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# Universality of the Cold Dark Matter Power Spectrum

Co-workers:

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• Arises before re-expansion of structures



\* paper in prep.

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If you have not seen our talk **Kinetic Field Theory: An Introduction** I recommend to watch this first, before continuing

• Density fluctuation power spectrum in KFT

$$\mathcal{P}_{\delta}(k,t) = \hat{\rho}\hat{\rho}(1+\mathrm{i}\hat{S}_{I}+...)Z_{0}[\mathbf{J},\mathbf{K}]_{\mathbf{J}=0=\mathbf{K}}$$

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• Solution to Hamiltonian e.o.m.

$$\vec{q}(t) = \vec{q}^{(i)} + g_{qp}(t,0)\vec{p}^{(i)} + \int_0^t dt' g_{qp}(t,t')\vec{f}(t')$$

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• Linear growth on large scales

$$g_{qp}(t,t') = D_{+}(t) - D_{+}(t')$$

$$\lim_{t \to 0} \text{ linear growth factor}$$

• Gaussian initial momentum correlations

$$\hat{C}_{pp}(\Delta \vec{q}) := \langle \vec{p}(\vec{q}) \otimes \vec{p}(\vec{q} + \Delta \vec{q}) \rangle = \int_{k} \frac{\vec{k} \otimes \vec{k}}{k^{4}} P_{\delta}^{(i)}(k) e^{i\vec{k}\Delta \vec{q}}$$
initial density fluctuation

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, power spectrum

# **CDM** Power spectrum

General expression

$$\mathcal{P}_{\delta}(k,t) = \langle \mathrm{e}^{-\mathrm{i}\vec{k}\cdot[\vec{q}_{1}(t)-\vec{q}_{2}(t)]} \rangle$$

### **CDM** Power spectrum

General expression

$$\mathcal{P}_{\delta}(k,t) = \langle \mathrm{e}^{-\mathrm{i}\vec{k}\cdot[\vec{q}_{1}(t)-\vec{q}_{2}(t)]} \rangle$$

For particles on straight Zel'dovich trajectories

$$\mathcal{P}_{\delta}(k,t) = e^{-\frac{\sigma_1^2}{3}g_{qp}^2(t,0)k^2} \int d^3q \left( e^{g_{qp}^2(t,0)\vec{k}^{\intercal}\hat{C}_{pp}(\vec{q})\vec{k}} - 1 \right) e^{i\vec{k}\cdot\vec{q}}$$
  
initial momentum correlations

### Large scale limit $k \to 0$

$$\mathcal{P}_{\delta}(k,t) = e^{-\frac{\sigma_1^2}{3}g_{qp}^2(t,0)k^2} \int d^3q \left( e^{g_{qp}^2(t,0)\vec{k}^{\mathsf{T}}\hat{C}_{pp}(\vec{q})\vec{k}} - 1 \right) e^{i\vec{k}\cdot\vec{q}}$$

Expansion for small modes  $\rightarrow$  linear power spectrum for large scales

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$$\mathcal{P}_{\delta}(k,t) \sim g_{qp}^2(t,0) \int \mathrm{d}^3 q \ \vec{k}^{\mathsf{T}} \hat{C}_{pp}(\vec{q}) \vec{k} \ \mathrm{e}^{\mathrm{i}\vec{k}\cdot\vec{q}}$$

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$$=g_{qp}^{2}(t,0)P_{\delta}^{(1)}(k), \quad \text{for } k \to 0$$

$$\mathcal{P}_{\delta}(k,t) = e^{-\frac{\sigma_1^2}{3}g_{qp}^2(t,0)k^2} \int d^3q \left( e^{g_{qp}^2(t,0)\vec{k}^{\mathsf{T}}\hat{C}_{pp}(\vec{q})\vec{k}} - 1 \right) e^{i\vec{k}\cdot\vec{q}}$$

Expansion of the exponent for small scales \*

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Expansion of the exponent for small scales  $^{\star}$ 

$$g_{qp}^{2}(t,0)\vec{k}^{\mathsf{T}}\hat{C}_{pp}(\vec{q})\vec{k} \sim g_{qp}^{2}(t,0)k^{2}\left[\frac{\sigma_{1}^{2}}{3} + \vec{q}^{\mathsf{T}}\hat{A}\vec{q} + O\left(q^{4}\right)\right]$$

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\* we gave a mathematical proof that this expansion correct (paper in prep.)

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with the Hessian

$$\hat{A} = \frac{\sigma_2^2}{15} \begin{pmatrix} 1 & 0 & 0\\ 0 & 1 & 0\\ 0 & 0 & 3 \end{pmatrix}$$

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with the Hessian  $\hat{A} = \frac{\sigma_2^2}{15} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 3 \end{pmatrix}$ 

and the moments of the initial power spectrum

$$\sigma_n^2 := \int \mathrm{d}k \ k^{2n-2} P_{\delta}^{(\mathrm{i})}(k)$$

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Expansion for large modes  $\rightarrow$  power law for small scales

$$\mathcal{P}_{\delta}(k,t) \sim \int \mathrm{d}^{3}q \, \mathrm{e}^{-g_{qp}^{2}(t,0)k^{2}\vec{q}^{\mathsf{T}}\hat{A}\vec{q}+\mathrm{i}\vec{k}\cdot\vec{q}}$$

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$$= \frac{1}{k^{3}} \frac{(2\pi)^{3/2}}{\sqrt{3}} \left(\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}\right)^{3/2} e^{-\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}}, \quad \text{for } k \to \infty$$

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$$= \frac{1}{k^{3}} \frac{(2\pi)^{3/2}}{\sqrt{3}} \left(\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}\right)^{3/2} \mathrm{e}^{-\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}}, \quad \text{for } k \to \infty$$

This result has also been found recently by **Chen and Pietroni** in the Lagrangian formulation of fluid mechanics arXiv:2002.11357

#### **Amplitude Time Evolution**

$$k^{3}\mathcal{P}_{\delta}(k,t) \sim \frac{(2\pi)^{3/2}}{\sqrt{3}} \left(\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}\right)^{3/2} e^{-\frac{5}{2g_{qp}^{2}(t,0)\sigma_{2}^{2}}}$$



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# Initial Hot Dark Matter $P_{\delta}^{(i)}(k) \propto k \cdot e^{-k}$

• Independent on the **shape** of the initial power spectrum

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# Thank you for watching my talk!

- Please visit also the talks by Elena Kozlikin and Carsten Littek
- I am looking forward to meet you in the discussion session!

Contact me

Email: Sara.Konrad@stud.uni-heidelberg.de LinkedIn: linkedin.com/in/sara-konrad-00a9551b3

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• Lagrangian Dynamics:  $\vec{r}(\vec{x},t)$ 

r(x, t)initial coordinate coordinate at time t

• Ansatz for intially rotation free velocity field

$$\vec{r}(\vec{x},t) = a(t)\vec{x} + b(t)\vec{p}(\vec{x})$$
  
initial momentum

• Being compatible with linear growth on large scales

$$\vec{r}(\vec{x},t) = a(t)\vec{x} + a(t)D_{+}(t)\vec{p}(\vec{x})$$
   
 
$$\linear growth factor$$

Gaussian initial momentum correlations

$$\hat{C}_{pp}(\vec{x} - \vec{y}) := \langle \vec{p}(\vec{x}) \otimes \vec{p}(\vec{y}) \rangle = \int_{k} \frac{\vec{k} \otimes \vec{k}}{k^{4}} P^{ini}(k) e^{i\vec{k}(\vec{x} - \vec{y})}$$
 initial density fluctuation power spectrum