

LYMAN- α CONSTRAINTS on COSMIC HEATING from DARK MATTER ANNIHILATION and DECAY

Wenzer Qin

with Greg Ridgway, Hongwan Liu, and Tracy Slatyer

based on [2008.01084]

MIT Center for Theoretical Physics, BSM Journal Club

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Massachusetts
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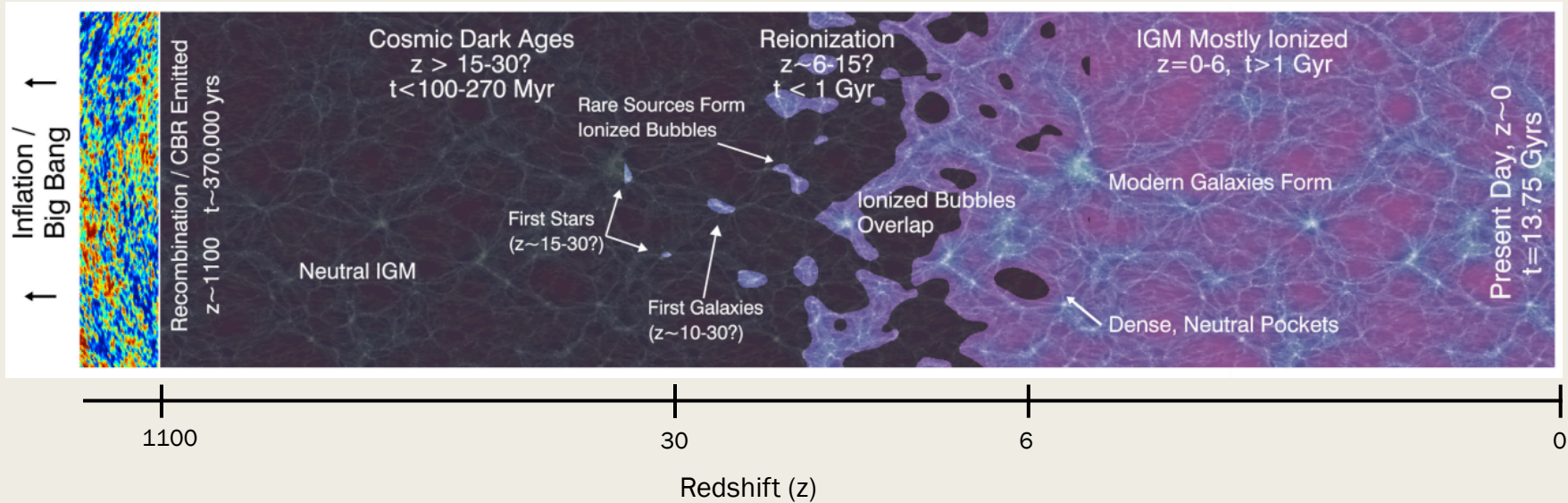
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Outline

- Early universe
- Effect of dark matter (DM) on ionization and temperature
- **DarkHistory** code package
 - *Choosing a reionization model*
- Data
 - *Reionization histories from Planck 2018*
 - *IGM temperature measurements from Ly α*
- New constraints
- Conclusion

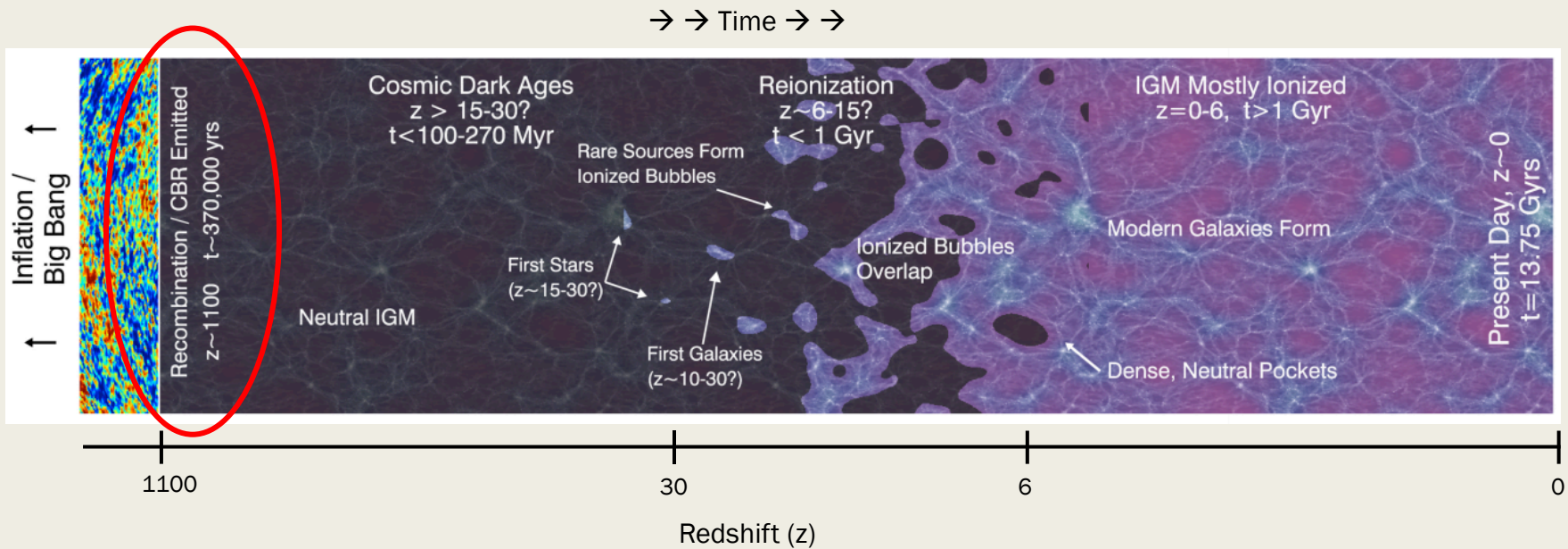
Timeline of the early universe

→ → Time → →



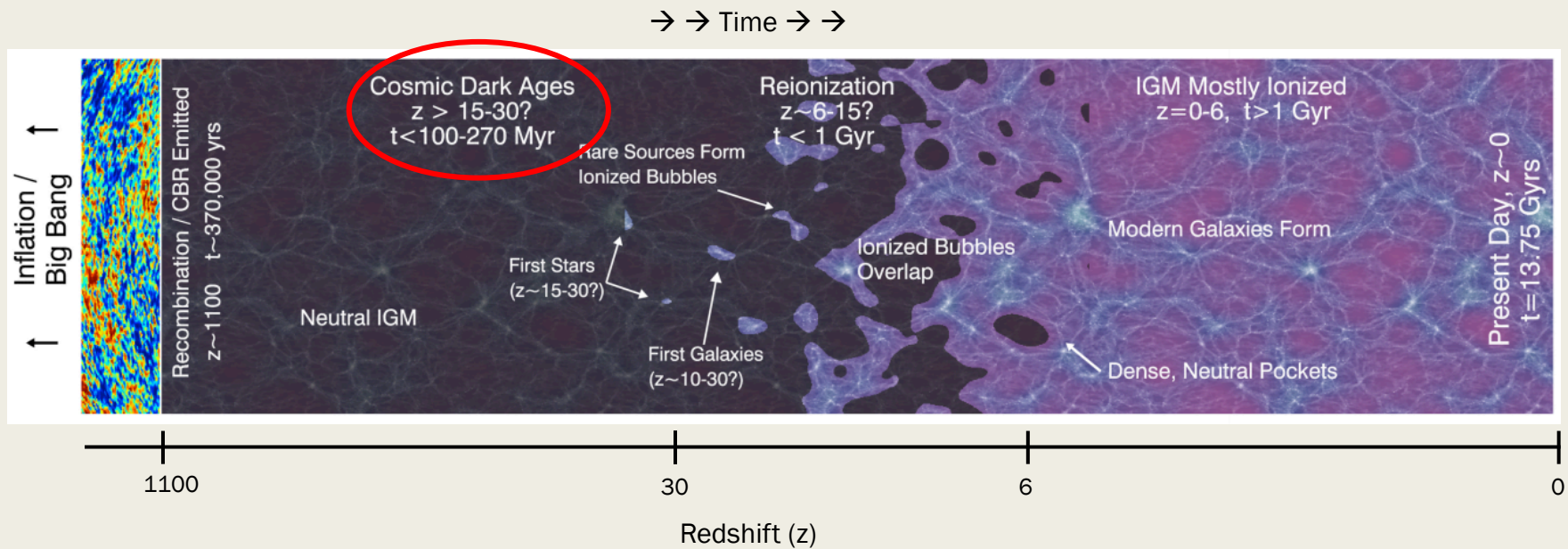
Robertson et al. (2010)

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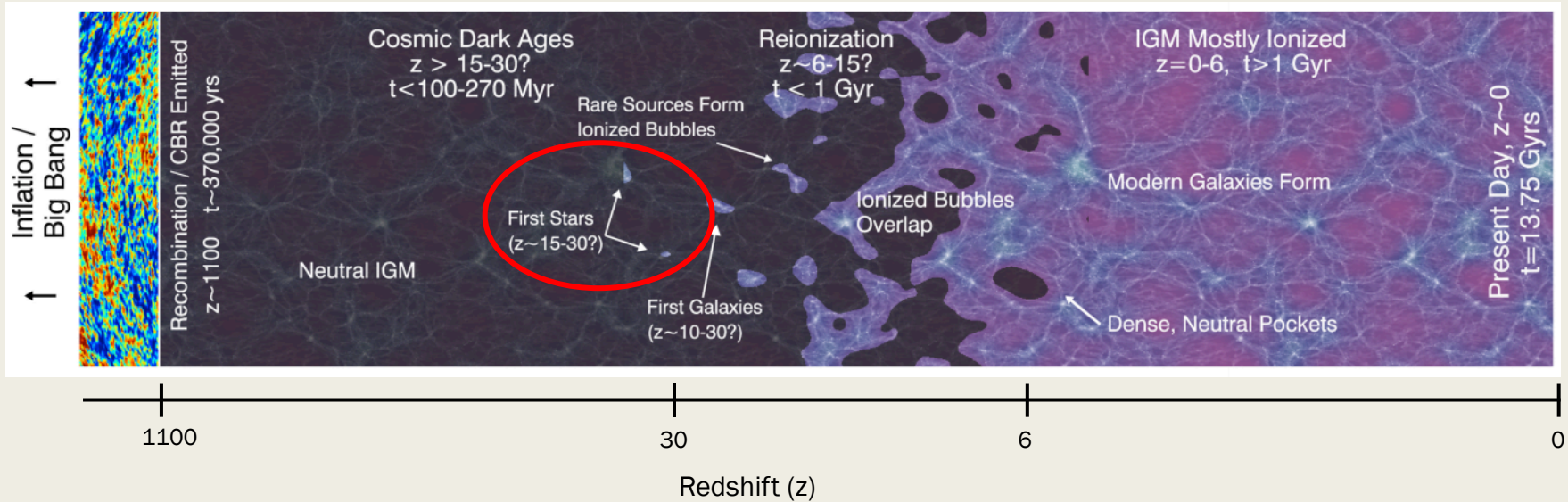
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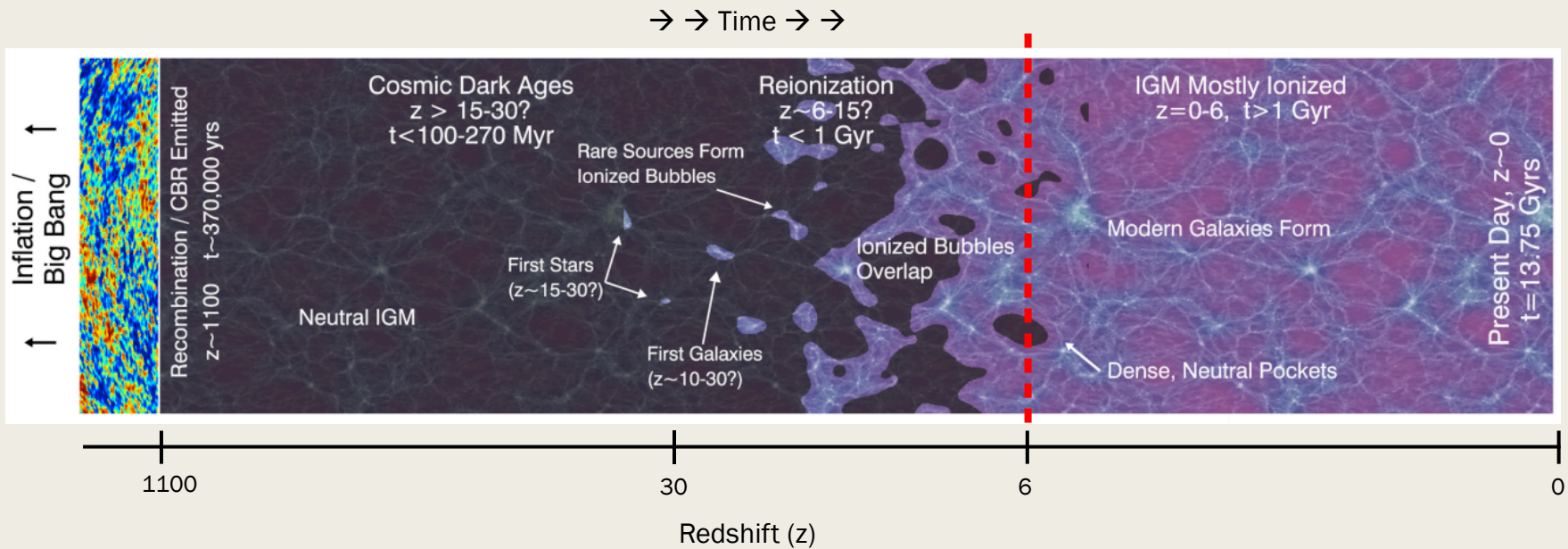
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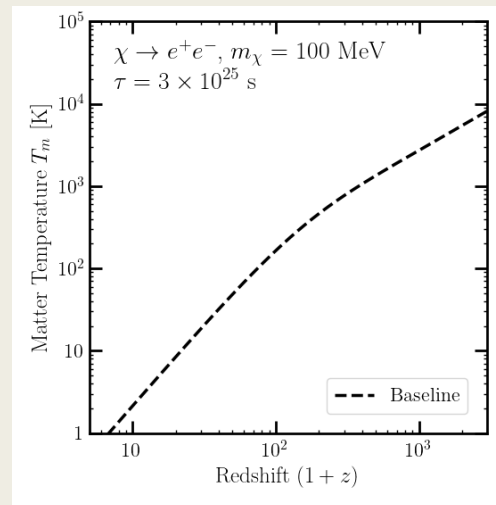
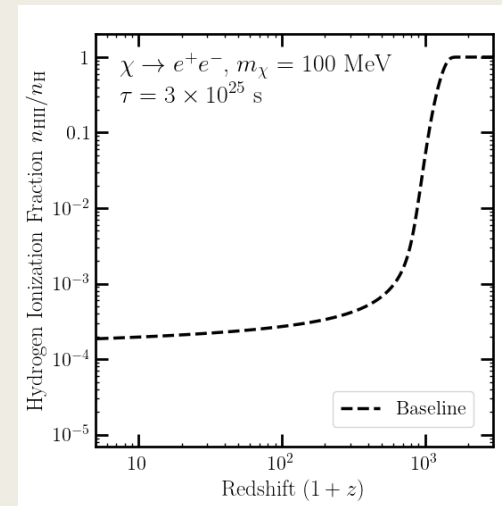
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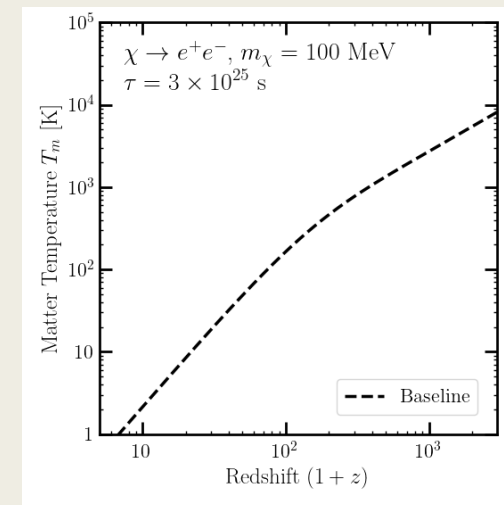
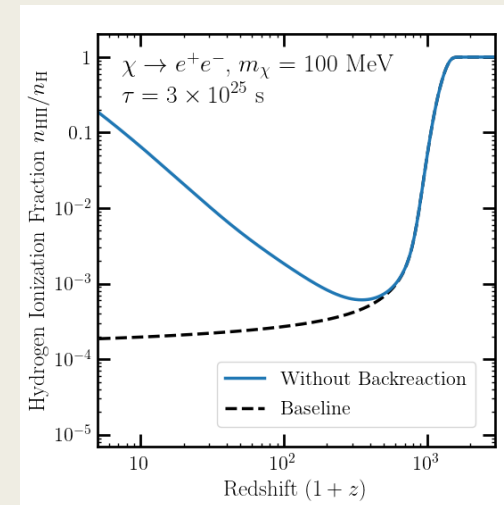
Cosmic heating and ionization

- If DM interacts with Standard Model particles, then decays/annihilations can inject energy into the IGM at early times
- Extra energy causes
 1. *Extra ionization: Detectable in the CMB power spectrum*
 2. *Extra heating: can probe with Lyman- α forest measurements*



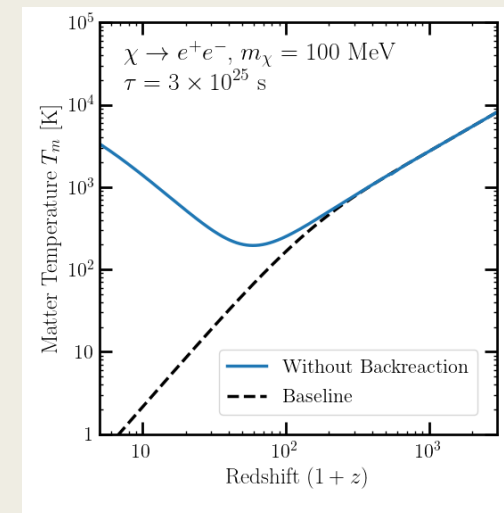
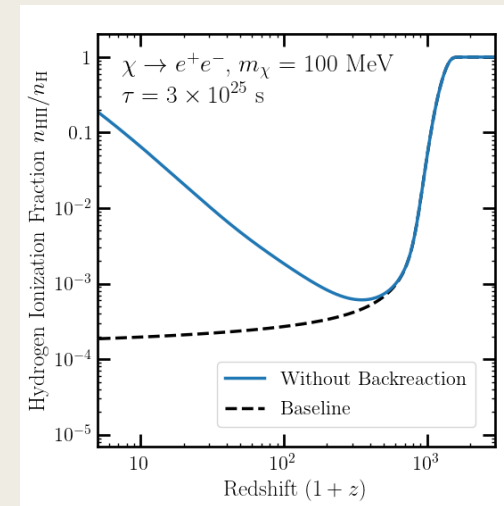
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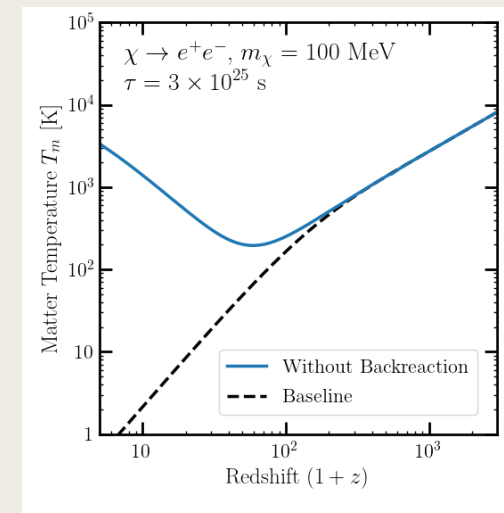
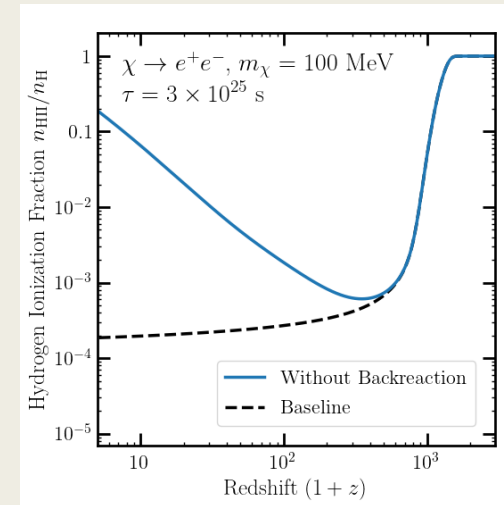
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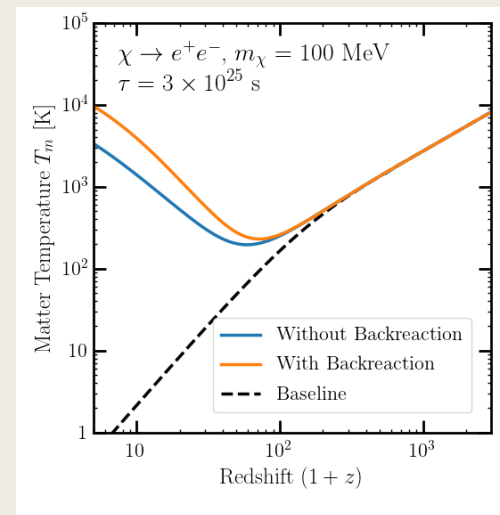
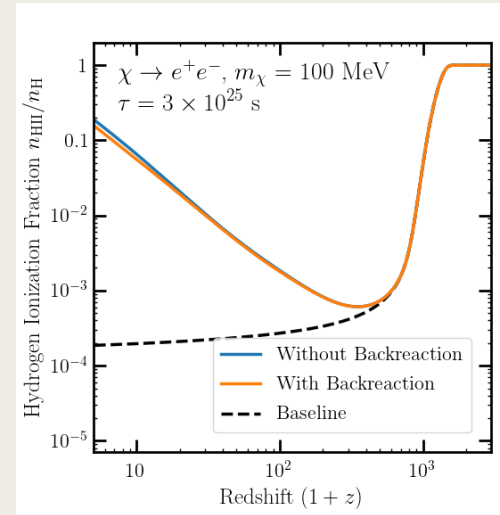
DarkHistory

- Python code package
 - Download at <https://github.com/hongwanliu/DarkHistory/>
- Simultaneously solves for evolution of free electron fraction and gas temperature, including exotic sources of energy injection
- Accounts for ‘backreaction’, where changes in ionization/temperature modify subsequent energy-loss processes
- Model-independent




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Evolution of H ionization

$$\dot{x}_{\text{HII}} = \dot{x}_{\text{HII}}^{\text{atom}} + \dot{x}_{\text{HII}}^{\text{DM}} + \dot{x}_{\text{HII}}^{\star}$$



Recombination,
collisional ionization

DM energy injection

Astrophysical sources,
use Planck constraints

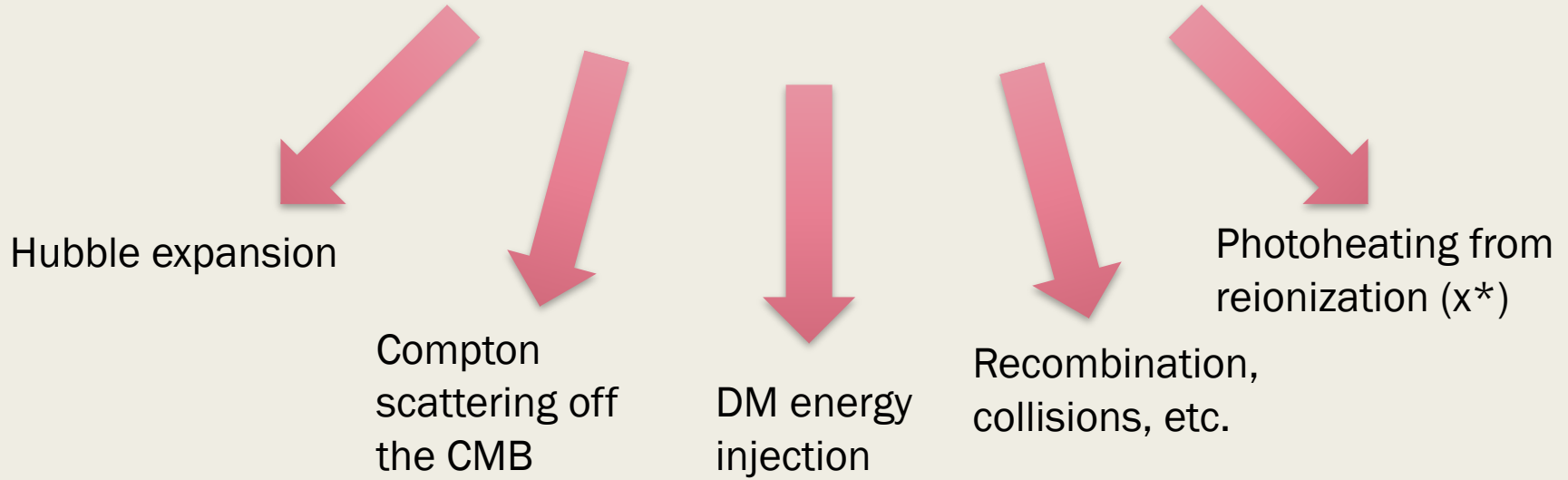
Evolution of He ionization

$$\dot{x}_{\text{HeII}} = \dot{x}_{\text{HeII}}^{\text{atom}} + \dot{x}_{\text{HeII}}^{\text{DM}} + \dot{x}_{\text{HeII}}^{\star}$$
$$x_{\text{HeIII}} = 0$$

- Analogous equation for Hell
- Assume there is no HeIII, which is a good approximation before Hell reionization

Evolution of temperature

$$\dot{T}_m = \dot{T}_{\text{adia}} + \dot{T}_C + \dot{T}_{\text{DM}} + \dot{T}_{\text{atom}} + \dot{T}^*$$



Photoheating model

How to treat photoheating rate?

- ‘Conservative’: Include no photoheating at all.
- ‘Photoheated’: Use a two-stage reionization model
 1. *Before reionization is complete, assume photoheating rate is proportional to photoionization rate \dot{x}^* ; parametrized by ΔT*
 - Restrict $\Delta T > 0$ K (‘photoheated-I’)
 - Restrict $\Delta T > 2 \times 10^4$ K (‘photoheated-II’), based on analytical arguments + simulations ¹
 2. *After reionization, assume IGM is in photoionization equilibrium \rightarrow gives calculable heating rate ²*

¹ Miralda-Escudé and Rees (1994), McQuinn (2012)

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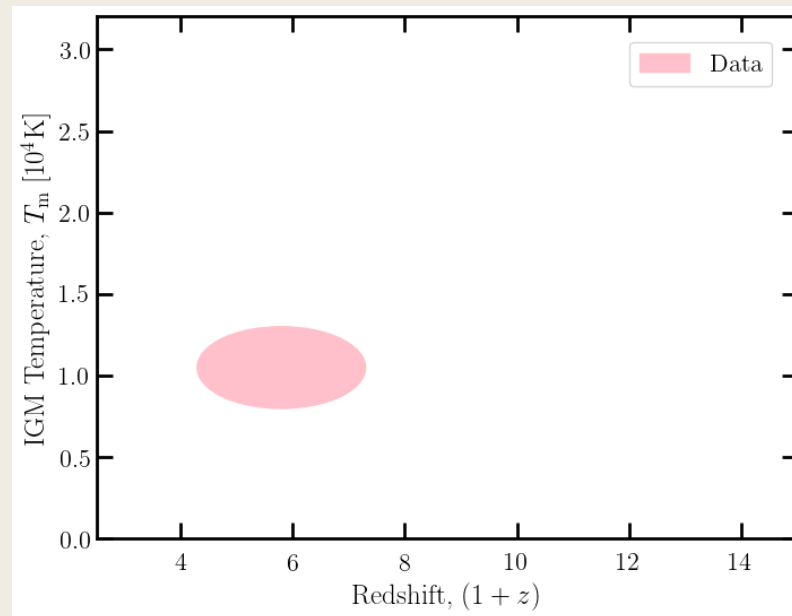
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Example histories

- For ‘conservative’ method, use modified chi-squared
 - *Only penalize model where overheats*
 - *Most robust because adding photoheating will only increase IGM temperature, producing stronger constraints*
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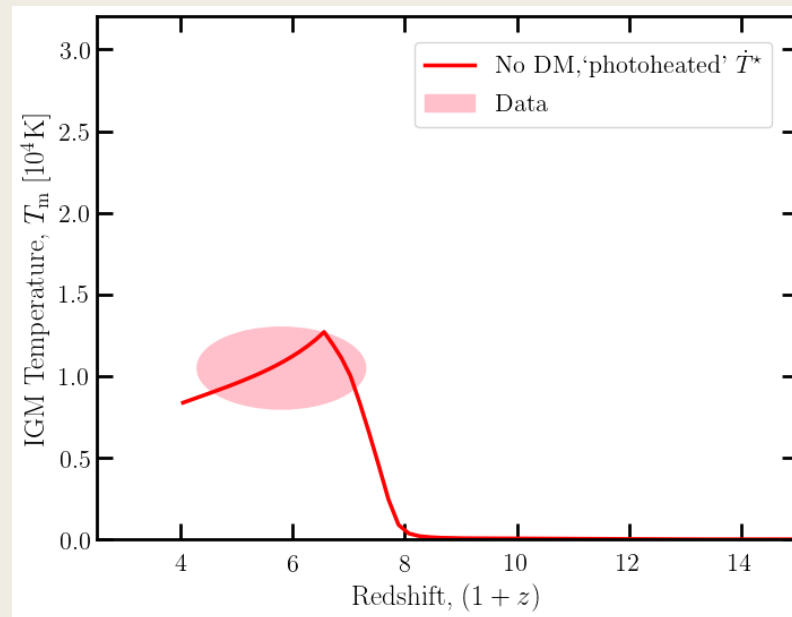
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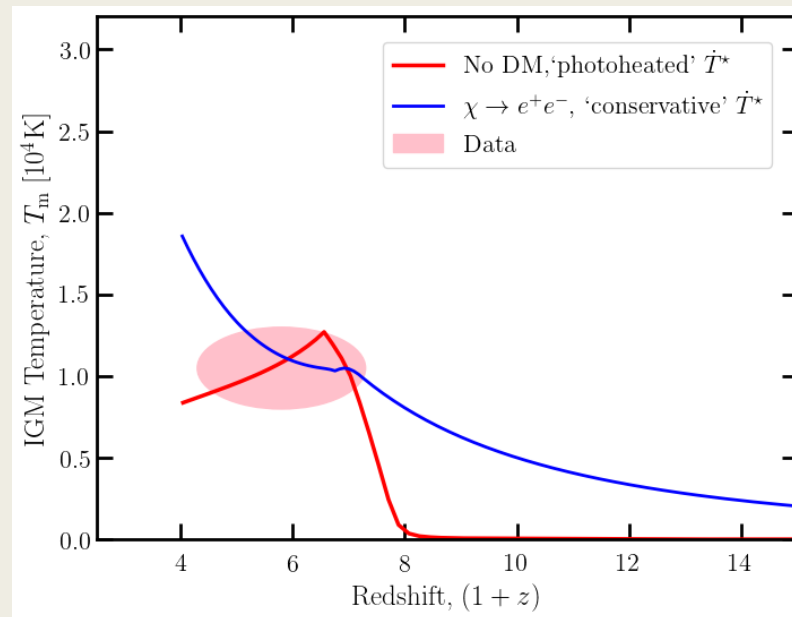
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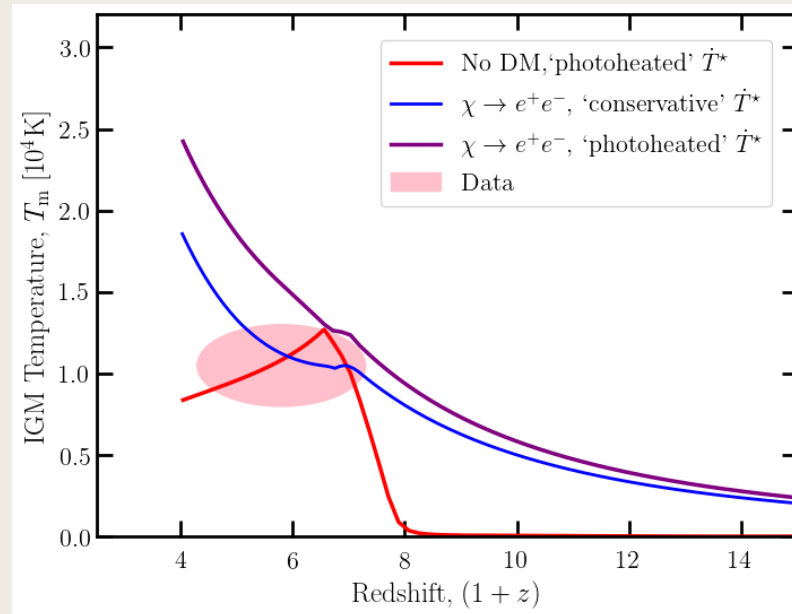
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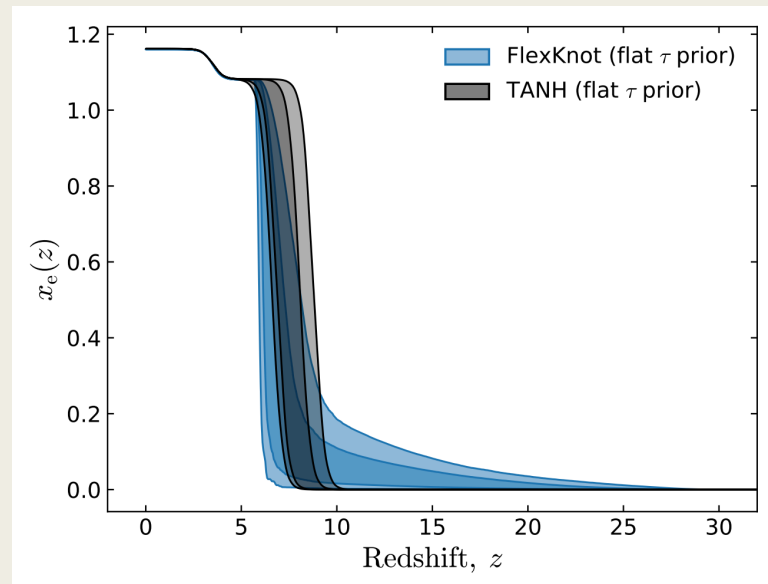


DATA



Reionization sources

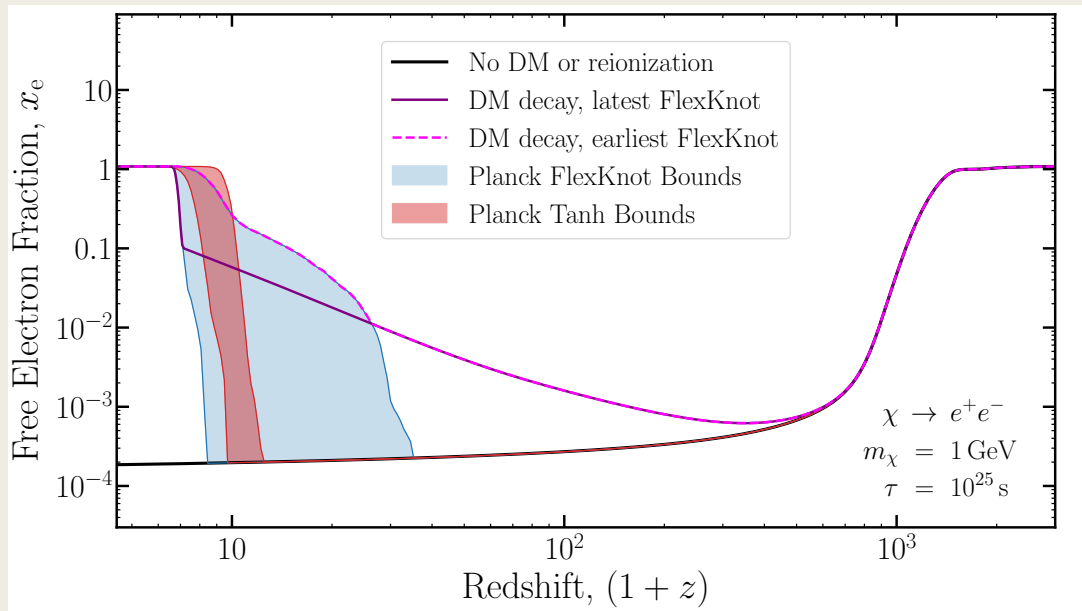
- Planck 2018: Constrains free electron fraction using either a ‘Tanh’ or ‘FlexKnot’ model
- DM produces extra ionization at early times, too little at late times—need astrophysical sources *
- Instead of modeling the astrophysics, match the ionization due to DM at early times onto the Planck reionization histories at later times



Planck (2018)

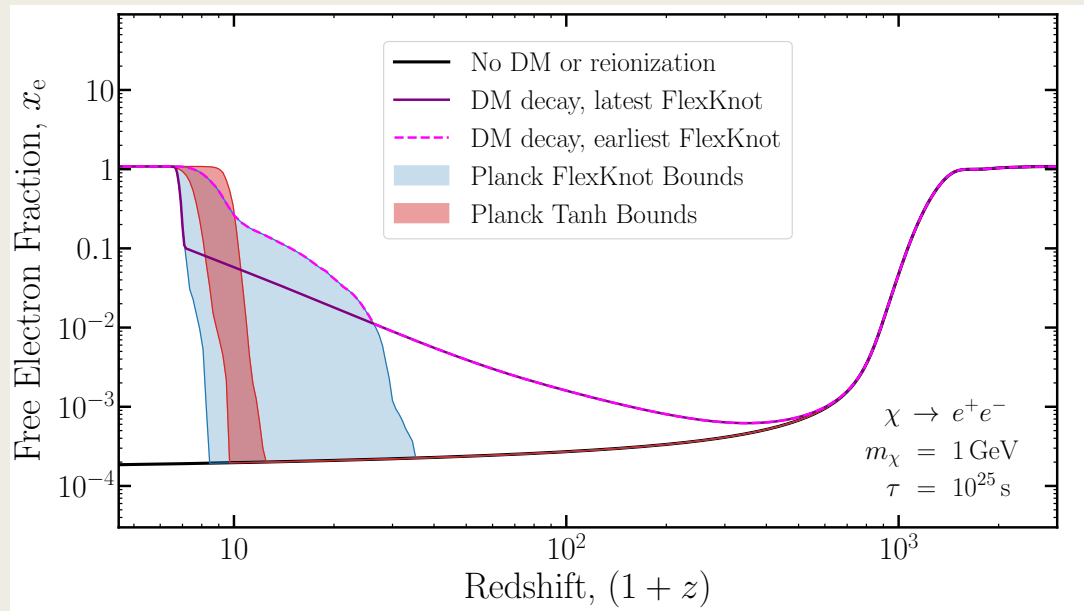
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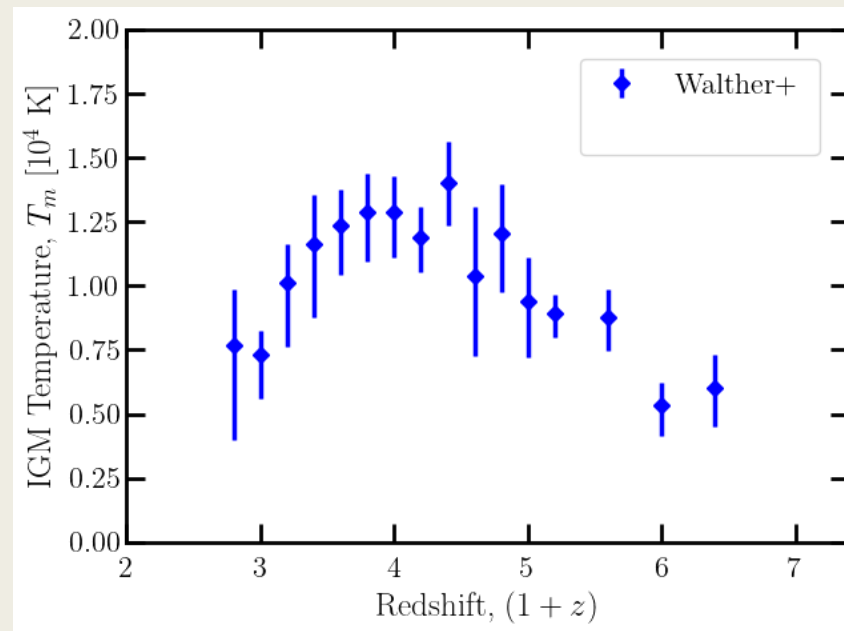
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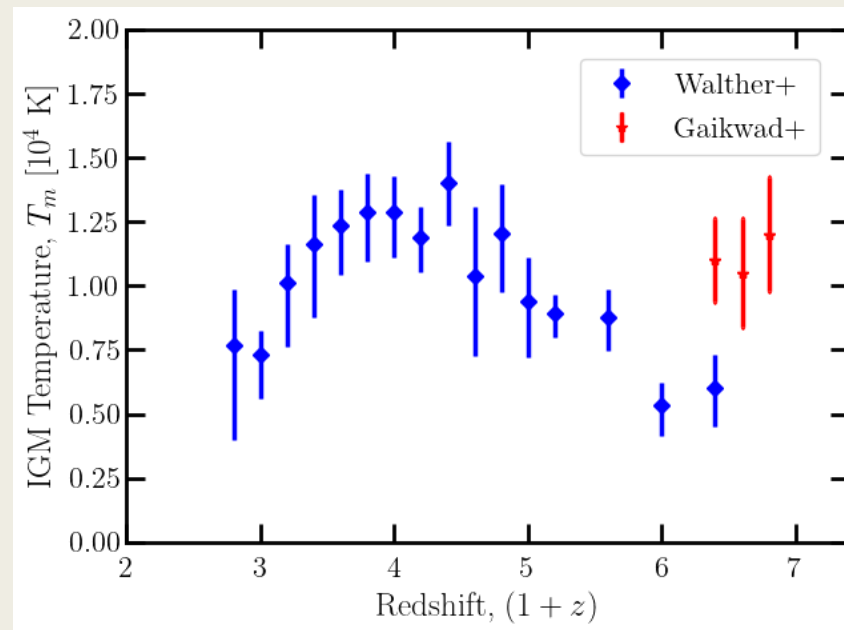
Temperature measurements

- Walther et al. (2019): Compared measured Ly α forest power spectra to hydrodynamical simulations to infer thermal evolution of IGM
- Gaikwad et al. (2020): Fit Ly α transmission spikes to simulations results



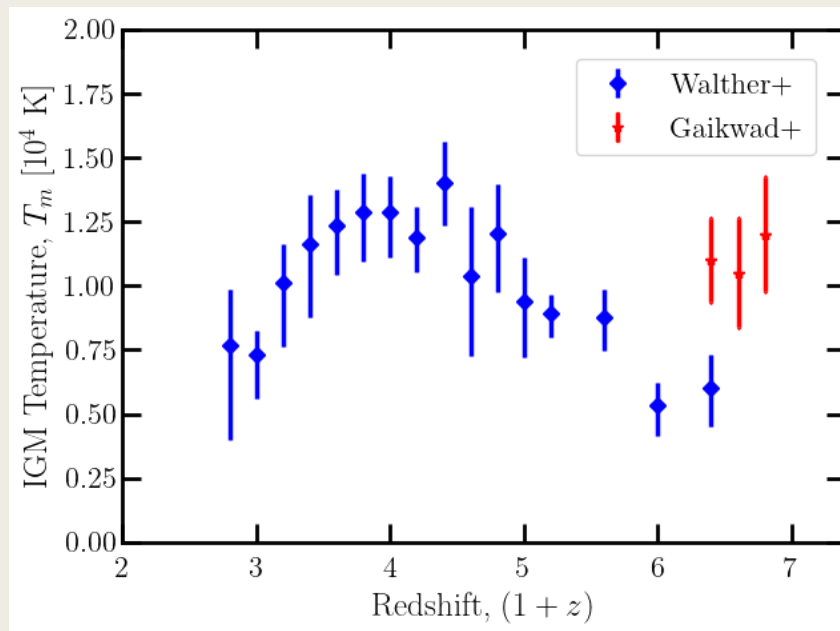
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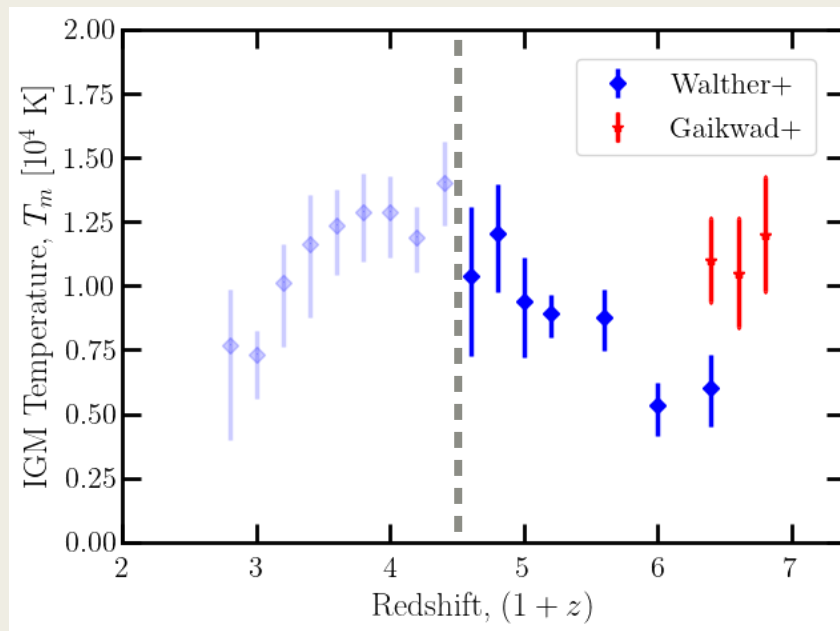
Temperature measurements

- Combined datasets:
 - *DarkHistory is not yet equipped to deal with Hell reionization so we only consider data points at $1+z > 4.6$*
 - *At $1+z \sim 6-7$, the two data sets are in tension; we discard the Walther+ results which could artificially strengthen our constraints*



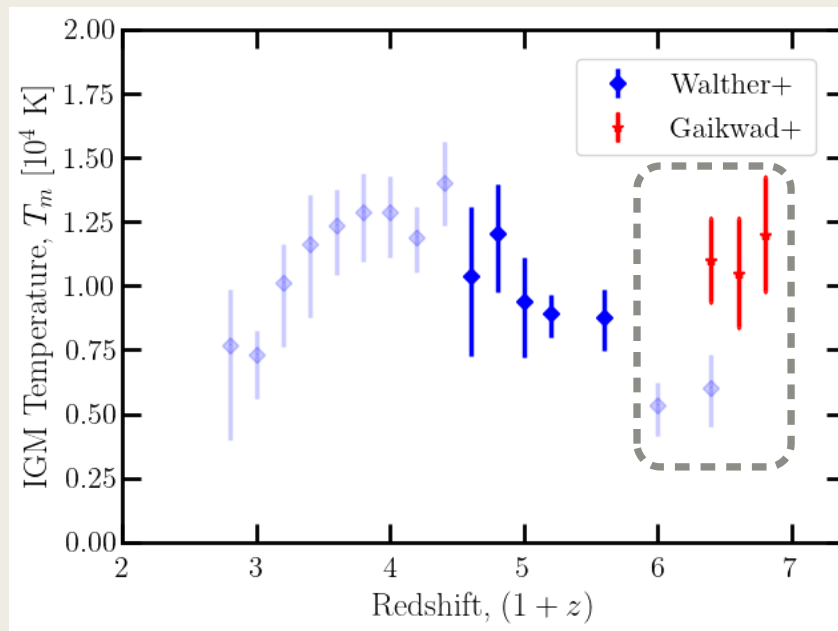
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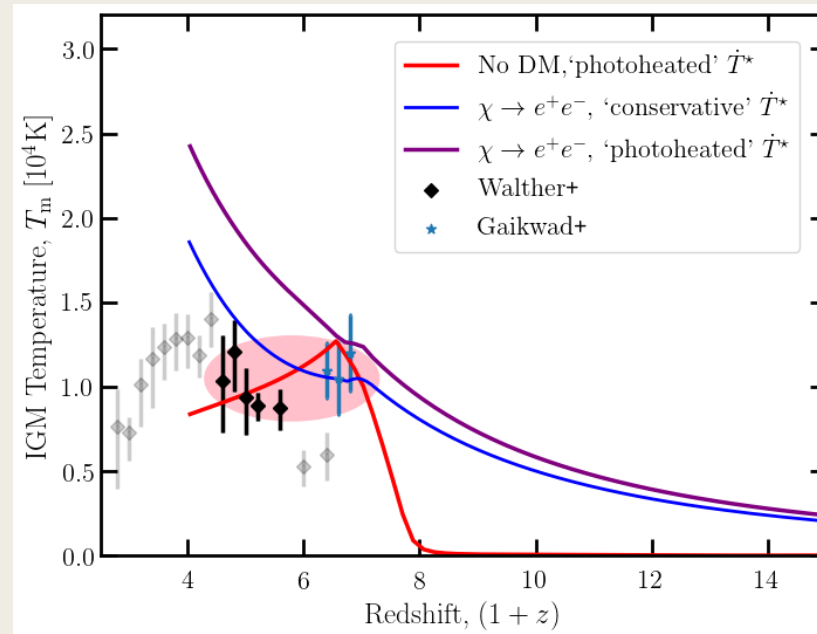
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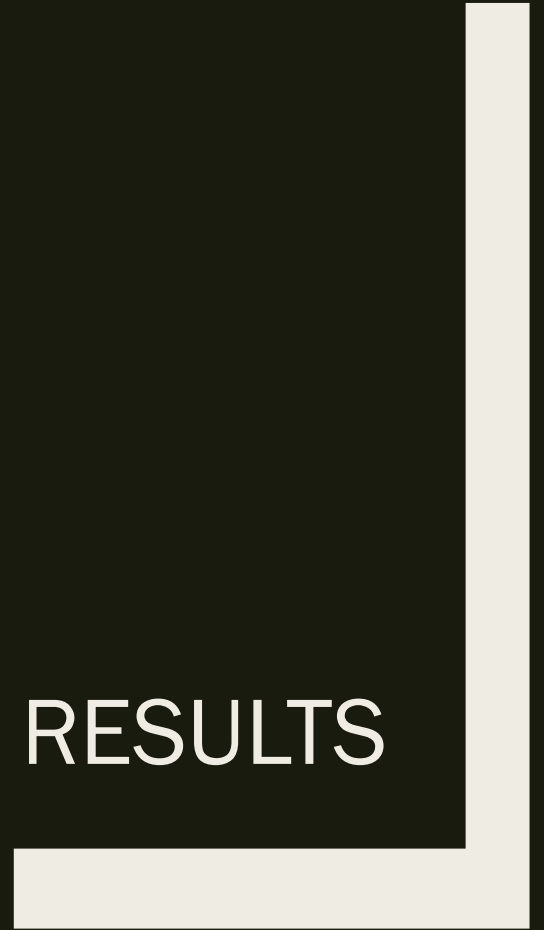


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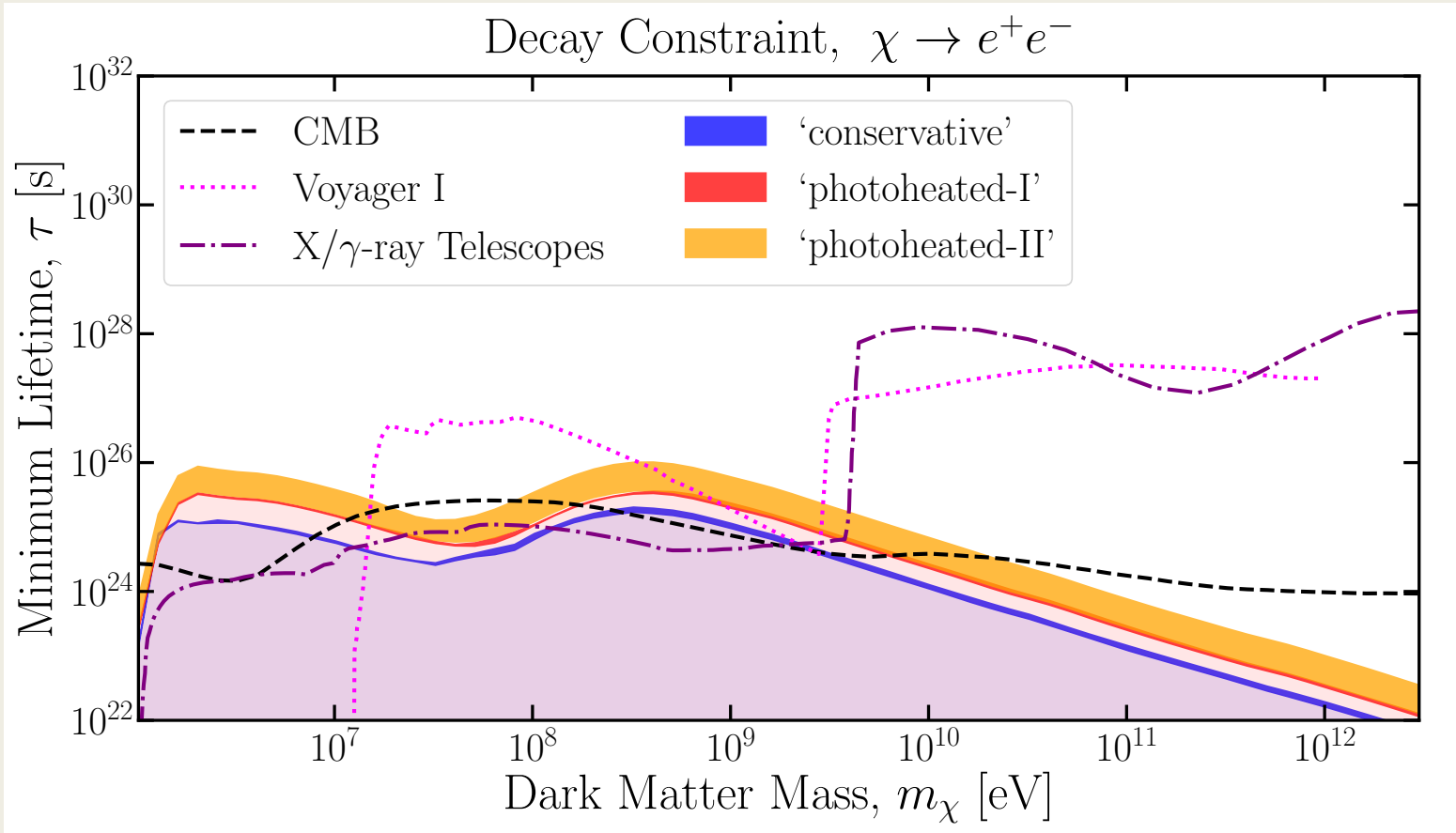
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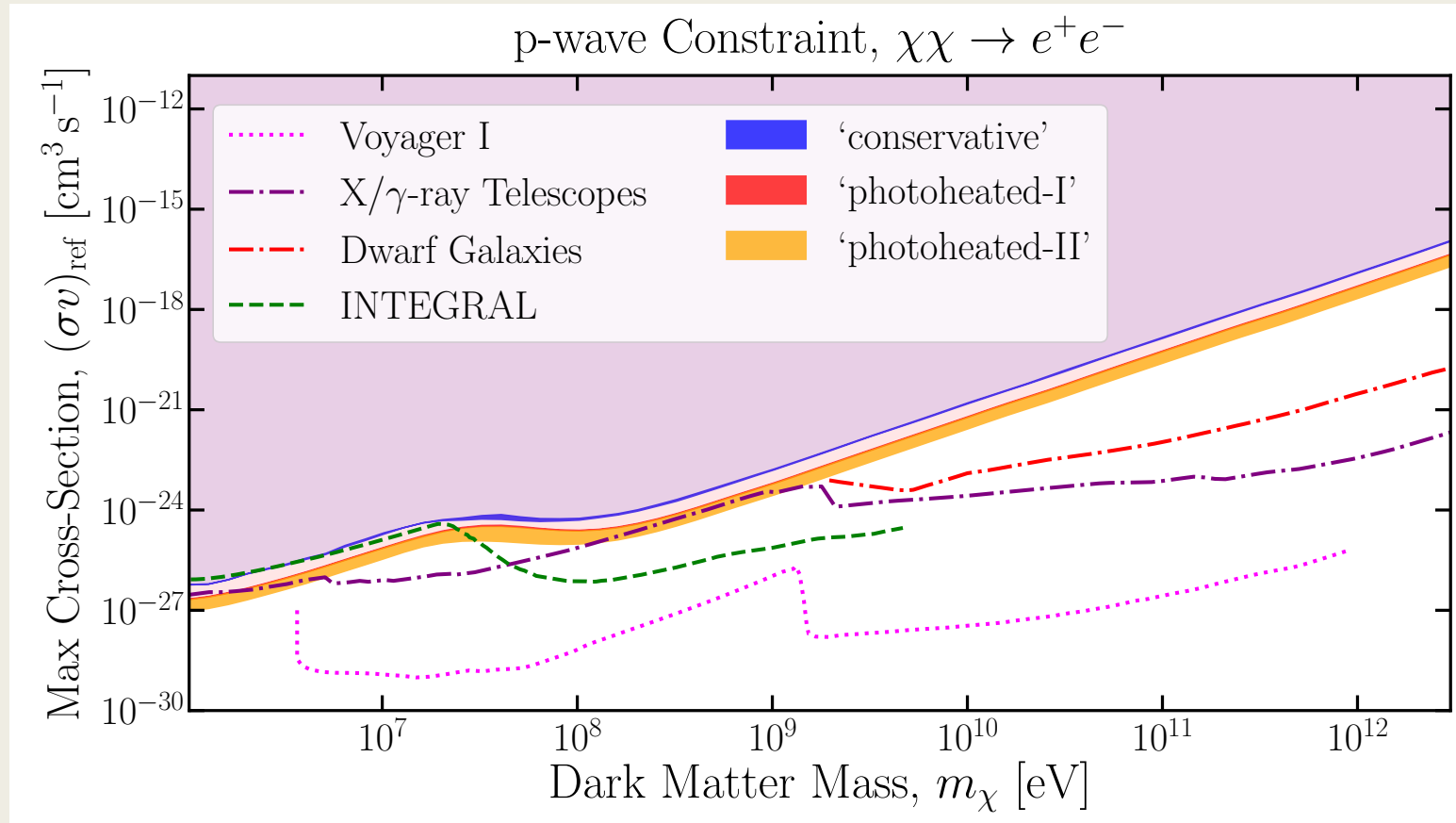
RESULTS



Decay to e^+e^-



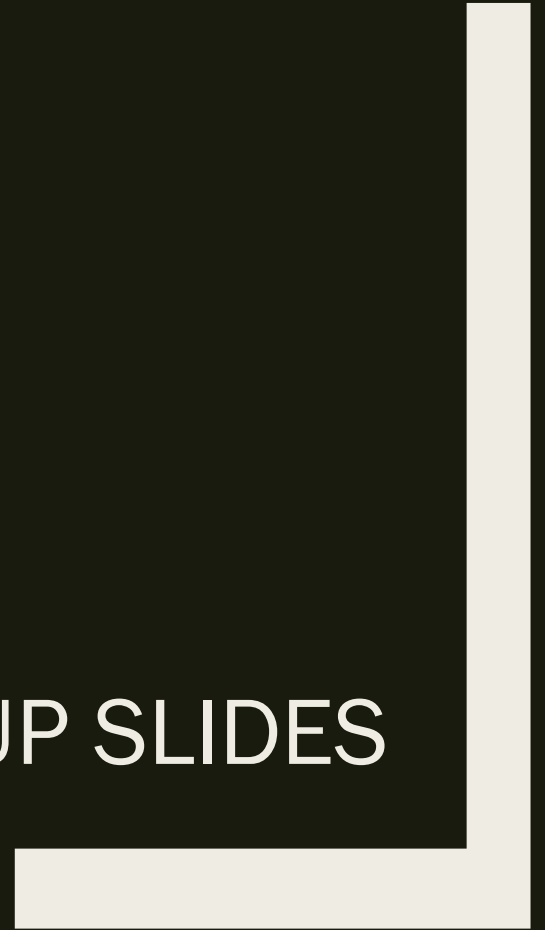
p-wave annihilation to e^+e^-



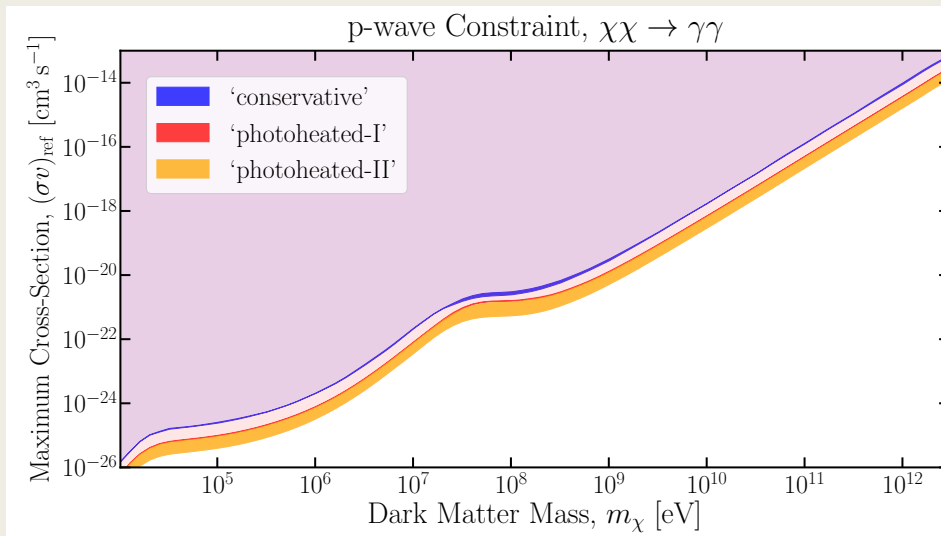
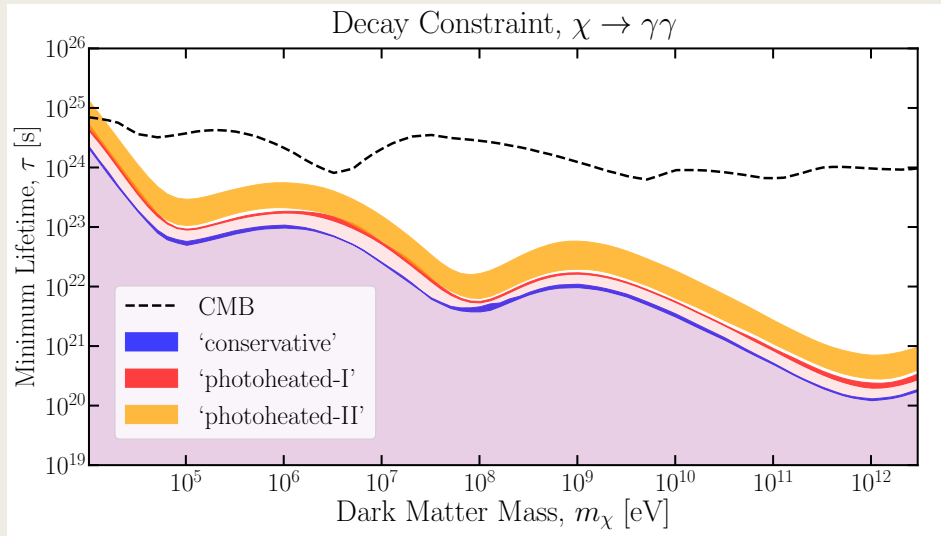
Summary

- We can self-consistently construct ionization and IGM temperature histories...
 - *in the presence of DM energy injection by using `DarkHistory` to solve evolution equations*
 - *in the presence of reionization sources by using Planck measurements*
- We use Ly α measurements of IGM temperature to constrain DM decay and p-wave annihilation
 - *'Conservative' constraints assume no photoheating*
 - *'Photoheated' constraints use two-stage model that strengthens limits by factor of 2-8*
- Could similarly use future 21cm observations for $z \sim 20$
- As uncertainties on IGM temperature measurements shrink, and reionization and photoheating models become more constrained, limits could strengthen considerably

BACKUP SLIDES



Photon constraints



Photoheating model: Stage 2

- ‘Photoheated’: Use a two-stage reionization model
 1. *Before reionization is complete, assume photoheating is proportional to photoionization; parametrized by ΔT*
 2. *After reionization, assume IGM is in photoionization equilibrium*
 - Parametrized by α_{bk} (defined by $J_\nu \propto \nu^{-\alpha_{bk}}$)
 - Restrict $-0.5 < \alpha_{bk} < 1.5$

$$\dot{T}^* = \begin{cases} \dot{x}_{\text{HII}}^* (1 + \chi) \Delta T, & x_{\text{HII}} < 0.99 \\ \sum_i \frac{E_i x_i}{3(\gamma_i - 1 + \alpha_{bk})} \alpha_{A,i} n_{\text{H}}, & x_{\text{HII}} \geq 0.99 \end{cases}$$



Sum is over HI and Hel (not Hell).

E_i is the ionization potential,

γ_i is the power-law index for photoionization cross section at threshold,

α_{Ai} is the case-A recombination coefficient.

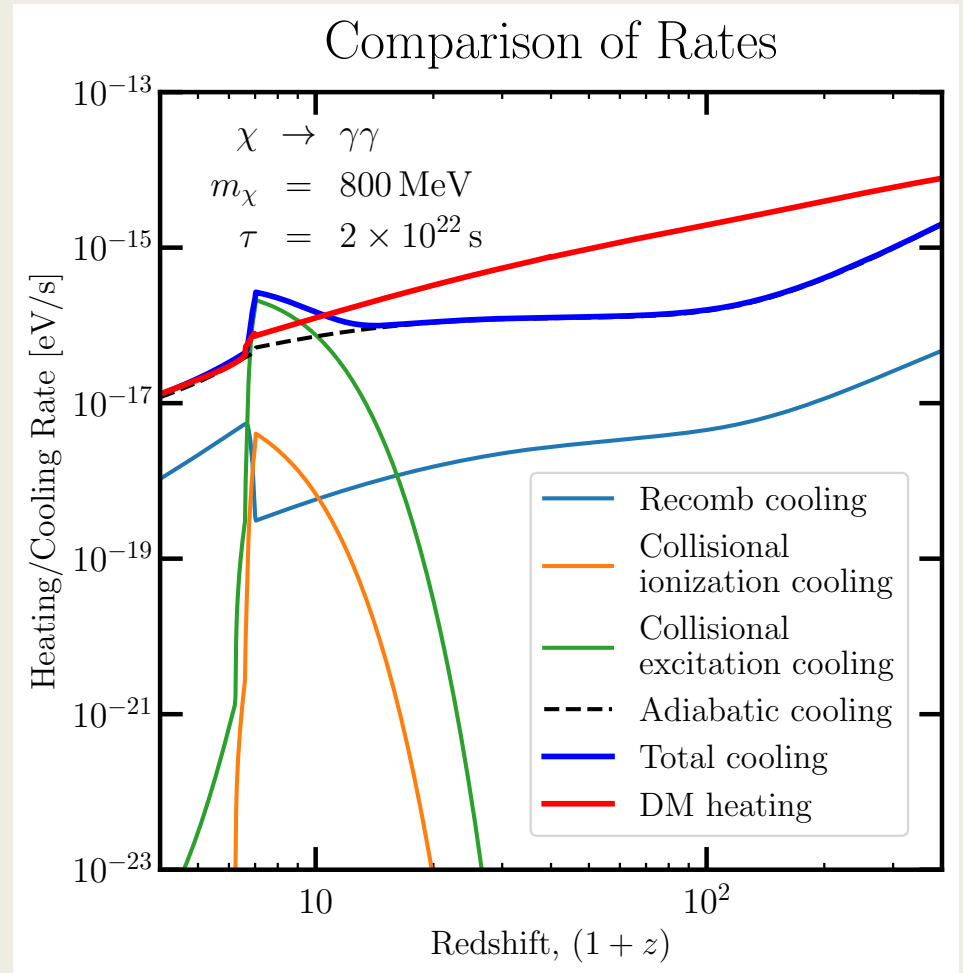
Heating/cooling terms

$$\dot{T}_{\text{adia}} = -2HT_{\text{m}}$$

$$\dot{T}_{\text{C}} = -\Gamma_{\text{C}}(T_{\text{CMB}} - T_{\text{m}})$$

$$\Gamma_{\text{C}} = \frac{x_{\text{e}}}{1 + \chi + x_{\text{e}}} \frac{8\sigma_{\text{T}}a_{\text{r}}T_{\text{CMB}}^4}{3m_{\text{e}}}$$

$$\dot{T}_{\text{DM}} = \frac{2f_{\text{heat}}(z, \mathbf{x})}{3(1 + \chi + x_{\text{e}})n_{\text{H}}} \left(\frac{dE}{dV dt} \right)^{\text{inj}}$$



Ionization terms

- Atomic processes

- Before z^* , case-B scenario

$$\dot{x}_{\text{HII}}^{\text{atom}} = 4 \mathcal{C}_{\text{H}} \left[(1 - x_{\text{HII}}) \beta_{\text{H}}^B e^{-E_{\text{H}}/T_{\text{CMB}}} - n_{\text{H}} x_{\text{e}} x_{\text{HII}} \alpha_{\text{H}}^B \right]$$

$$\dot{x}_{\text{HeII}}^{\text{atom}} = 4 \sum_s \mathcal{C}_{\text{HeII},s} \left[g_s (\chi - x_{\text{HeI}}) \beta_{\text{HeI},s}^B e^{-E_{\text{HeI},s}/T_{\text{CMB}}} - n_{\text{H}} x_{\text{e}} x_{\text{HeII}} \alpha_{\text{HeI},s}^B \right]$$

- After $z < z^*$, case-A scenario

$$\dot{x}_{\text{HII}}^{\text{atom}} = n_{\text{H}} (1 - x_{\text{HII}}) x_{\text{e}} \Gamma_{\text{eHI}} - n_{\text{H}} x_{\text{e}} x_{\text{HII}} \alpha_{\text{HII}}^A$$

- Dark matter energy injection

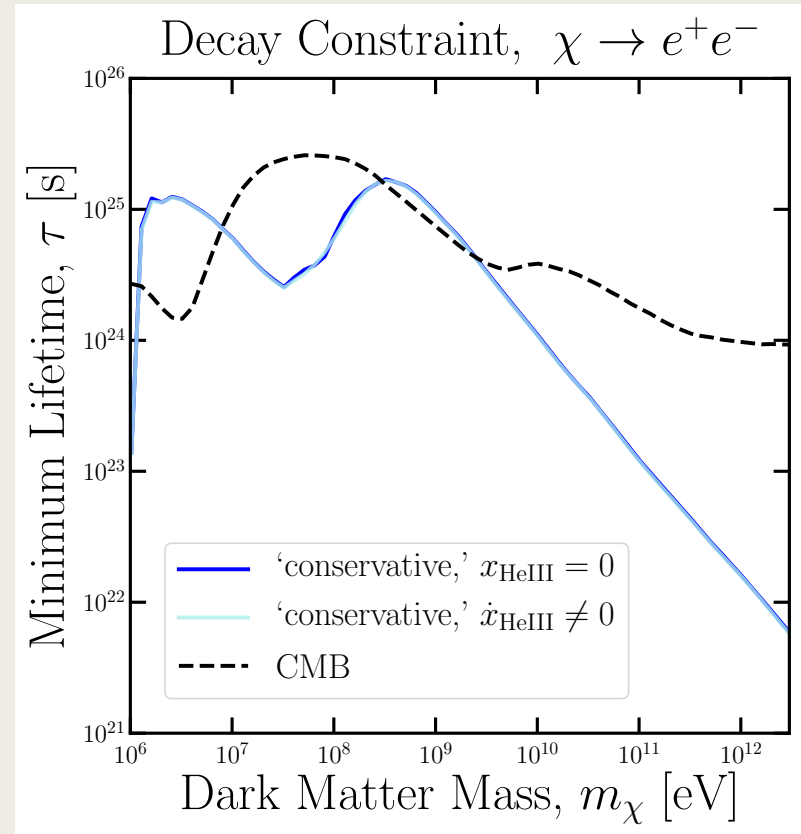
$$\dot{x}_{\text{HII}}^{\text{DM}} = \left[\frac{f_{\text{H ion}}(z, \mathbf{x})}{E_{\text{H}} n_{\text{H}}} + \frac{(1 - \mathcal{C}_{\text{H}}) f_{\text{exc}}(z, \mathbf{x})}{0.75 E_{\text{H}} n_{\text{H}}} \right] \left(\frac{dE}{dV dt} \right)^{\text{inj}}$$

$$\dot{x}_{\text{HeII}}^{\text{DM}} = \frac{f_{\text{He ion}}(z, \mathbf{x})}{E_{\text{HeI}} n_{\text{He}}} \left(\frac{dE}{dV dt} \right)^{\text{inj}}$$

$$\dot{x}_{\text{HeIII}}^{\text{DM}} = 0$$

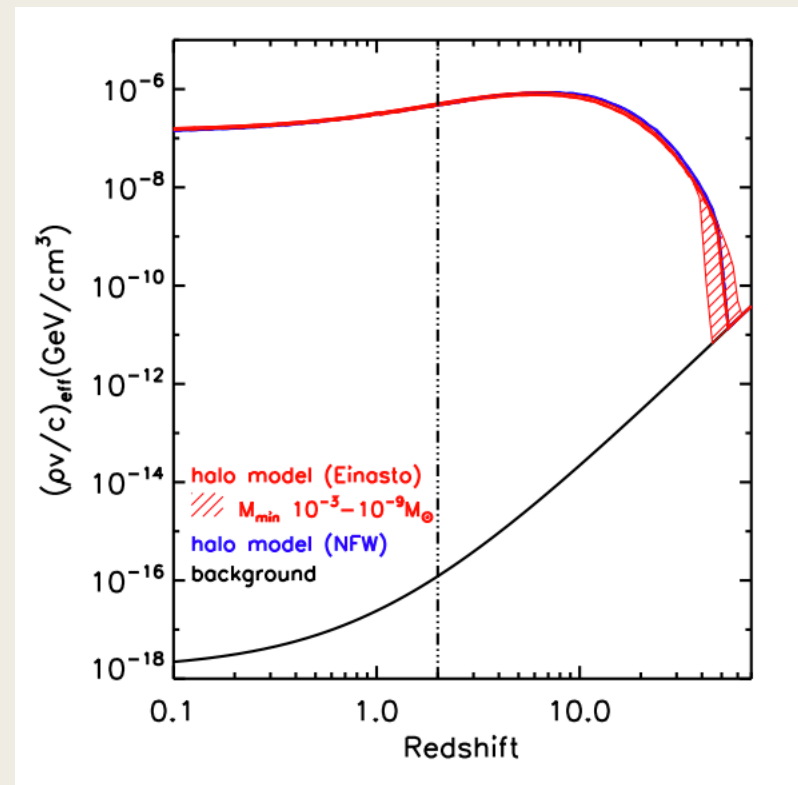
Effects of including HeIII

- Currently, DarkHistory assumes there is never HeIII
- However, still allows energy to be deposited into the IGM by
 1. Photoionization of $\text{HeII} \rightarrow \text{HeIII} + e^-$
 2. Resulting electron thermalizes with IGM
 3. Does NOT keep track of change to HeII and HeIII fractions
- This treatment is not self-consistent; however, the effect is small



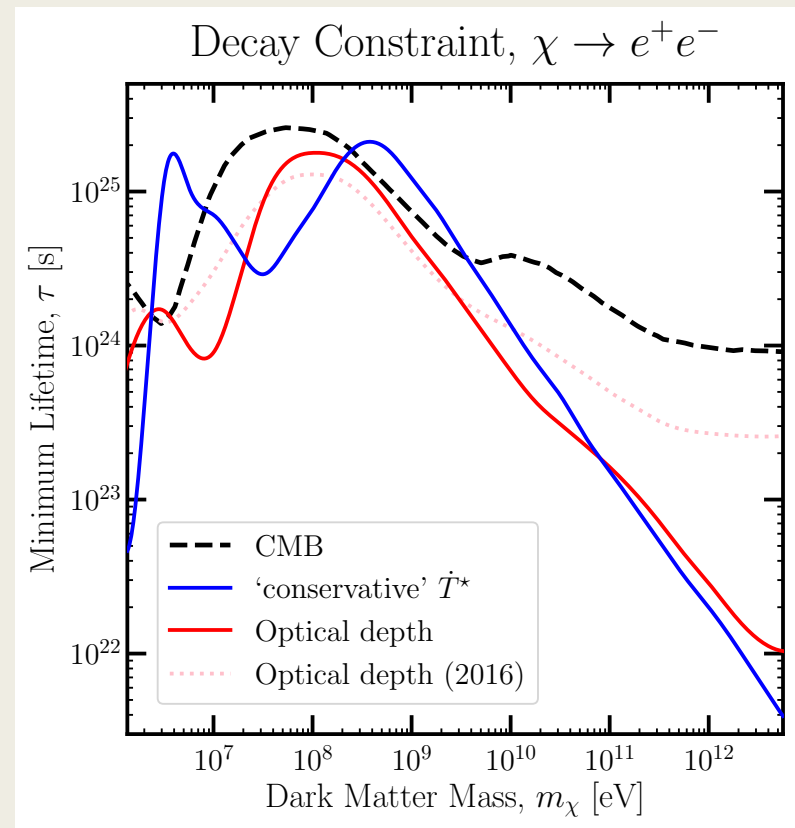
Uncertainty from p-wave boost factor

- For p-wave, main contribution to boost factor is from largest halos
 - *Well studied by N-body simulations*
- Uncertainty on boost factor is small, hence uncertainty on our constraints is small,
 - *Results in < 0.5% variation in our results*



Comparison to optical depth constraints

- Are the histories along our constraints excluded by their contribution to the optical depth?
 - *Planck 2018: For tanh function reionization,*
 $\tau = 0.0519^{+0.0030}_{-0.0079}$
- Temperature constraints are comparable to ionization constraints, and stronger in many places



Test statistics

- For the ‘conservative’ method, we use a modified chi-squared that only penalizes data points that overheat the IGM
- Can also think of this as a standard chi-squared test with a flexible background heating model
 - *Model contributes a non-negative amount to the temperature at each measured redshift*
 - In bins where contribution from DM is under measurement, model is set to match data exactly
 - In bins where contribution from DM exceeds measurement, model is fixed to 0
 - *# of model parameters = # of bins where DM is under is under measurement*
 - *Hence, degrees of freedom = # of bins where DM exceeds measurement*