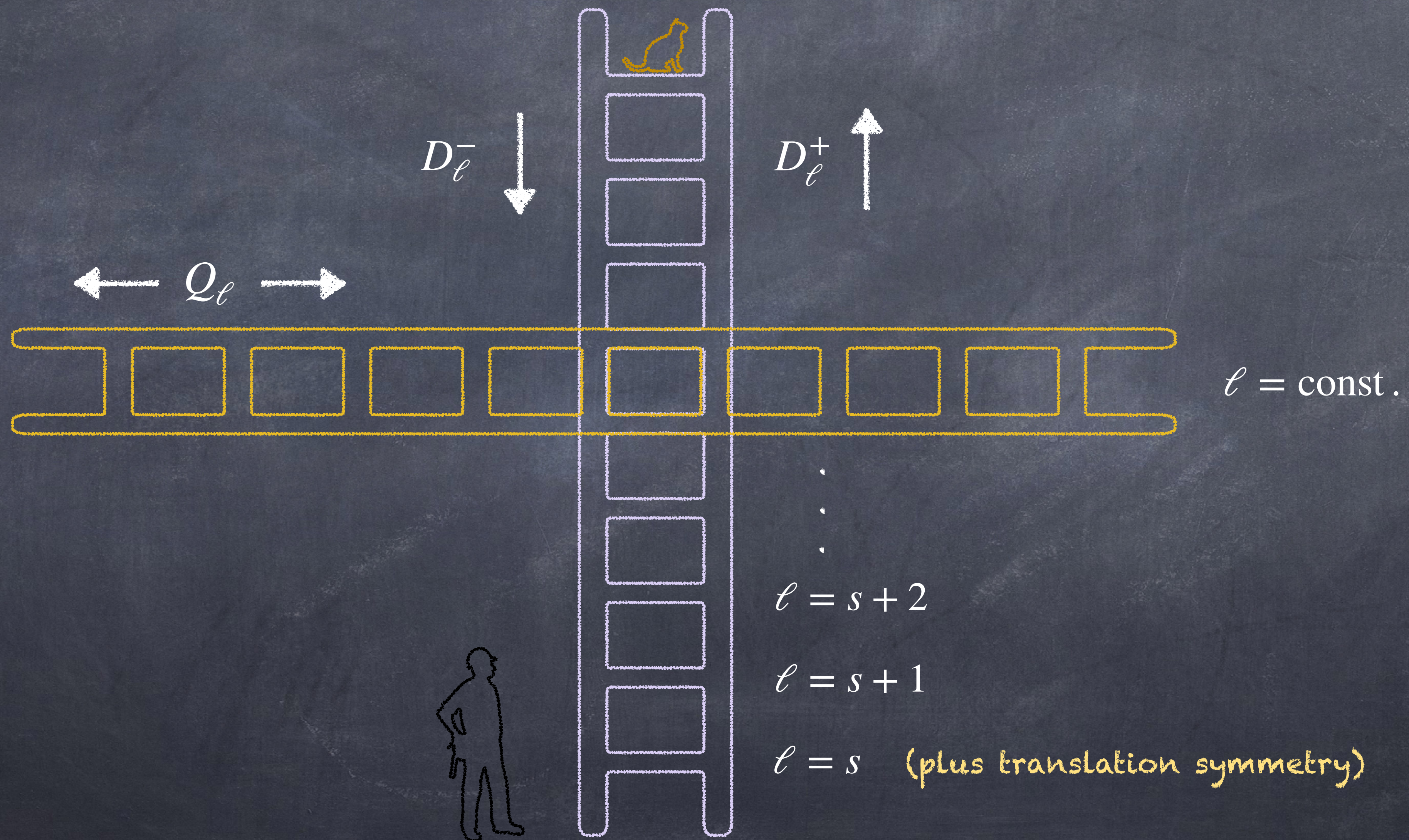
See also Compton and Morrison (2003.08023)

Vanishing Love

- Charge conservation implies $\text{Love} = 0$
- $P_\ell = 0$ for growing (infinity) and constant (horizon) modes
- $P_\ell \neq 0$ for decaying (infinity) and divergent (horizon) modes
- Spontaneous symmetry breaking:
 $Q_\ell \phi_\ell^{(g)} = 0, \quad Q_\ell \phi_\ell^{(d)} \neq 0$
- **Conclusion:** a decaying term diverges at the horizon
 - This also implies no (static) hair



Ladders Up and Down



Geometric Interpretation

From Schwarzschild to AdS

Static scalar: $S = \frac{1}{2} \int d^3x \sqrt{g} \phi \square \phi$

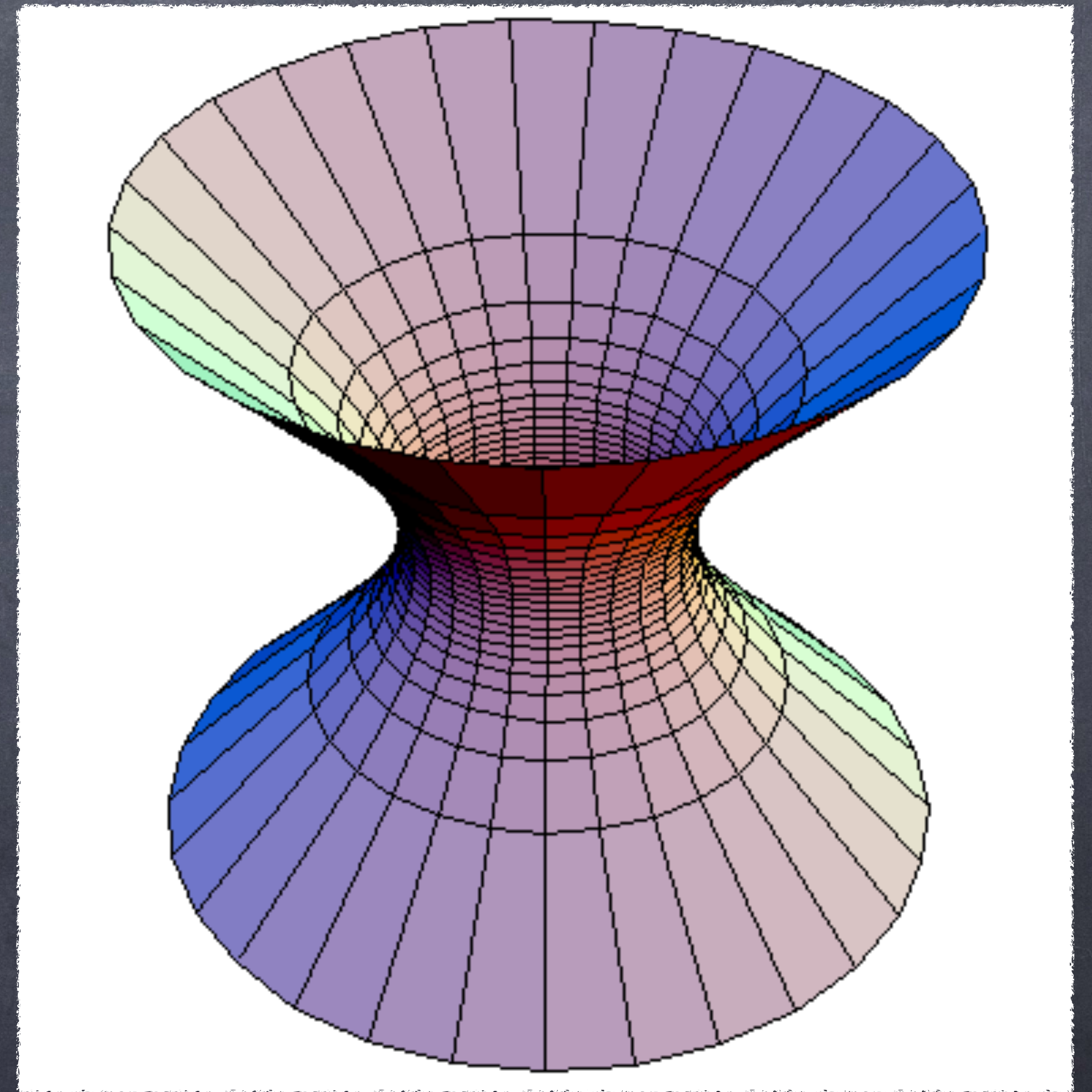
Conformal transformation:

$$\tilde{g}_{ij} = L^4 \Delta^{-2} g_{ij}, \quad \tilde{\phi} = L^{-1} \sqrt{\Delta} \phi \quad L: \text{arbitrary scale}$$

so that

$$S = \frac{1}{2} \int d^3x \sqrt{\tilde{g}} \left(\tilde{\phi} \tilde{\square} \tilde{\phi} + \frac{r_s^2}{4L^4} \tilde{\phi}^2 \right)$$

The metric \tilde{g}_{ij} is nothing other than AdS_3



Geometric Interpretation

From Schwarzschild to AdS

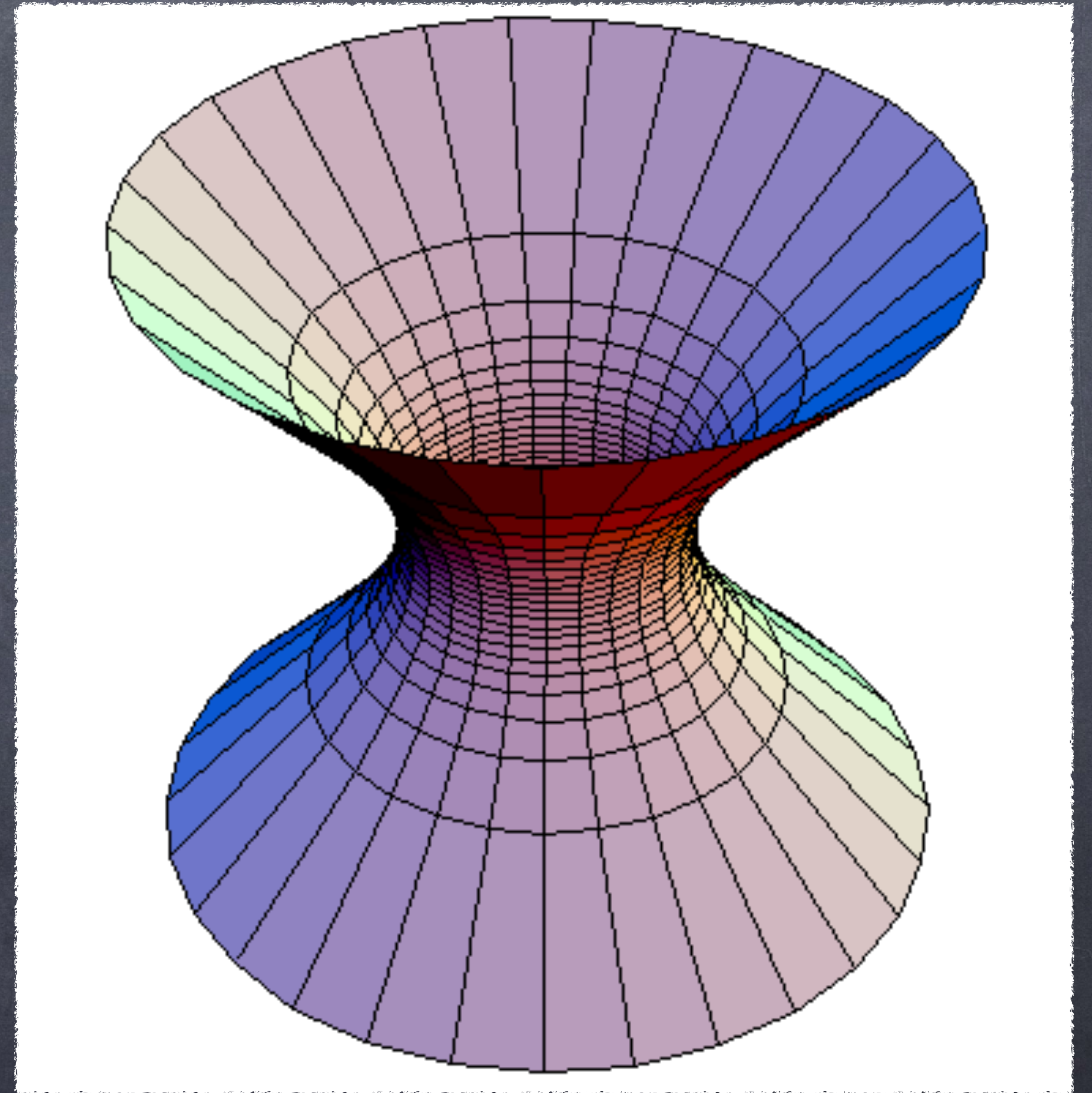
$$S = \frac{1}{2} \int d^3x \sqrt{\tilde{g}} \left(\tilde{\phi} \tilde{\square} \tilde{\phi} + \frac{r_s^2}{4L^4} \tilde{\phi}^2 \right)$$

Killing vectors of AdS_3 generate isometry:

$$\delta\phi = -2\Delta \cos\theta \partial_r \phi + (r_s - 2r) \partial_\theta (\sin\theta \phi)$$

Decompose in spherical harmonics:

$$\delta\phi_\ell \sim D_{\ell-1}^+ \phi_{\ell-1} + D_{\ell+1}^- \phi_{\ell+1}$$



To Kerr: Spin Ladder

Teukolsky equation:

$$\partial_r \left(\Delta \partial_r \phi_{\ell}^{(s)} \right) + s(2r - r_s) \partial_r \phi_{\ell}^{(s)} + \left(\frac{a^2 m^2 + is(2r - r_s)am}{\Delta} - (\ell - s)(\ell + s + 1) \right) \phi_{\ell}^{(s)} = 0$$

Admits ladders in ℓ and spin!

$$E^+ \equiv \partial_r, \quad E_s^- \equiv \Delta \partial_r - s(r_+ + r_- - 2r) - 2iam \frac{r_+ - r_-}{r_+ - r_-}$$

Relates solutions to Klein-Gordon, Maxwell, and Einstein

vanishing scalar Love \longrightarrow vanishing gravitational Love

NB: also have ladders for Regge-Wheeler and Teukolsky for more direct proof

IR symmetries

Point-particle EFT:

treat BH as a point, encode structure in higher-order operators

γ : worldline einbein
 g : (monopole) scalar charge
 λ_ℓ : Love numbers

$$S = -\frac{1}{2} \int d^4x (\partial\phi)^2 + \int d\tau \gamma \left[\frac{1}{2} \gamma^{-2} \dot{x}^\mu \dot{x}_\mu - \frac{\mu^2}{2} - g\phi + \sum_{\ell=0}^{\infty} \frac{\lambda_\ell}{2\ell!} \left(\partial_{(a_1} \cdots \partial_{a_\ell)_T} \phi \right)^2 \right]$$

UV symmetry in flat-space limit: $\delta\phi = r^2 \cos\theta \partial_r \phi + r \partial_\theta (\sin\theta \phi)$

Punchline: only the bulk $(\partial\phi)^2$ term is invariant

Summary

- Vanishing Love numbers and **isospectral QNMs** both indicate hidden symmetries of GR (/massless fields on Schwarzschild)
- We find **symmetries of Einstein-Hilbert** underlying these:
 - QNMs: EM duality on Schwarzschild
 - Love = 0: Ladder symmetries (shift sym + conformal Killing vector/ladder syms)

